

# Free energy sources in current sheets formed in collisionless plasma turbulence

Neeraj Jain<sup>1</sup>, Jörg Büchner<sup>1</sup>, Amir Chatraee<sup>2</sup>

<sup>1</sup>*Center of Astronomy and Astrophysics, Technical University Berlin,  
Hardenbergstraße 36, Berlin, Germany.*

<sup>2</sup>*Freie University, Kaiserswerther Str. 16-18, 14195 Berlin, Germany*

## **Collaborators:**

Uwe Motschmann, Horia Comișel (TU-Braunschweig)

Zoltán Vörös, Yasuhito Narita (Space Research Institute, Graz)

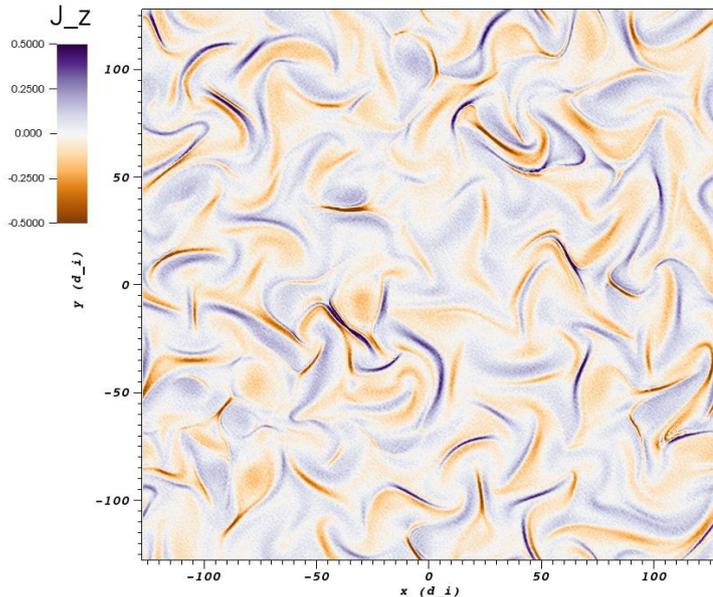
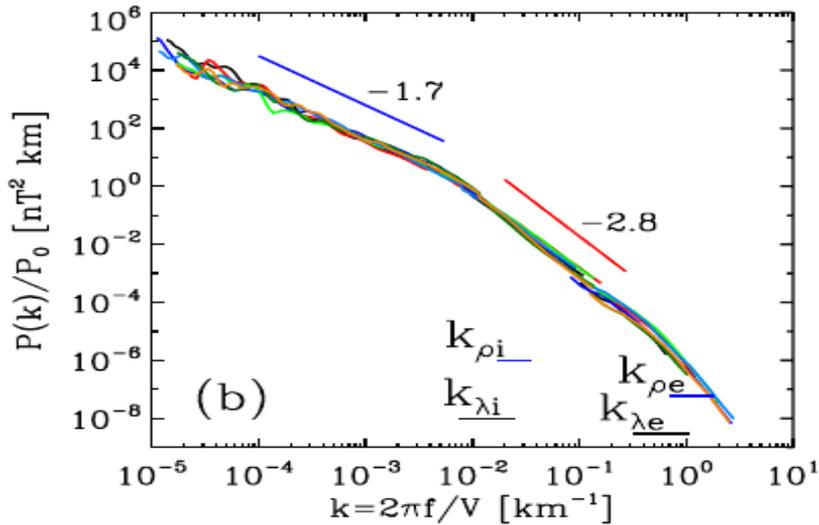
# Collisionless Plasma Turbulence

- Space and astrophysical plasmas are usually collisionless and in turbulent state.
  - Earth's magnetosphere, Solar wind, ISM, star atmosphere etc.
- Turbulence acts to dissipate energy into heat by transferring it to dissipation scales.
  - Slower-than-adiabatic-cooling of the solar wind is accounted for by its turbulent heating.
- In space and astrophysical plasmas with rare collisions,
  - Collisional dissip. scale  $\ll$  kinetic scales (gyro-radius, inertial length)
- Energy cascades down all the way to kinetic scales before hitting collisional dissipation scale.

What happens at kinetic scales?

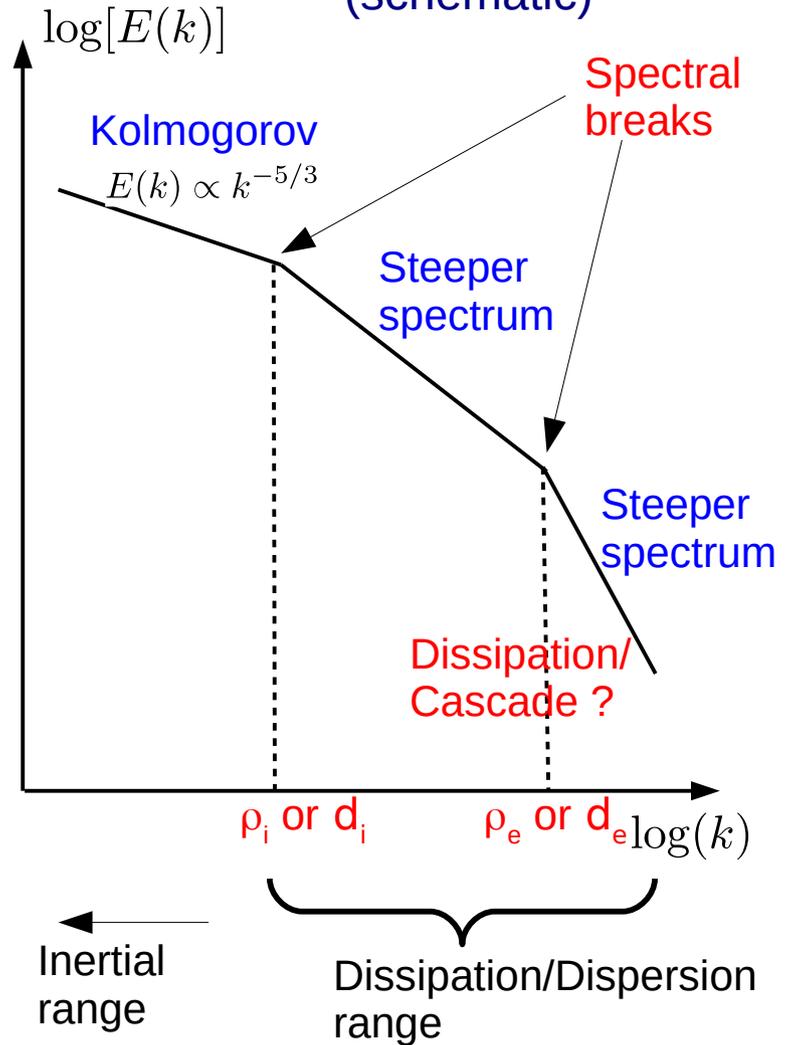
# Spectral breaks forms in energy spectrum at kinetic scales.

