

MODELING THE EARTH'S MAGNETIC FIELD BY LOCAL MULTISCALE METHODS

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Scaling- and Wavelet-Kernels:

$$\Phi_J(\xi, \eta) = \sum_{n=0}^{\infty} \sum_{k=0}^{2n+1} (\Phi_J)^{\wedge}(n) Y_{n,k}(\xi) Y_{n,k}(\eta),$$

$$\Psi_J(\xi, \eta) = \Phi_{J+1}(\xi, \eta) - \Phi_J(\xi, \eta).$$

Scaling- and Wavelet-Kernels:

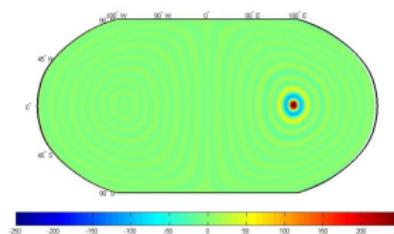
$$\Phi_J(\xi, \eta) = \begin{cases} \text{weak sing.,} & 1 - \xi \cdot \eta \geq 2^{-J}, \\ \text{regularization,} & 1 - \xi \cdot \eta < 2^{-J}, \end{cases}$$

$$\Psi_J(\xi, \eta) = \Phi_{J+1}(\xi, \eta) - \Phi_J(\xi, \eta).$$

Multiscale Representation:

$$F(\xi) = \int_{\Omega} \Phi_{J_0}(\xi, \eta) G(\eta) d\omega(\eta) + \sum_{J=J_0}^{\infty} \int_{\Omega} \Psi_J(\xi, \eta) G(\eta) d\omega(\eta).$$

Shannon
Wavelet ($J=4$):



1. Bayer, Freeden, Maier: J. Atm. Solar-Terr. Phys. 63 (2001)
2. Freeden, Schreiner: J. Geod. 78 (2006)
3. Holschneider, Chambodout, Mandea: Phys. Earth Plant. Int. 135 (2003)

Mie-Decomposition:

$$\begin{aligned} b &= p_b + q_b = \nabla \wedge L P_b + \textcolor{red}{L Q_b}, \\ j &= p_j + q_j = \nabla \wedge L P_j + L Q_j. \end{aligned}$$

Toroidal Scalar:

$$Q_b(x) = - \int_{\Omega} L_{\eta}^* G(\Delta^*; \xi \cdot \eta) \cdot b(r\eta) d\omega(\eta).$$

Helmholtz-Mie-Kombination:

$$j(x) = \xi \underbrace{\frac{\Delta_{\xi}^* Q_b(x)}{r}}_{=J_{rad}(x)} + \nabla_{\xi}^* \left(\frac{1}{r} \frac{\partial}{\partial r} (r Q_b(x)) \right) + L_{\xi}^* \Delta_x P_b(x), \quad r = |x|, \xi = \frac{x}{|x|}.$$

Mie-Decomposition:

$$b = p_b + q_b = \nabla \wedge L P_b + \textcolor{red}{L Q_b},$$

$$j = p_j + q_j = \nabla \wedge L P_j + L Q_j.$$

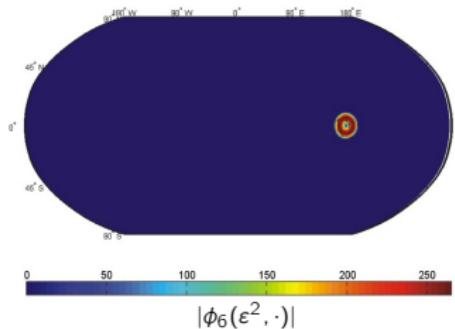
Toroidal Scalar:

$$\textcolor{red}{Q_b(x)} = - \int_{\Omega} L_{\eta}^* G(\Delta^*; \xi \cdot \eta) \cdot b(r\eta) d\omega(\eta).$$

Regularized Green Function:

$$G^J(\Delta^*; \xi \cdot \eta) = \begin{cases} \frac{1}{4\pi} \ln(1 - \xi \cdot \eta) + \frac{1}{4\pi} (1 - \ln(2)), & 1 - \xi \cdot \eta \geq 2^{-J}, \\ T_3^J(\xi \cdot \eta), & 1 - \xi \cdot \eta < 2^{-J}, \end{cases}$$

