On the origin of the 1/f spectrum in the solar wind



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Introduction

Aim: building the low-frequency (1/f) spectrum observed in the solar wind Method: couple an MHD turbulence model with a model of solar wind Results: the coupling of the upward and downward Alfvén waves builds a mixed weak & strong turbulence which generates a $k_{\perp}^{-5/3}$ spectrum together with a 1/f spectrum for the largest eddies which could be at the origin of the 1/f spectrum observed at larger distances

The low-frequency magnetic spectrum

The break between the f⁻¹ and f^{-5/3} range is thought to mark the boundary between the large scales dominated by expansion and the smaller scales dominated by nonlinear coupling (Tu et al 1984)



MHD with mean field B° Slow transport downward from air- Fast upward transport from solar ocean interface surface AIR OCEAN CORDNA

Homogeneous turbulence

Physical/Numerical model

1. Reduced MHD (b/B° <<1) •No gradients // B° (x) except for linear propagation (B°.V) •No parallel comp. (b//=u//=0), •Incompressible limit (∇_{\perp} .b₁= ∇_{\perp} .u₁=0) => Quasi 2D, but still b,u depend on x Two regimes (initial conditions or forcing): • WEAK if $B^{\circ}. \nabla >> b_{\perp}. \nabla_{\perp} => no // cascade, E(k_{\perp}) \propto k_{\perp}^{-2}$ • STRONG if $B^{\circ}.\nabla \approx b_{\perp}.\nabla_{\perp} => E(k_{\perp}) \propto k_{\perp}^{-5/3}$, $E(k_{\perp}) \propto k_{\perp}^{-2}$

2. Shell Reduced MHD

 solves for û(x,k⊥) (1/2 FFT) (one wavenumber per "shell" k⊥=2ⁿ k^o) • allows reaching $\text{Re} \approx 10^6$

Strong and weak forcing: 3D spectrum Strong forcing Weak forcing 105



Dashed : $\chi = 1$ isocontour; dotted: theoretical $\chi = 1$ isocontour

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Dotted: reduced spectra obtained from 3D spectra after suppressing excitation in region $\chi < 1/2 \Rightarrow$ suppresses k_{\perp}^{-2} weak scaling (right)

Strong and weak forcing: k//-m tails



•Energy in $\chi \leq 1$ region : $\mathbf{E}_{3}(\mathbf{k}_{\perp},\mathbf{k}_{\ell}) \propto \mathbf{k}_{\ell} \cdot \mathbf{m}^{-n}$ due to $1/f^{n}$ spectrum of large \perp eddies • Δ marks boundary of "strong" regime $\Delta f = 1/t_{cor}(k_{\perp}) = 1/t_{NL} = kz^{\pm}$

Time scale diagnostics $\chi = kz^2/k_{1/2}B^{\circ}$



Solar Wind turbulence

Physical/Numerical model (solar wind case)

1. Numerical model

·Choose a (numerical) 1D (radial) stationary wind solution (solving energy equations

etc ...) U. B. n. V. •Use Shell-Reduced MHD + variable phase speed U±Va including expansion weakening on perp. gradients

2. New physics

•Stratification ; z + >> z- (imbalanced turbulence)

Snapshots: from surface to Alfvén radius ($\approx 19R_s$)



Solar Wind at 1RA: reduced perpendicular spectra



•Both upward and downward spectra $\propto k_{\perp}^{-5/3}$ •Large imbalance (weak reflection): $E^- \approx E^+/10 \Longrightarrow t_{NL^+} \propto 1/z_- > t_{NL^-} \propto 1/z_+$

Solar Wind at 1RA: the 1/f sound of large eddies $E^{\pm}(f, k_{\perp}) \ (k^{\circ} \le k_{\perp} \le 2^{20} \ k^{\circ})$



Origin of the frequency spectrum

•R << RAlfvén: the absolute frame coincides with the plasma frame, hence the frequency spectrum (as well as the autocorrelation of the signal) reflects - either the nonlinear clock of eddies (strong regime) - or a shorter time (weak turbulence) - or shorter times scales of smaller eddies

 R≈R_{Alfvén} : the frequency spectrum reflects partially the internal clock of the eddies AND the parallel spatial structure of the signal

Diagnostics of weak and strong turbulence

Figure below shows that •for the major upward species z+ the cascade is weak (and becomes weaker as distance increases) •for z- the cascade is strong



Conclusion

Time diagnostics shows that in the subalfvénic region the upward wave suffers weak coupling, while the downard wave suffers strong coupling. However, probably due to the recycling of downward into upward waves at the transition region, both wave species show a perpendicular "strong" spectrum $\propto k_{\perp}^{-5/3}$. At the Alfvén point the integrated spectrum of the dominant component (z+ or b) scales as 1/f at low f (figure below) How evolves the 1/f spectrum at increasing R ? NB The S-RMHD model cannot explore this, since has to be radial



As B° rotates (Parker spiral) the kr1 law may become the observed f⁻¹

We are presently checking this hypothesis using the Expanding Box Model (3D MHD with comobile coordinates) (Grappin Velli Mangeney 1993)