Solar wind obs. at 1 AU  
[Alexandrova et al. (2009)]



CS forms at kinetic scales

Kinetic scale energy spectrum (schematic)

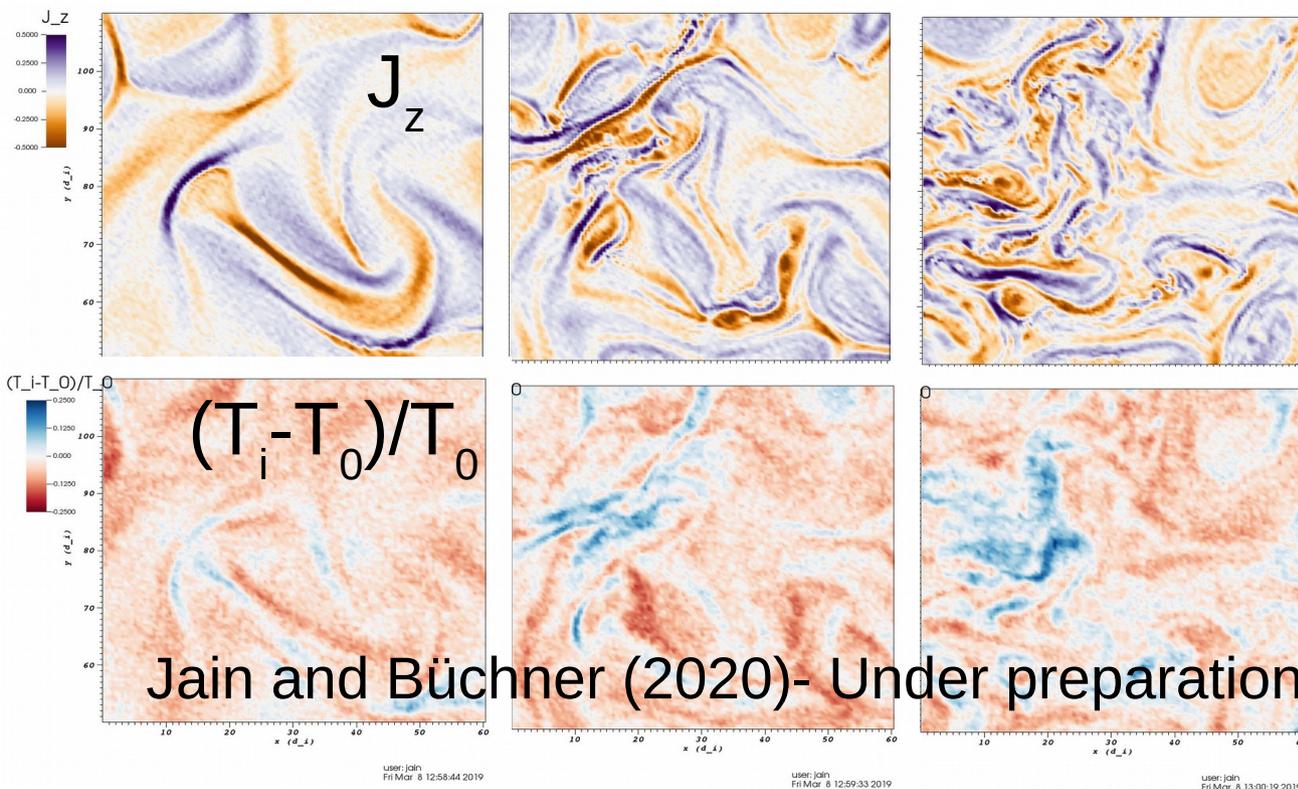


What about dissipation?

# Dissipation takes place in and around kinetic scale current sheets formed in collisionless plasma turbulence.

- Current sheets form self-consistently at kinetic scales. [Franci et al. (2015), Servidio et al. (2014)]
- Simulations and space observations show that the dissipation is concentrated in and around these current sheets [Matthaeus et al. (2015)].

Time evolution



What are the physical processes responsible for dissipation?

# Kinetic plasma processes responsible for dissipation in current sheets are not well understood.

Several processes were considered.

- Magnetic reconnection
  - generates parallel electric field which can accelerate particles
  - In plasma turbulence basically a time-dependent tearing-like instability process
- Fermi acceleration:
  - energization by multiple bounces in contracting magnetic islands
  - magnetic islands are generated by magnetic reconnection
- Stochastic acceleration:
  - energization by chaotic orbits of ions for large enough fluctuation amplitudes
  - plasma instabilities can influence fluctuation amplitudes
- Plasma instability in CS:
  - can produce anomalous dissipation by generating their own turbulence
  - can generate plasma waves which participate in dissipation and/or cascade

Plasma instabilities in CS directly or indirectly influence the dissipation/cascade and resulting spectral behavior of plasma turbulence.

# Free energy sources in CS

- Plasma instabilities are driven by free energy sources.
  - Spatial gradients of fluid variables
  - non-Maxwellian features of the distribution function
- Free energy sources in current sheets formed in collisionless plasma turbulence need to be characterized to understand the role of plasma instabilities in collisionless dissipation.
- In a current sheet, current density,  $J = e n (u_i - u_e)$ , is confined in small thickness.
  - Gradients in at least one of  $n$ ,  $u_i$  and  $u_e$  are present in current sheets.  
(Focus of this talk)
  - In addition non-Maxwellian features, like, temperature anisotropy can also be present. (Work under progress)