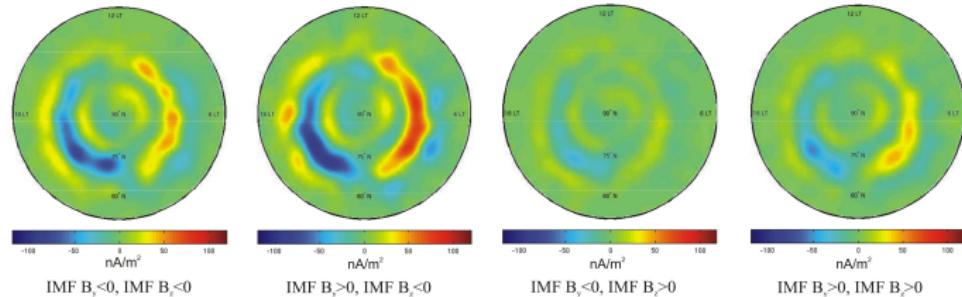
$$J_{rad}(x) = \lim_{J \rightarrow \infty} \int_{\Omega} \underbrace{\left(-\frac{1}{r} L_{\eta}^* \Delta_{\xi}^* G^J(\Delta^*; \xi \cdot \eta) \right)}_{= \phi_J(\xi, \eta)} \cdot b(r\eta) d\omega(\eta).$$



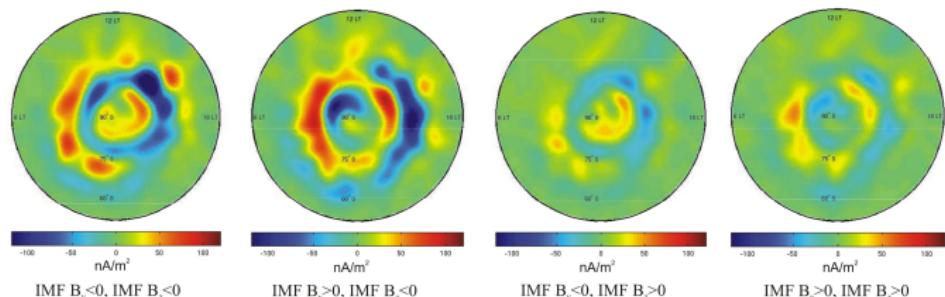
CHAMP Data: Jan. 2001 - Dec. 2007

Subtraction of CHAOS-3 model, Binning via Huber's M-estimation

N-Winter



S-Summer

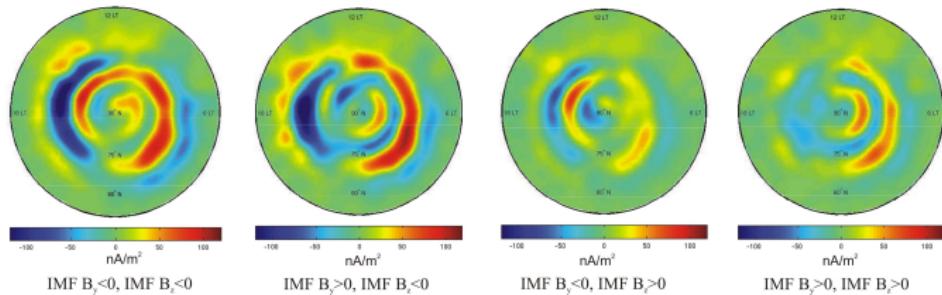


1. Gerhards: PhD Thesis, Geomathematics Group, TU Kaiserslautern
2. Papitashvili, et al.: Geophys. Res. Lett. 29 (2002)
3. Weimer: J. Geophys. Res. 106 (2001)

