

# Log-log linearity of the asymptotic distribution – a valid indicator of multi-fractality?



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# Introduction

It has been claimed (Cheng et al. 1994, Cheng 1999, etc.), that the presence of anomalies of spatial fields (typical: geochemical) resulting from its multi-fractal nature, is indicated by

$$\text{area}(Z > z) \sim z^{-\beta}$$

for high threshold  $z$ , i.e. a power-law behaviour.

Call  $A(z) := \text{area}(Z>z)$  the excursion set (Adler 1981)

statement equivalent to  $\lim_{z \rightarrow \infty} \frac{\ln A_z}{\ln z} = -\beta$

Question: Does multifractality imply the power-law behaviour?



# Multifractal

- reminder: locally,  $|Y(x+h) - Y(x)| \sim h^{\beta(x)}$   
 $\beta$  – local singularity or Hölder exponent; (Hausdorff-)fractal dimension of set of  $x$  with same  $\beta(x)$ : fractal spectrum. (In environmental applications: box-dim.)
- Notation:

$V_i(s)$  = i-the cell of size  $s$ , covering domain  $D$ ; their areas  $A(s)$

$m_i(s)$  = mass in cell  $i$  of size  $s$  =  $A(s) EZ (Z \in V_i(s))$

$p_i(s)$  = probability measure =  $m_i(s)/\sum m_i(s) = E_i Z / \sum E_i Z$ .

$S(s) := \sum m_i(s)$

$\mu_i(q,s) := p_i(s)^q / \sum p_i(s)^q = m_i(s)^q / S_q(s)$ ,  $S_q(s) := \sum m_i(s)^q$ .

- According Chhabra & Jensen (1989),

$$\alpha(q) = \lim_{s \rightarrow 0} \frac{\sum_i \mu_i(q,s) \log p_i(s)}{\log s} \quad f(\alpha(q)) = \lim_{s \rightarrow 0} \frac{\sum_i \mu_i(q,s) \log \mu_i(q,s)}{\log s}$$

- $a(q,s) := \sum \mu_i(s) \ln p_i(s) = \dots = (1/S_q(s)) \sum m_i(s)^q - \ln S(s)$   
 $f(q,s) := \sum \mu_i(s) \ln \mu_i(s) = \dots = (q/S_q(s)) \sum m_i(s)^q - \ln S_q(s)$
- $\alpha(q)$  and  $f(\alpha(q))$  estimated from slopes of graphs  $a(q,s)$  and  $f(q,s)$  vs.  $\ln s$ .



# Generation of multifractals by multiplicative cascade

*Ansatz: multiplicative binomial cascade*

1-dim example (next slide):

- 1) Domain, length=1. Contains  $n_0$  events.
- 2) Divided into 2 equal halves. Left half:  $n_1$  events, right half:  $n_0-n_1$ .  
 $E[n_1]=n_0/2=E[n_0-n_1]$ ,  $n_1 \sim f(n_1) = \text{Bin}(n_0, 1/2)$ .  
→ Sample  $n_1$  from  $f$ .
- 3)...i)  $2^i$  bins in  $i$ -th cascade,  $n_{i+1} \sim \text{Bin}(n_i, 1/2)$ .

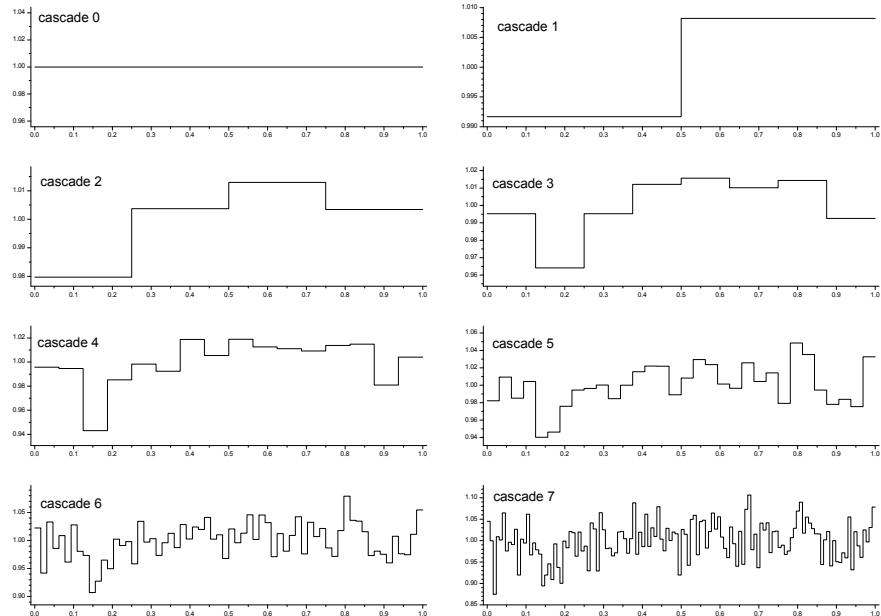
Technicalities:

- a) De Moivre (1733), CLT:  $\text{Bin}(n, p) \rightarrow N(np, \sqrt{np(1-p)})$
- b)  $n_{i+1} \equiv (n_i/2)(1-\gamma_i)$ ,  $n_i - n_{i+1} \equiv (n_i/2)(1+\gamma_i)$ :  **$\gamma_I = \text{splitting factor}$** .  
Construction ensures mass conservation.
- c) Transition: events  $n \rightarrow$  intensity  $z_i = n_i/2^{-i}$ ,  $2^{-i}$ =length of bin  $i$ -th cascade.
- d) Traditional random **p-model**:  $|\gamma_i| \equiv \gamma$ , random  $+\gamma$  or  $-\gamma$ .

History: Kolmogorov (1940s), De Wijs (1950s), Mandelbrot (1960-70s), Schertzer & Lovejoy (1980s), Agterberg & Chen (1990s-)

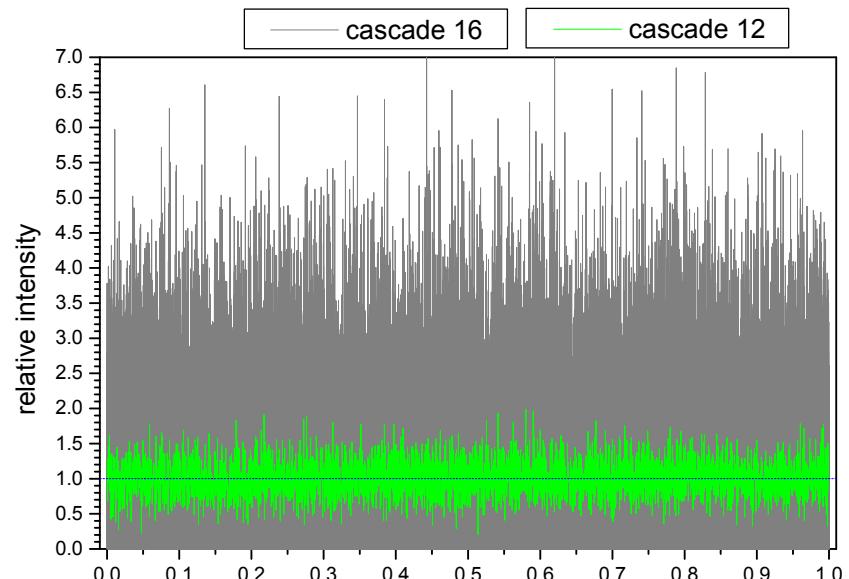


[from Bossew et al. 2008]



First 8 cascades

**Example:**  
1-dim, splitting factor  $\gamma=0.6$



Cascades 12 and 16 ( $2^{12}=4096$  and  $2^{16}=65536$  bins)

# Properties of cascades

De Wijs 2-dim cascade (p model) is asymptotically LN.

After k generations,  $\sigma = \ln(\text{GSD})$ :

$$\sigma^2 = \frac{k}{4} \left( \ln \left( \frac{1+\gamma}{1-\gamma} \right) \right)^2$$

Minimum singularity exponent:

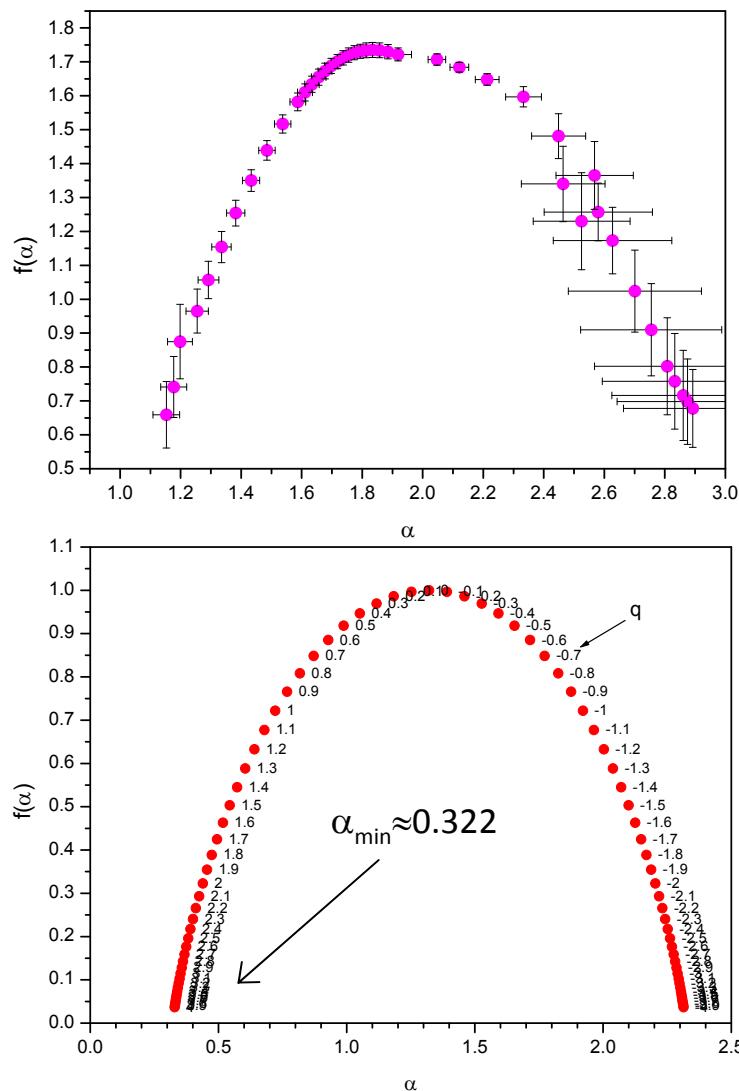
$$\alpha_{\min} = \log_2 \left( \frac{2}{1+\gamma} \right)^2 = 2 \left( 1 - \frac{\ln(1+\gamma)}{\ln 2} \right)$$

1-dim:  $\alpha_{\min} = 1 - \frac{\ln(1+\gamma)}{\ln 2}$  for  $\gamma=0.6$ :  $\alpha_{\min} \approx 0.322$

De Wijs 1951; Agterberg 2007a,b; and other.



# examples



Friedmann-Radon potential (Austria)  
(indoor Rn concentration normalized for  
anthropogenic factors)

1-dim multiplicative cascade, splitting  
factor  $\gamma=0.6$

Maximum of the curve:  $f(\alpha)=$  box dimension

# Cheng's derivation

Idea:

- Mass in cell  $i$  (size  $s$ ) scale as  $\mu_i(s) \sim s^{\alpha_i}$ ; therefore density or concentration,  $z_i = \mu_i/s^2$  in 2-dim case:  
 $z_i(s) \sim s^{\alpha_i-2}$ .
- Number of cells with  $\alpha_i = \alpha$ :  $N(s) \sim s^{-f(\alpha)}$ , therefore area(cells with  $\alpha$ )  $\sim s^{2-f(\alpha)}$ .
- Calculate area(cells from  $\alpha_{\min}$  to  $\alpha_{\min} + \delta$ ),  $\delta$ =small, i.e. the area of high anomalies (hot spots), defined by  $\alpha$  near  $\alpha_{\min}$ :

$$\text{area}(\alpha_{\min} \leq \alpha \leq \alpha_{\min} + \delta) \sim \int_{\alpha_{\min}}^{\alpha_{\min} + \delta} s^{2-f(\alpha)} d\alpha$$

- This can be approximated according to Cheng (1994); annex A, leading to

$$\log \text{area(cells from } \alpha_{\min} \text{ to } \alpha_{\min} + \delta) \approx \text{const.} - \beta \log z ,$$

i.e. a power-law behaviour (const. and  $\beta$  defined appropriately). However, the derivation (eq. (13) in the paper) is difficult to follow.