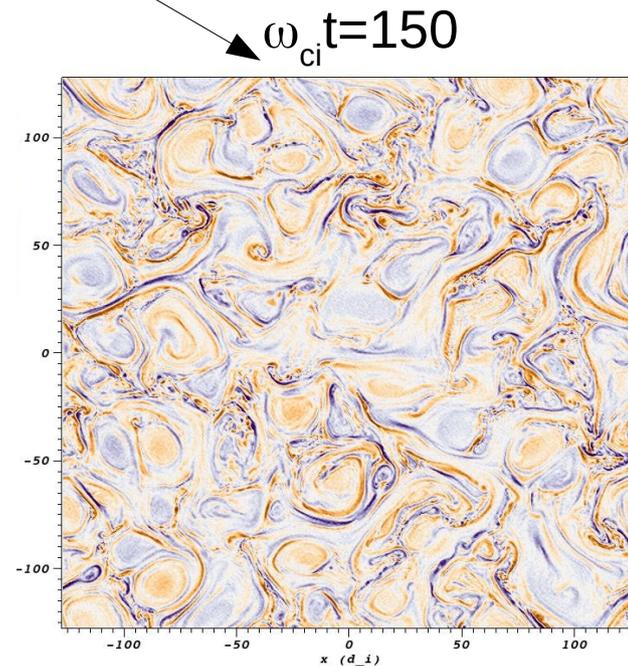
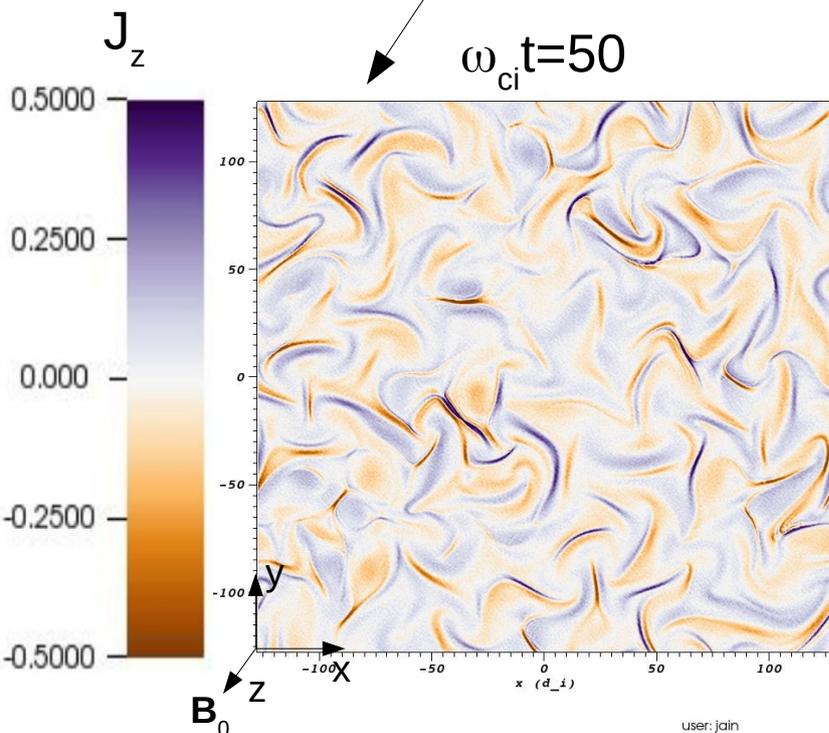
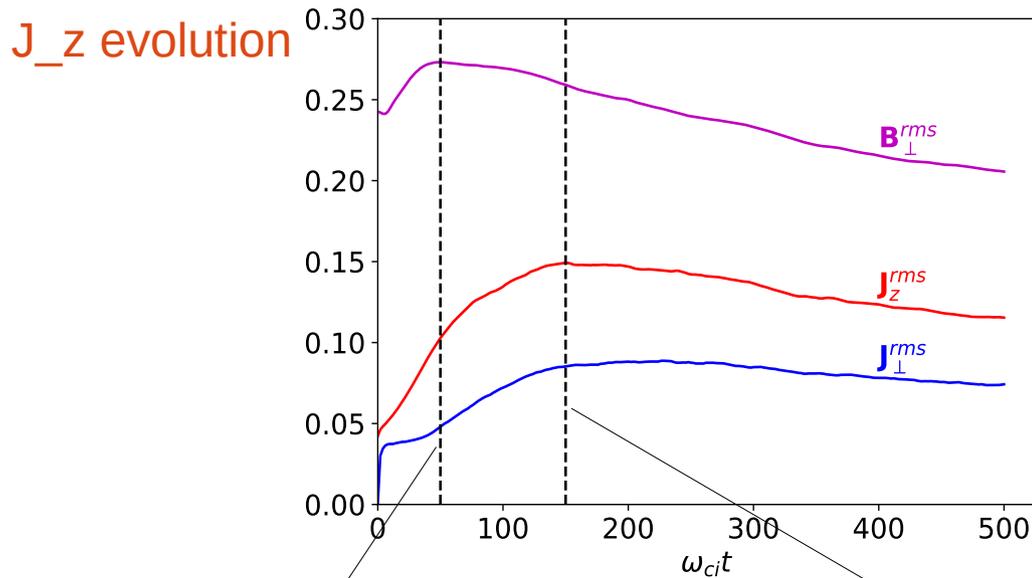
# We carry out hybrid simulations to address physics near ion-kinetic scales.

- 2-D (x-y plane) hybrid simulations by A.I.K.E.F. code [Mueller et al, 2011]
  - Ions as particles and electrons as inertialess fluid
- Initial conditions:
  - A uniform plasma ( $\beta_i = \beta_e = 0.5$ )
  - – random phased equal amplitude Alfvénic fluctuations in the range
$$k_{\min} d_i < k d_i < 0.2 < k_{\text{break}} d_i \approx 1$$
    - An external magnetic field ( $B_0$ ) perpendicular to the 2-D simulation plane.
- Periodic boundaries.
- Simulation parameters:
  - box size:  $256 d_i \times 256 d_i$
  - Grid points: up to  $2048 \times 2048$
  - Particles per cell: 500
- Normalization:
  - length by  $d_i$
  - time by  $\omega_{ci}^{-1}$
  - Magnetic field by  $B_0$

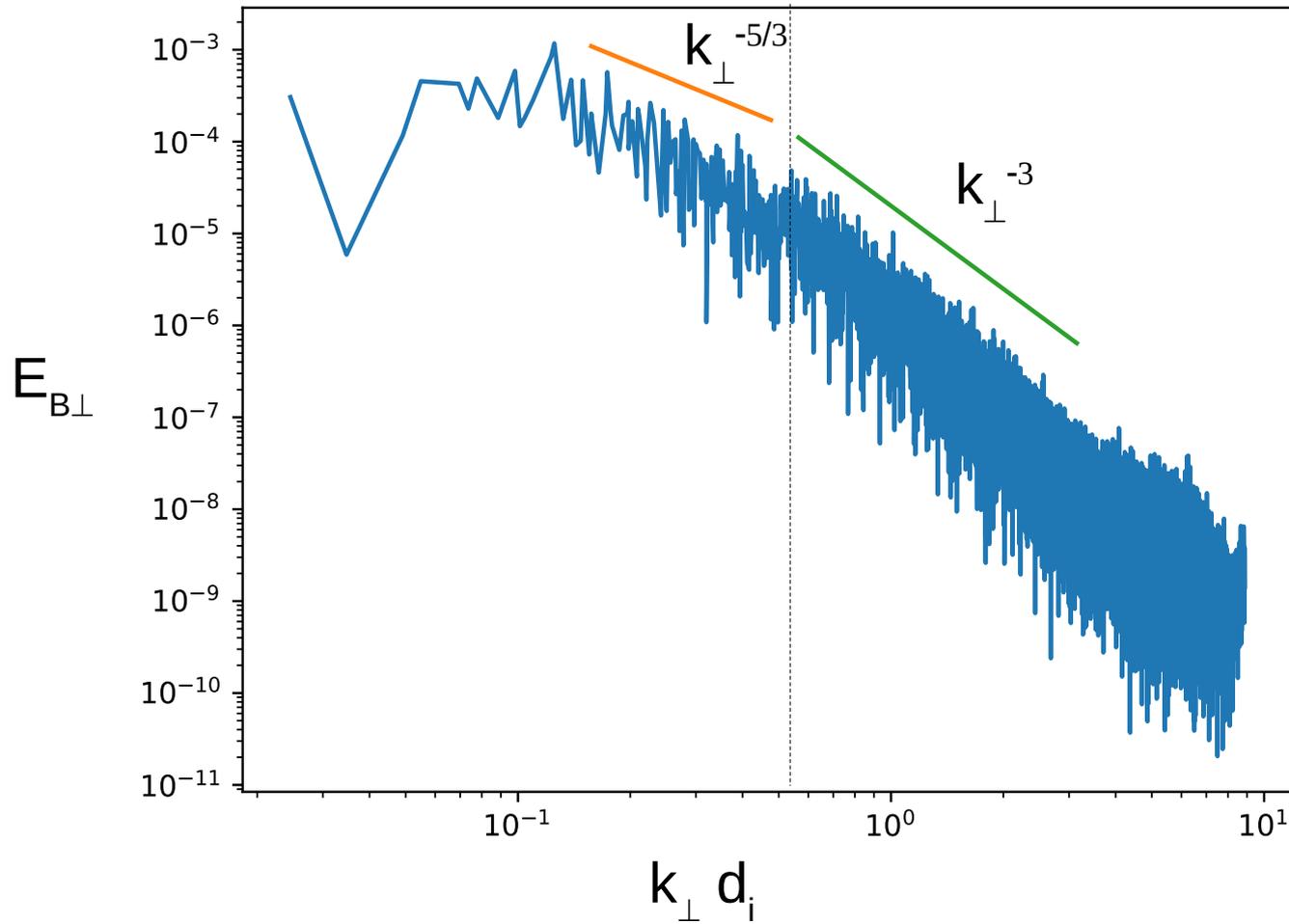
# Initial random fluctuations evolve into current sheets (CS) which later disintegrate.