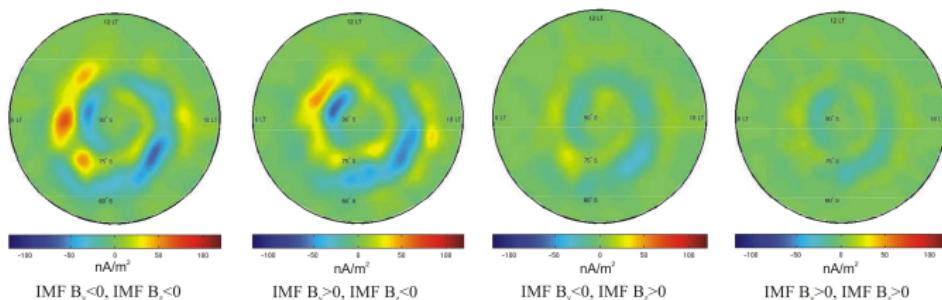
CHAMP Data: Jan. 2001 - Dec. 2007

Subtraction of CHAOS-3 model, Binning via Huber's M-estimation

N-Summer



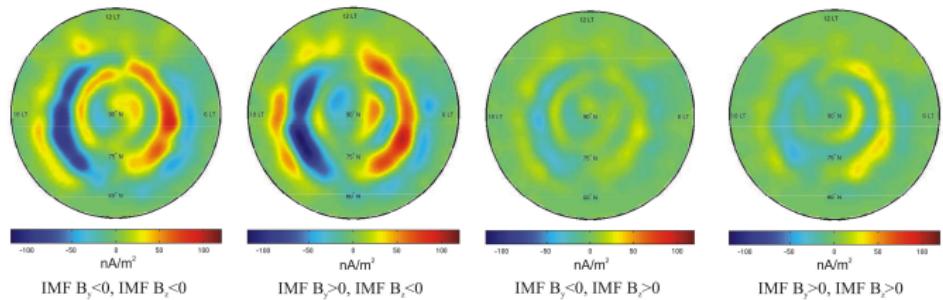
S-Winter



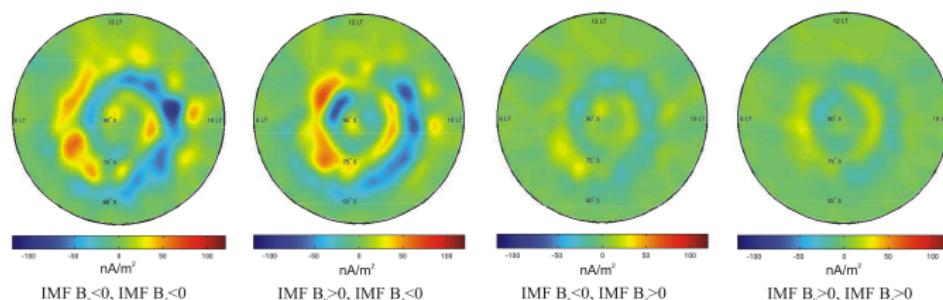
CHAMP Data: Jan. 2001 - Dec. 2007

Subtraction of CHAOS-3 model, Binning via Huber's M-estimation

N-Equinox



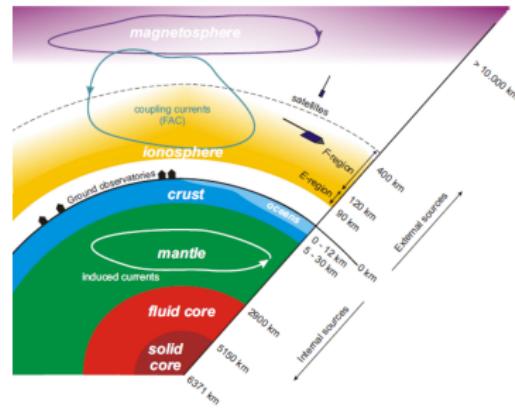
S-Equinox



SEPARATION OF SOURCES

Mie-Decomposition:

$$\begin{aligned} b &= p_b + q_b = \nabla \wedge LP_b + LQ_b, \\ j &= p_j + q_j = \nabla \wedge LP_j + LQ_j. \end{aligned}$$



Separation of the Poloidal Part:

$$\nabla \wedge p_b^{int}(x) = \begin{cases} q_j(x), & x \in \Omega_R^{int}, \\ 0, & x \in \Omega_R^{ext}, \end{cases}$$

$$\nabla \wedge p_b^{ext}(x) = \begin{cases} 0, & x \in \Omega_R^{int}, \\ q_j(x), & x \in \Omega_R^{ext}. \end{cases}$$

Potential Representation:

$$p_b^{int}(x) = \nabla U^{int}(x), \quad x \in \Omega_R^{ext}, \quad p_b^{ext}(x) = \nabla U^{ext}(x), \quad x \in \Omega_R^{int}.$$

