

# Coupling subglacial hydrology to basal friction in an Antarctic ice sheet model

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and Frank Pattyn





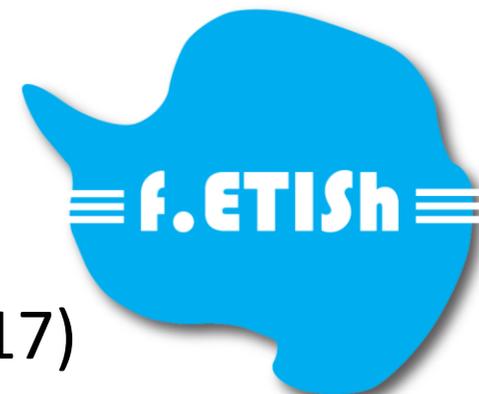
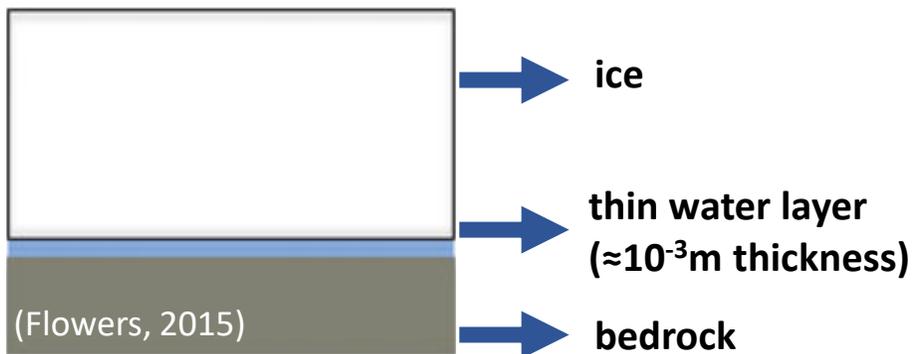
# TAKE-HOME MESSAGE

**Subglacial hydrology and basal sliding** influence the sensitivity of the ice-sheet system on centennial time scales.

In this study we consider different parametrizations and representations of effective pressure and till water content at the base. We also consider different ways to **couple subglacial hydrology to basal sliding**.



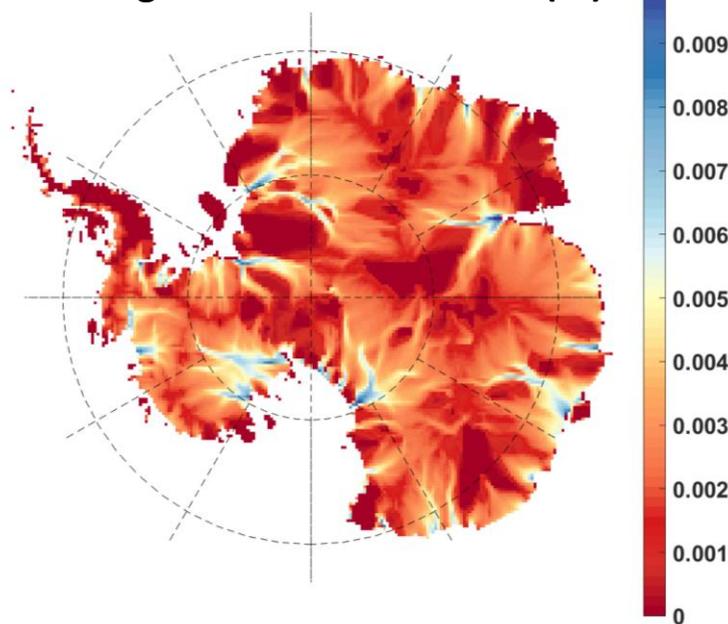
# OUR ICE-SHEET MODEL



## f.ETISH (Pattyn, 2017)

- Hybrid SSA-SIA thermomechanical coupled model
- **Coupled** to a basic subglacial hydrology model (Lebrocq et al., 2009)
  - Water flow = Thin film of water
  - **Different coupling** parametrizations
  - In a **power-law** or a **Coulomb friction law basal sliding**

Subglacial water thickness (m)



Subglacial water thickness obtained from an optimization in a power-law basal sliding with  $m = 2$ .



# COUPLING HYDROLOGY AND SLIDING

4 representations of water content at the base ... coupled with 5 different sliding laws

Sliding controlled by effective pressure ( $N_{eff}$ )

$N_{eff}$  function of the subglacial water pressure

**OPTION 1 (Tsai et al., 2015)**  
→ Details on slide 5

**OPTION 2 (Bueller and Brown, 2009)**  
→ Details on slide 6

$N_{eff}$  function of the amount of liquid water in the till

**OPTION 3 (Bueller and van Pelt, 2015)**  
→ Details on slide 7

Sliding controlled by flux of subglacial water

**OPTION 4 (Goeller et al., 2013)**  
→ Details on slide 8



In a Power-law basal sliding with  $m = 1, 2$  or  $3$



$$\tau_b = \beta^2 v_b$$

$\tau_b$  = driving stress  
 $v_b$  = basal speed

$$\beta^2 = A^{-1} |\tau_d|^{1-m} N_{eff}^p$$

$A$  = basal sliding factor (optimized)  
 $m$  = sliding law exponent  
 $p$  =  $N_{eff}$  exponent



In a Coulomb friction law basal sliding

with 2 different parametrizations of grounding line flux (according to Schoof, 2007a or Tsai et al., 2015)



$$\tau_c = c_0 + (\tan\phi) N_{eff}$$

$\tau_c$  = yield stress  
 $C_0$  = till cohesion (= 0 is further considered)  
 $\Phi$  = till friction angle (optimized)