Jain and Büchner (2020)  
*(under preparation)*

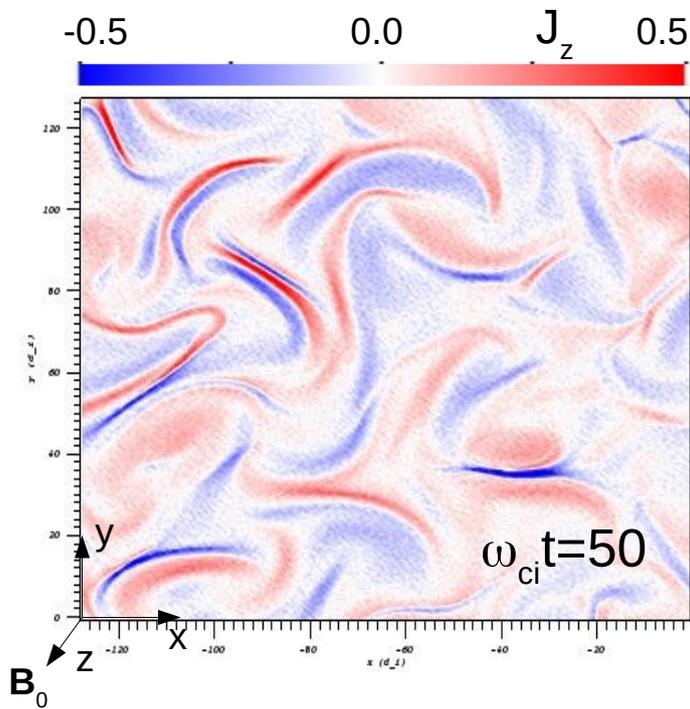
In order to identify free energy sources in current sheets, we analyze them at  $\omega_{ci}t=50$  before they disintegrate.



Spectrum of simulated turbulence develops an ion-scale break.



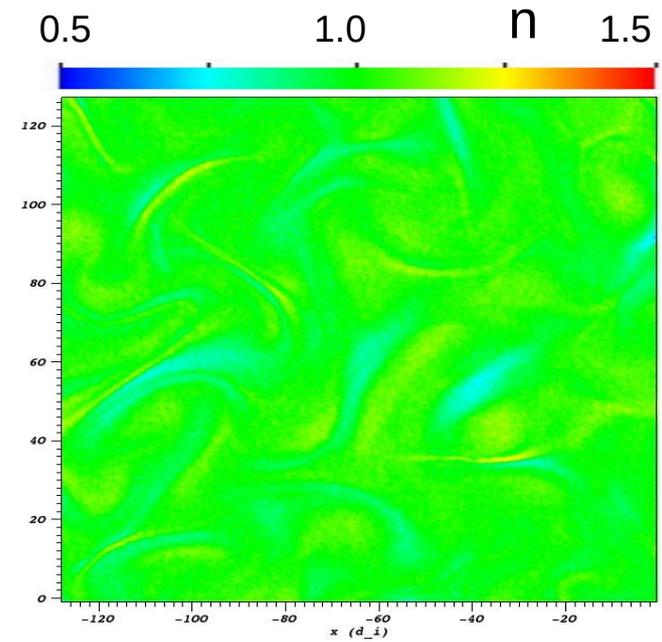
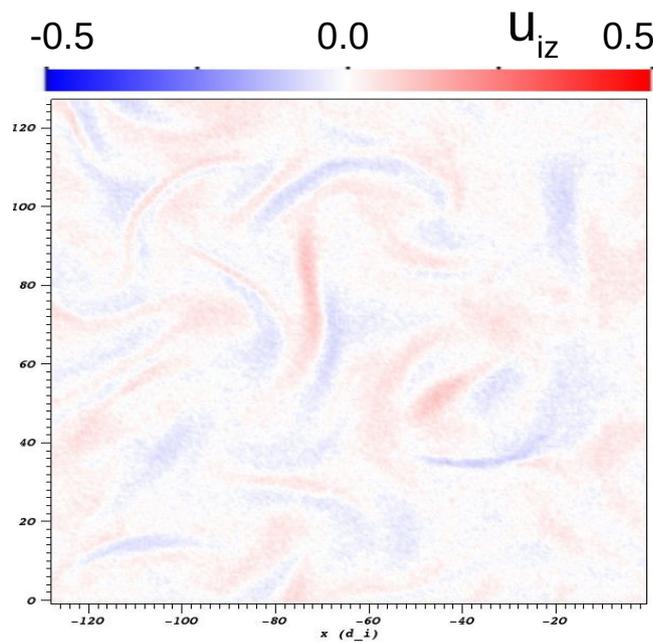
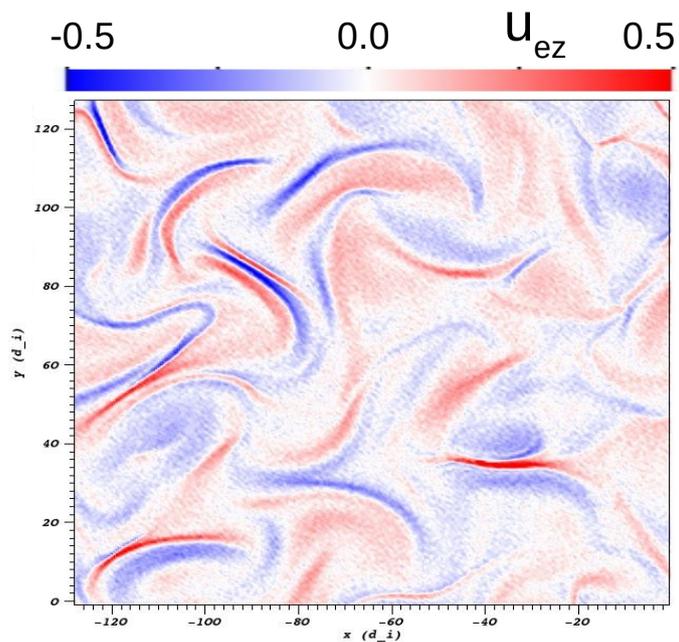
# Current sheets are primarily due to electron shear flow.