Vector Spherical Harmonics:

$$\nabla_x \frac{1}{R} \left(\frac{R}{|x|} \right)^{n+1} Y_{n,k} \left(\frac{x}{|x|} \right) = -\frac{1}{R^2} \left(\frac{R}{|x|} \right)^{n+2} \tilde{o}_\xi^{(1)} Y_{n,k}(\xi),$$

$$\nabla_x \frac{1}{R} \left(\frac{|x|}{R} \right)^n Y_{n,k} \left(\frac{x}{|x|} \right) = \frac{1}{R^2} \left(\frac{R}{|x|} \right)^{n-1} \tilde{o}_\xi^{(2)} Y_{n,k}(\xi).$$

Sphärische Operatoren:

$$\tilde{o}^{(1)} = o^{(1)} \left(D + \frac{1}{2} \right) - o^{(2)}, \quad o_\xi^{(1)} = \xi, \quad D = \left(-\Delta^* + \frac{1}{4} \right)^{\frac{1}{2}},$$

$$\tilde{o}^{(2)} = o^{(1)} \left(D - \frac{1}{2} \right) + o^{(2)}, \quad o_\xi^{(2)} = \nabla_\xi^*,$$

$$\tilde{o}^{(3)} = o^{(3)}, \quad o_\xi^{(3)} = L_\xi^*.$$

Decomposition:

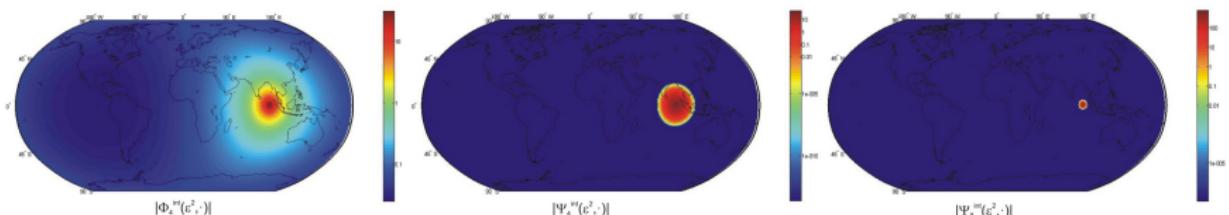
$$b = p_b^{int} + p_b^{ext} + q_b = \tilde{o}^{(1)} B^{int} + \tilde{o}^{(2)} B^{ext} + \tilde{o}^{(3)} B^q.$$

Multiscale Representation of the Internal Magnetic Field:

$$p_b^{int}(x) = \lim_{J \rightarrow \infty} \int_{\Omega} \Phi_J(\xi, \eta) b(r\eta) d\omega(\eta),$$

with

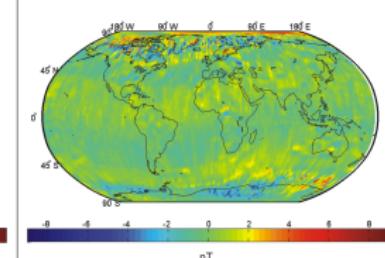
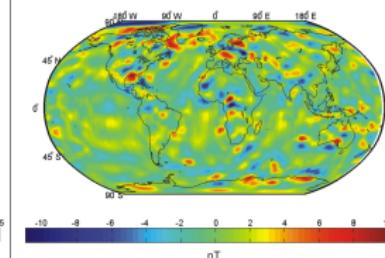
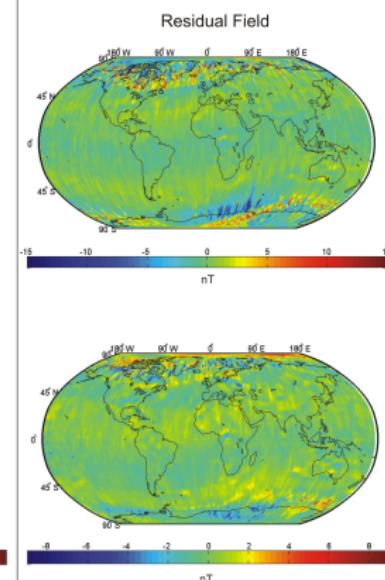
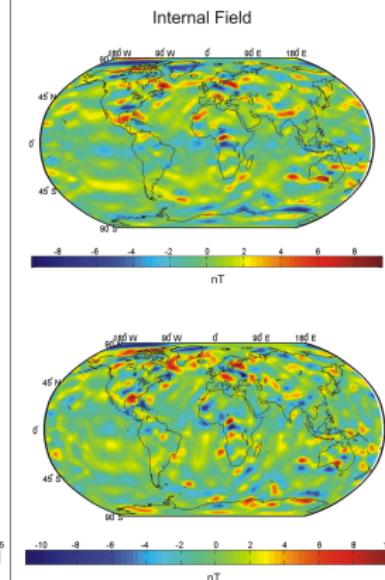
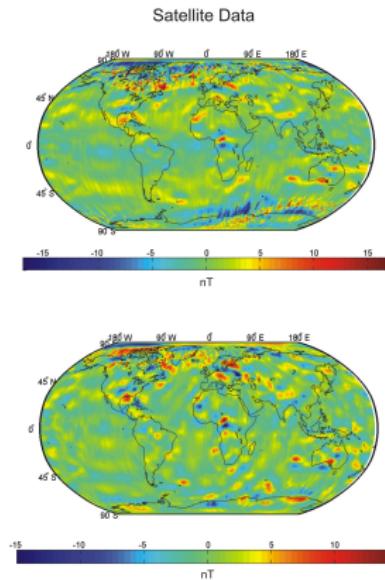
$$\begin{aligned} \Phi_J(\xi, \eta) &= \xi \otimes \eta \left(\frac{1}{2} \Delta_{\xi}^* G^J(\Delta^*; \xi \cdot \eta) + \frac{1}{8\pi} S^J(\xi \cdot \eta) \right) - \frac{1}{4\pi} \xi \otimes \nabla_{\eta}^* S^J(\xi \cdot \eta) \\ &+ \frac{1}{4} \nabla_{\xi}^* \otimes S_{\nabla^*}^J(\eta, \xi) - \frac{1}{4\pi} \nabla_{\xi}^* S^J(\xi, \eta) \otimes \eta - \frac{1}{2} \nabla_{\xi}^* \otimes \nabla_{\eta}^* G^J(\Delta^*; \xi \cdot \eta). \end{aligned}$$