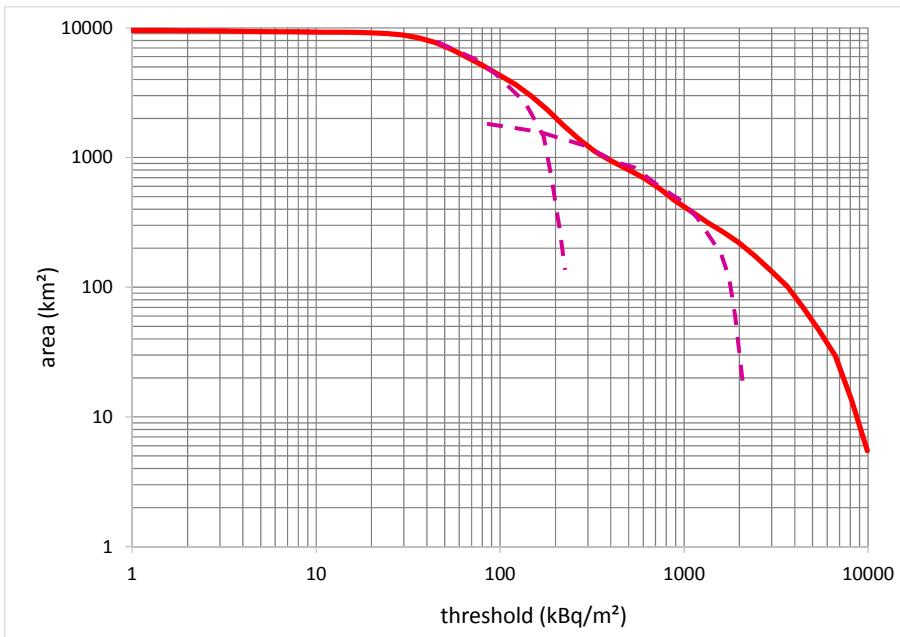
# LN distribution

- CDF of LN distribution  $\sim \text{area}(Z < z)$ :  $CDF = \frac{1}{2} \operatorname{erfc}\left[-\frac{\ln z - \mu}{\sigma\sqrt{2}}\right]$ ,
- $p_z := \text{prob}(Z > z) = 1 - \text{CDF} = \frac{1}{2} \operatorname{erfc}\left(\frac{\mu - \ln z}{\sigma\sqrt{2}}\right)$
- Asymptotic expansion:  $\operatorname{erfc}(x) = \frac{e^{-x^2}}{x\sqrt{\pi}} \sum_{k=1}^{\infty} (-1)^k \frac{(2k-1)!!}{(2x^2)^k}$  Abramowitz & Stegun 7.1.23
- with  $\operatorname{erfc}(x) - \frac{e^{-x^2}}{x\sqrt{\pi}} \sum_{k=0}^{n-1} (\dots) = O\left(x^{1-2n} e^{-x^2}\right)$
- $n=1$ :  $\operatorname{erfc}(x) \sim \frac{e^{-x^2}}{x\sqrt{\pi}} + O\left(\frac{e^{-x^2}}{x}\right)$
- $\ln$ :  $\ln \frac{e^{-x^2}}{x\sqrt{\pi}} = -x^2 - \ln x - \text{const}$

which is dominated by  $x^2$  for high  $x$ . Therefore, the log survival function is asymptotically shaped quadratic. – not linear!



# Mixture of populations; Example: Fukushima



"log-survival plot":