$$J_z = n (u_{iz} - u_{ez})$$

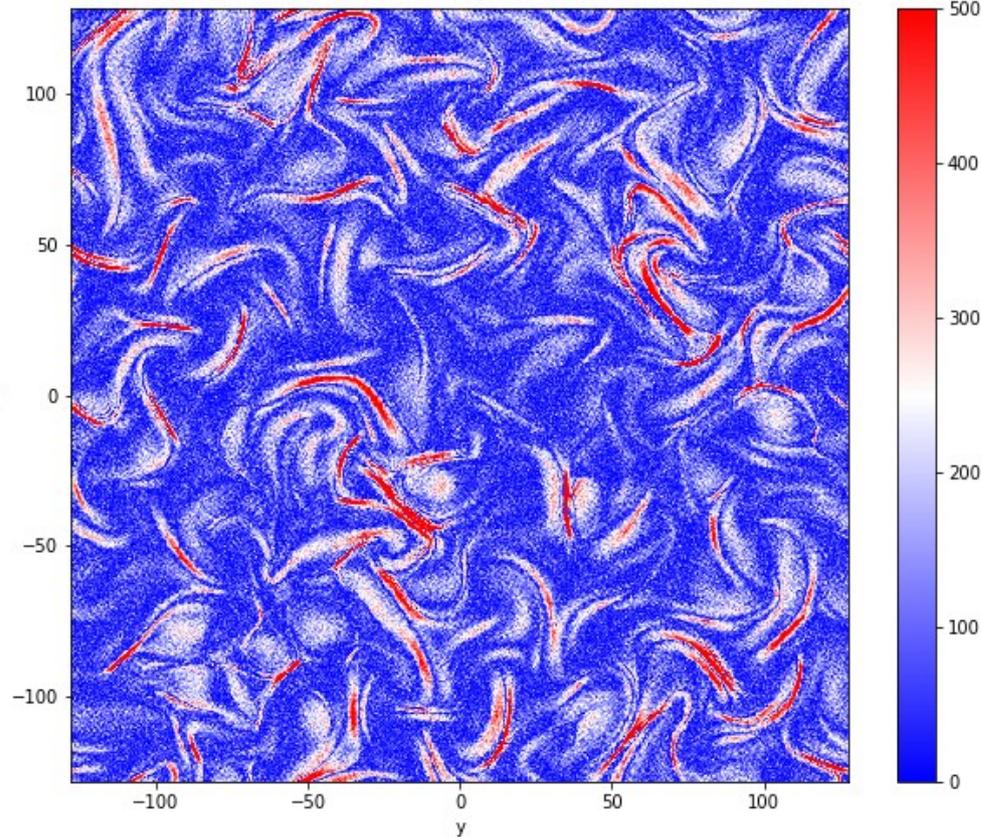
Contributions of  $u_{ez}$ ,  $u_{iz}$  and  $n$ ?

Results are shown on a quarter of the simulation domain.



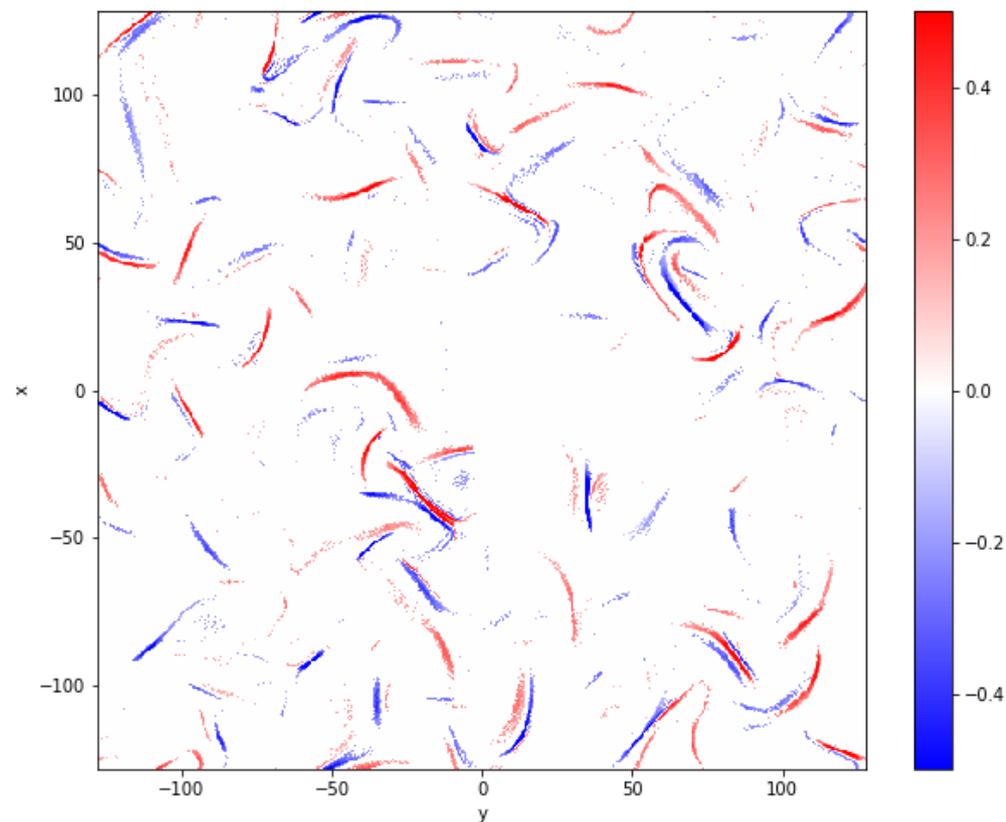
$|u_{ez}|/u_{iz,rms}$  has large value  $\sim 300$  in current sheets.

$|u_{ez}|/u_{iz,rms}$



Masked  $J_z$ :

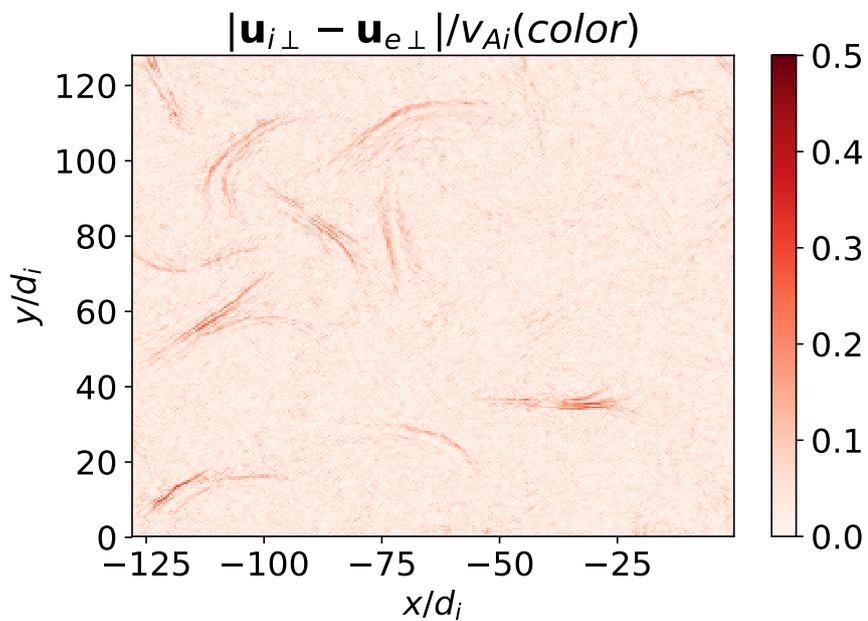
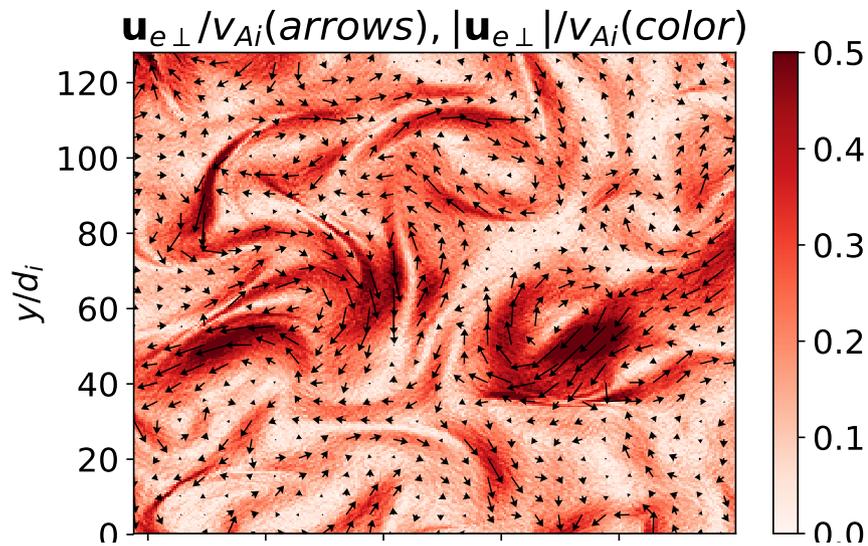
$J_z [ |u_{ez}|/u_{iz,rms} < 300 ] = 0$