1. Gerhards: PhD Thesis, Geomathematics Group, TU Kaiserslautern

CHAMP Data: June 2001 - Dec. 2001

Binning via Huber's M-estimation

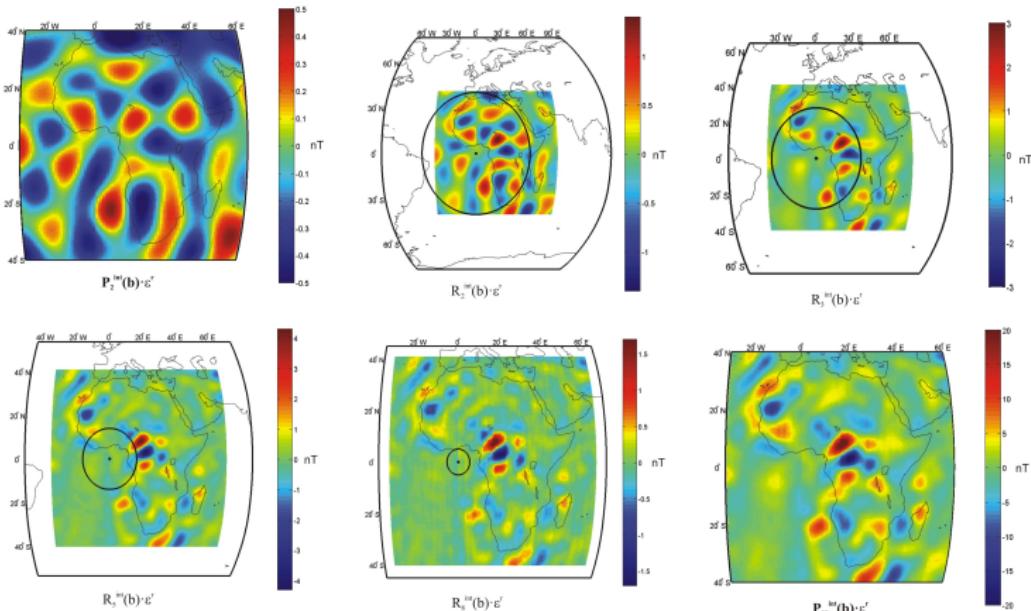


- Mayer, Maier: Geophys. J. Int. 167 (2006)

SEPARATION OF SOURCES

CHAMP Data: Jan. 2009 - Apr. 2010

Subtraction of CHAOS-4 model, Binning via Huber's M-estimation



1. Gerhards: PhD Thesis, Geomathematics Group, TU Kaiserslautern