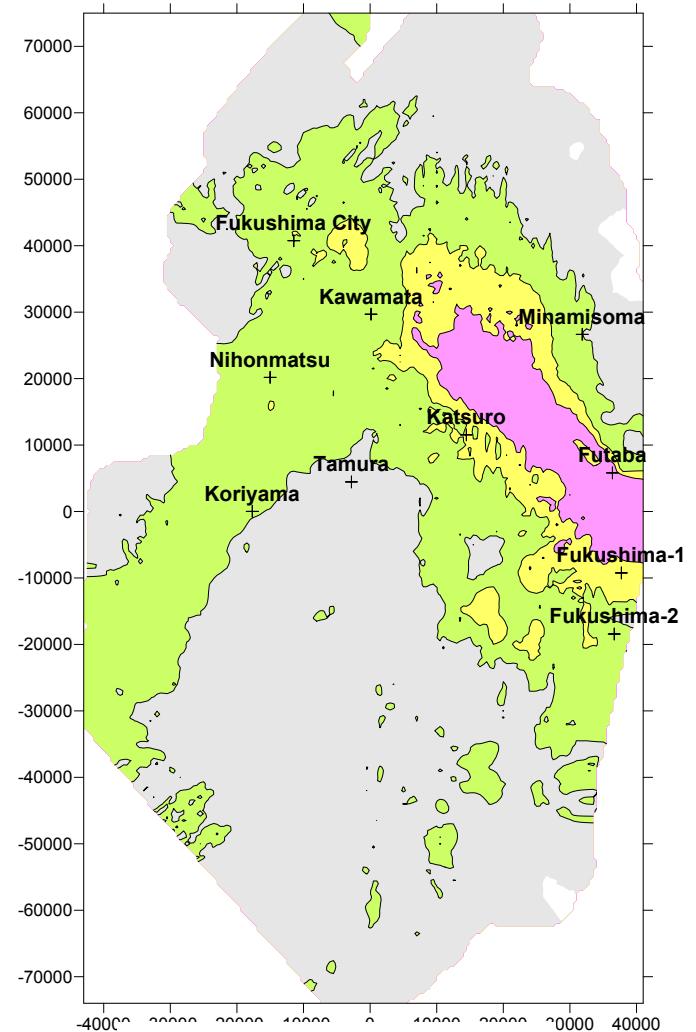
$\log(\text{area in which } Z>z)$  against  $\log(z)$   
e.g. area ( $^{137}\text{Cs} > 1 \text{ MBq/m}^2$ )  $\approx 400 \text{ km}^2$

Lognormal: would be falling parabola; here mixture of several (3-4) parabolas;

"Breakpoints" → separation of populations



[Bossew & Debayle, NARE-9, 2014]



breakpoints:  
green  $> 100 \text{ kBq/m}^2$ ,  
yellow  $> 330$ , pink  $> 900$

# Number and size of maxima

- Number  $M_z$  of maxima above  $Z=z$  for Gaussian fields, Adler 1981 (p.133, Theorem 6.3.1), 2-dim case:

$$EM_z = \frac{\text{area}(D)(\det \Lambda)^{1/2} z}{(2\pi)^{3/2} \sigma^3} \exp\left(-\frac{z^2}{2\sigma^2}\right) (1 + O(1/z))$$

$\Lambda$ =matrix of 2<sup>nd</sup> order spectral moments  $E(\nabla_i Z \nabla_j Z)$   
( $\approx$  inv FT of covariance, e.g. Hristopulos 2020, sec. 5.3.2)

- Expected total excursion area  $\Sigma \text{area}(z) = \text{area}(D) p_z$ .
- Mean area of maximal region  $> z$ :  $\langle \text{area}(B)(z) \rangle = \Sigma \text{area}(z) / M_z \sim 1/z^2$  or

$$\log \langle \text{area}(B)(z) \rangle \sim -\log z. \quad (\text{Bossew 2010})$$

LN fields:  $\log z \rightarrow z$

- So, the log-survival function  $\ln(1-\text{CDF})$  total area decreases by  $\sim z^2$ , while the mean size of individual maxima, by  $\sim -z$ .
- Mixture of populations: more complicated.



# Questions 1

- Multifractal generated by multiplicative cascades is asymptotically (with generation) LN;
- (1) Total excursion area( $Z>z$ ) behaves as falling parabola in log-log plot for high  $z$ ;
- (2) But according to analysis of Cheng: linear (=power law).
- Mean size of individual anomalous patches indeed falling linear.
- **Solution of discrepancy between (1) and (2)?**



## Questions 2

- Shall certain local maxima be understood (1) as realizations of separate processes, anomaly process on top of BG process? –
- Or (2) as “rare” realization of one process?
- (1) mixture of individual (LN) processes  
(2) extreme outcomes of cascades
- Possibly philosophical question, dependent on scale on which the field is viewed? Because sometimes anomalies appear only at “closer look”.



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# Thank you!



Bundesamt für Strahlenschutz