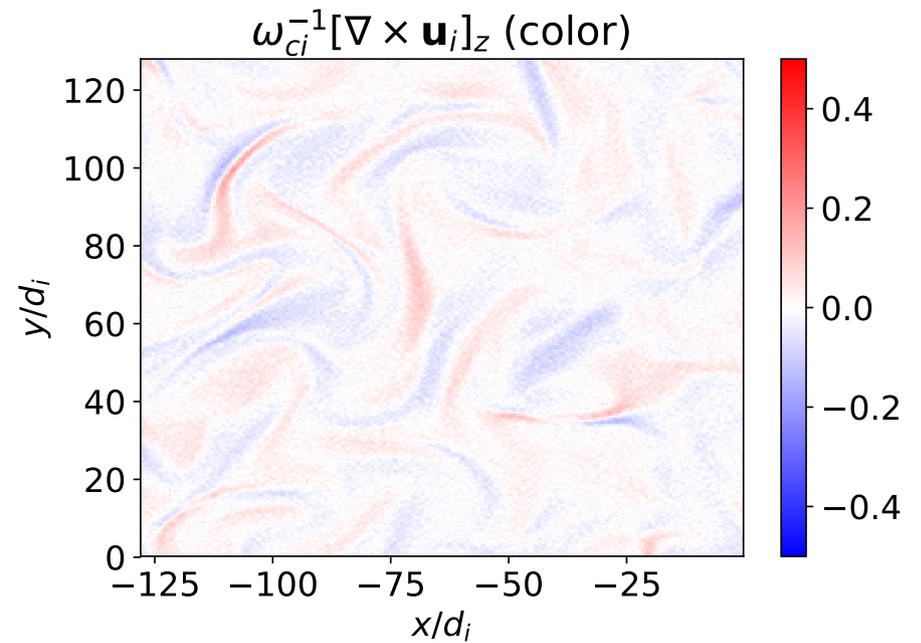
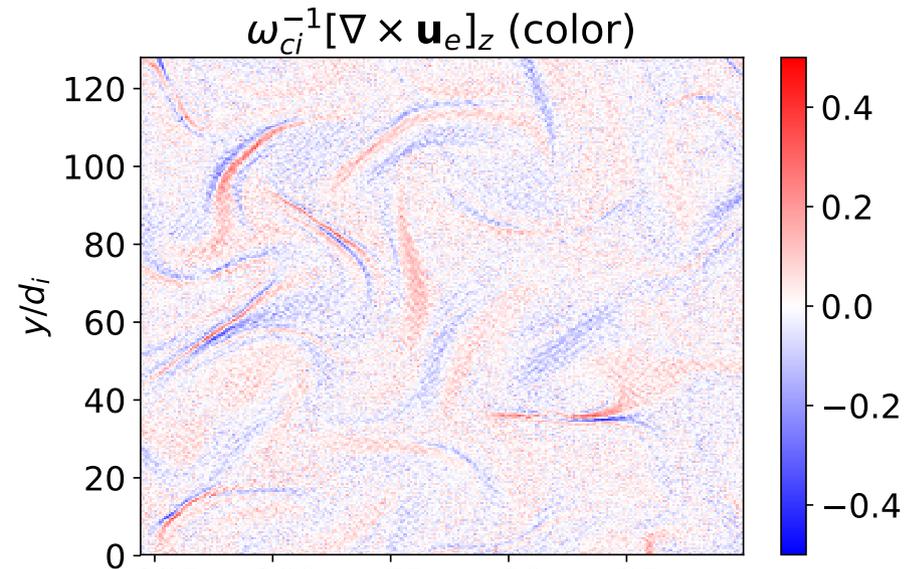
We have used large value of  $|u_{ez}|/u_{iz,rms}$  as condition in numerical algorithm developed for the identification and characterization of CS in turbulence.

# Perpendicular flows

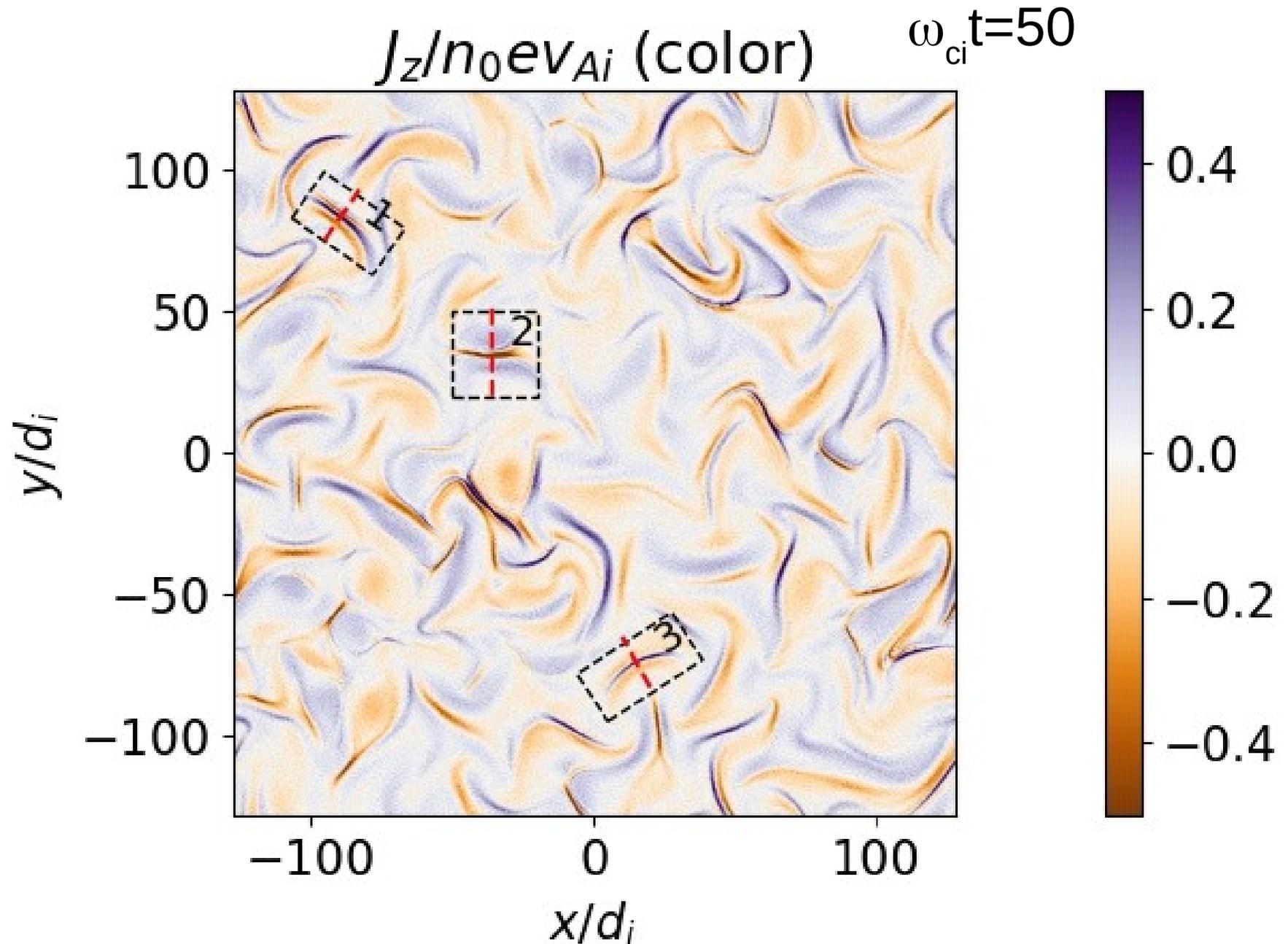
$\mathbf{u}_{i\perp} \approx \mathbf{u}_{e\perp}$  except at few locations



Perpendicular flows have shear



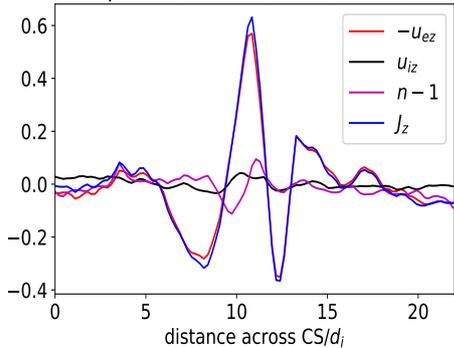
We select manually from the full simulation domain three current sheets for a closer inspection.



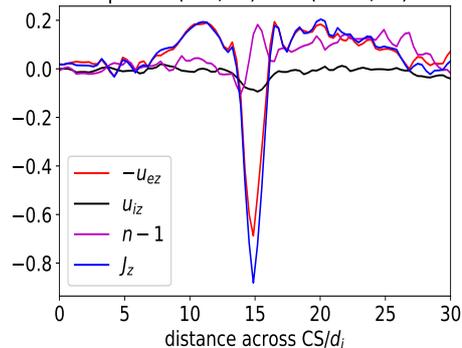
# Line-outs across current sheets

$J_z$ : Blue line,  
 electron quantities: red line,  
 ion quantities: black line

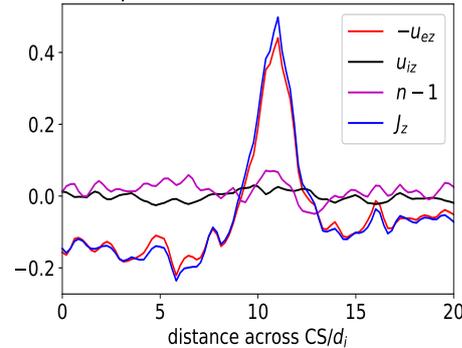
line-outs along the CS normal-1 joining points (-95.5,75) and (-83,93)



line-outs along the CS normal-2 joining points (-36,20) and (-36,1.52)



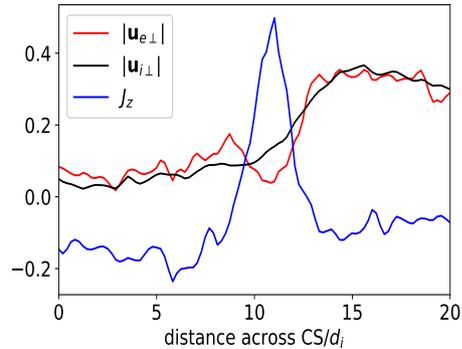
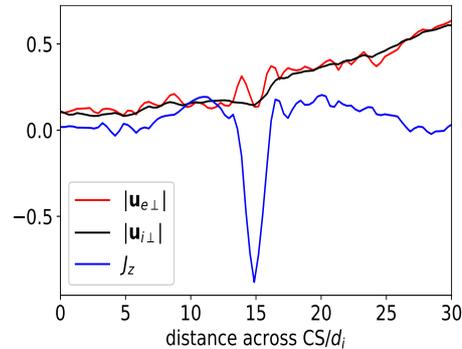
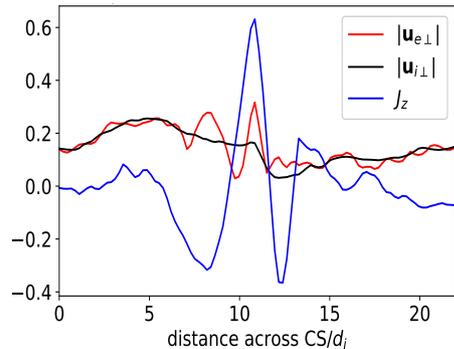
line-outs along the CS normal-3 joining points (20,-83) and (10,-65)



$J_z$  dominated by  $u_{ez}$

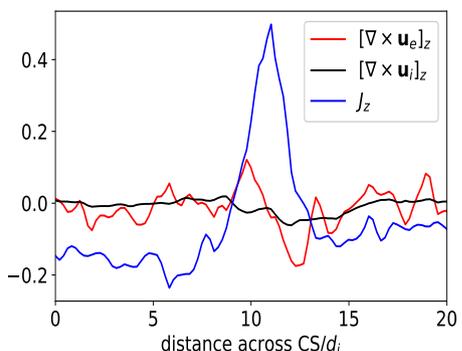
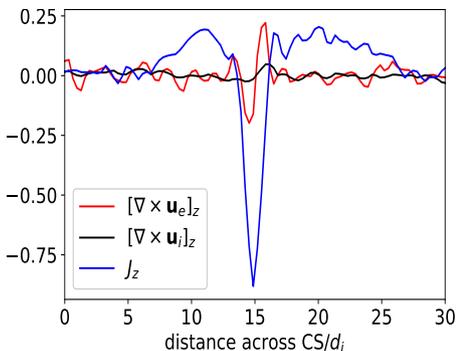
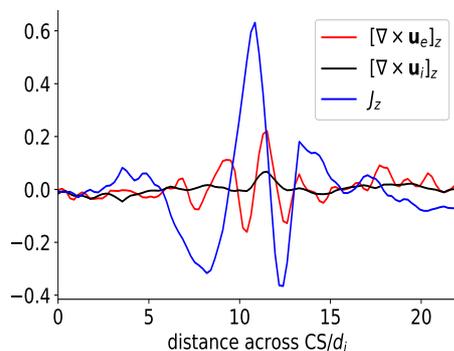
$\Delta n/n \sim 0.1$

Return current



$|\mathbf{u}_{i\perp}| \approx |\mathbf{u}_{e\perp}|$  outside CS

$|\mathbf{u}_{e\perp}|$  differs from  $|\mathbf{u}_{i\perp}|$   
 and varies faster  
 inside CS



electron vorticity larger  
 than ion vorticity,

Changes sign at center,  
 peaks at edges of CS

## Theoretical estimates support simulation results.

At the ion scales, magnetic field is pushed around by turbulent electron flows generating a parallel inductive electric field  $E_z$  which accelerate ions.

At the scales of current sheet  $\sim d_i = 2^{1/2} \rho_i$ , ions are approx. unmagnetized.

$$\frac{\partial u_{iz}}{\partial t} = \frac{e}{m_i} E_z \quad \text{Neglecting convective derivative}$$

Electron velocity adjusts to satisfy Ampere's law time derivative of which gives,

$$\frac{\partial u_{ez}}{\partial t} = \frac{e}{m_i} (E_z - d_i^2 \nabla^2 E_z) \quad \text{Not electron mom. eq.}$$

## Theoretical estimates support simulation results.

The two eq. gives,

$$\frac{|u_{ez}|}{|u_{iz}|} \sim \left| 1 - \frac{d_i^2}{L^2} \right|$$

For CS thickness  $L=0.5 d_i$ ,

$$|u_{ez}|/|u_{iz}| \sim 3.$$

Consistent with  
simulations

For  $L \ll d_i$ ,

$$|u_{ez}|/|u_{iz}| \sim d_i^2/L^2 \gg 1$$

Thinner the CS, more dominating  
is parallel electron flow

In hybrid models,  $\mathbf{u}_{e\perp}$  is always ExB drift. When strongly magnetized  $\mathbf{u}_{i\perp}$  is also ExB drift. In current sheets,  $\mathbf{u}_{i\perp}$  differ from  $\mathbf{u}_{e\perp}$  as ions are unmagnetized.

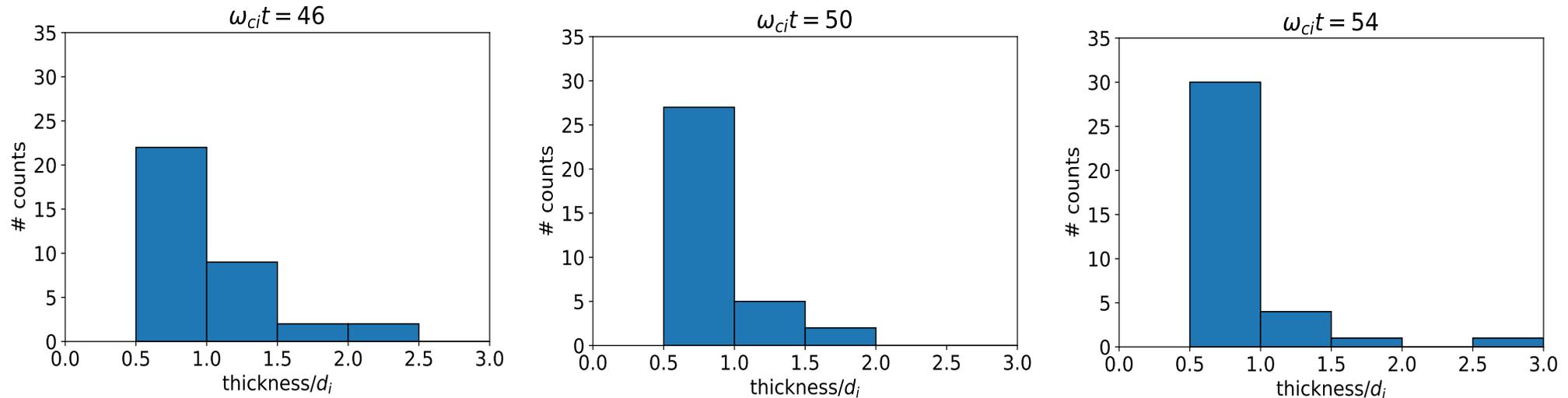
# Characterization of current sheets

- Gradient driven instabilities depend on scale lengths.
- Need of statistical characterization of CS in terms of thickness, length, aspect ratio
- Developed a python code implementing the Zhdankin et al (2013) algorithm of CS detection and characterization.
- CS detection depends on algorithm parameters.
- Varying the parameters give results statistically independent of them.

# Distribution of CS half-thickness

Chatraee, Jain, and Büchner (2020)  
*(under preparation)*

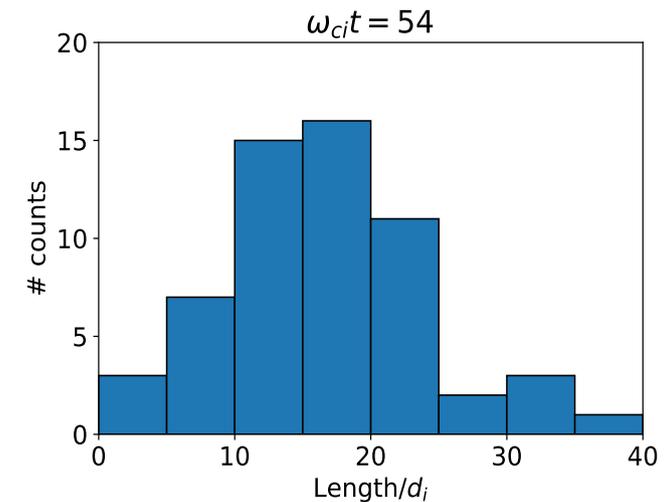
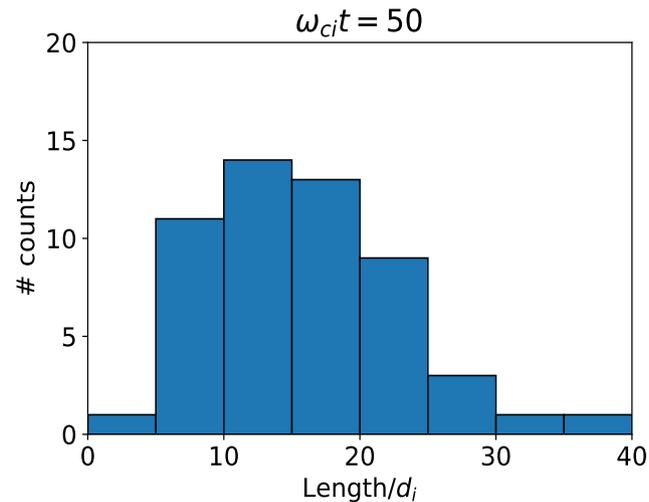
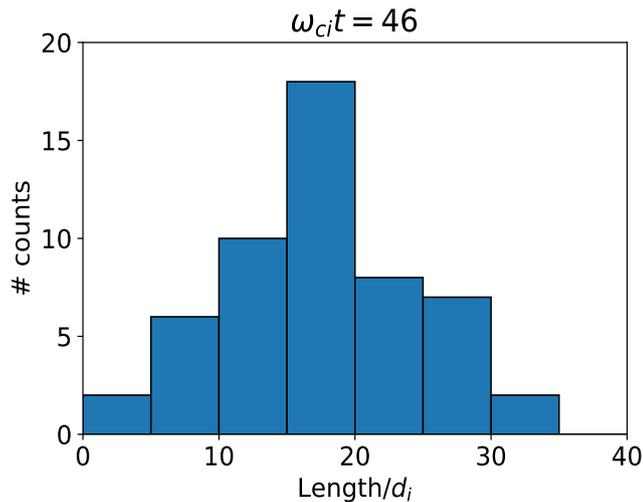
- Current sheets thin down to grid scales.



- Increasing grid resolution leads to thinner current sheets.
- Hybrid plasma model w/o electron inertia is not sufficient to address the physics of current sheets.
- Next step: Hybrid plasma simulations with electron inertia using a code CHIEF recently developed in our group.

# Distribution of CS length

- Distribution of CS length peaks around  $15 d_i$ .



- Expected half-thicknesses  $L_{CS} \ll d_i$
- Aspect ratio: Length/half-thickness  $\gg 15$ .
  - Condition  $k L_{CS} < 1$  of gradient driven instabilities is easily satisfied.

# Summary

- 2-D hybrid simulations of kinetic plasma turbulence show formation of current sheets which thins down to grid scale.
- Current sheets are formed mainly due to electron shear flow.
- In a current sheet:
  - parallel electron flow velocity dominates ion flow velocity
  - parallel electron vorticity is larger than ion vorticity
  - gradients in plasma density are weak ( $\sim 10\%$  variation).
  - Theoretical estimates support simulation results.
- Instabilities driven by parallel and perpendicular electron shear flow can grow in current sheets of collisionless plasma turbulence.
- Hybrid simulations including electron inertia are required to address the physics of the current sheets and instabilities therein.