In a Power-law basal sliding with  $m = 1, 2$  or  $3$



$$\tau_b = \beta^2 v_b$$

$\tau_b$  = driving stress  
 $v_b$  = basal speed



Not yet coupled to a Coulomb friction

$$\beta^2 = A(\phi)^m |\tau_d|^{1-m}$$

$A$  = basal sliding factor (function of subglacial water flux)  
 $\Phi$  = subglacial water flux  
 $m$  = sliding law exponent

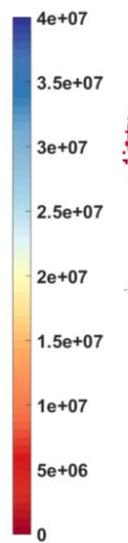


$$N_{\text{eff}} = p_o - p_w = (\rho \cdot g \cdot H) - p_w$$

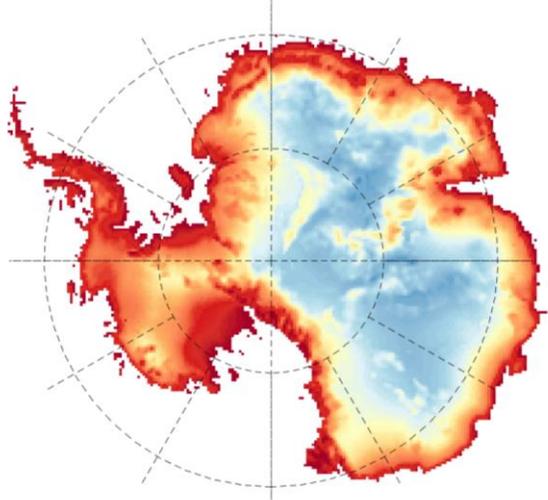
# OPTION 1 (Tsai et al., 2015)

$p_o$  = overburden pressure  
 $p_w$  = subglacial water pressure  
 $\rho$  = ice density  
 $g$  = gravity acceleration  
 $H$  = ice thickness

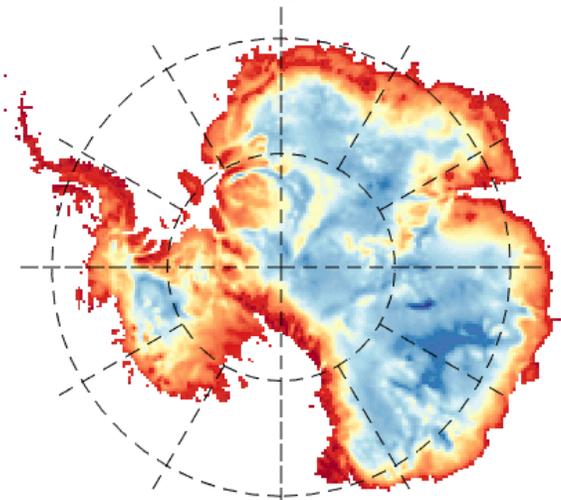
BRUXELLES



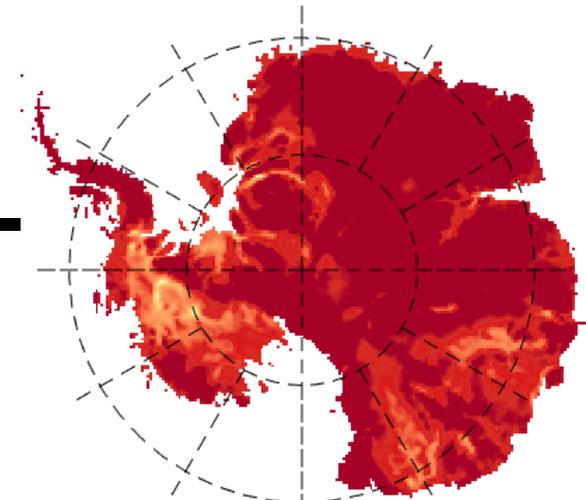
Effective pressure (Pa)



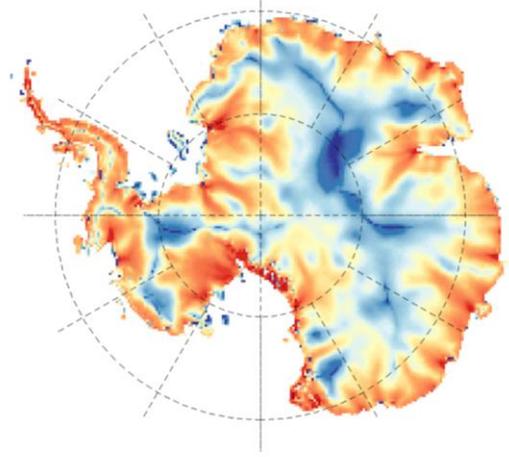
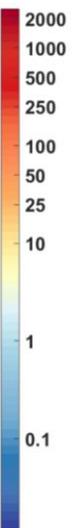
Overburden pressure (Pa)



Subglacial water pressure (Pa)



Basal sliding (m/a)



$p_w = -P_w \rho_{ss} g (b - z_{sl})$   
 $P_w$  = a fixed fraction of the  $p_o$   
 $\rho_{ss}$  = water density  
 $g$  = gravity acceleration  
 $b$  = bedrock elevation  
 $z_{sl}$  = sea level  
 This formula is valid for  $b - z_{sl} < 0$ ,  
 otherwise  $p_w = 0$ .

**Subglacial water pressure is linked to the bed below sea level.**

The figures show different parameters after an optimization in a power-law basal sliding with  $m = 2$ .

$$N_{\text{eff}} = p_o - p_w = (\rho \cdot g \cdot H) - p_w$$

# OPTION 2 (Bueler and Brown, 2009)

$p_o$  = overburden pressure

$\rho$  = ice density

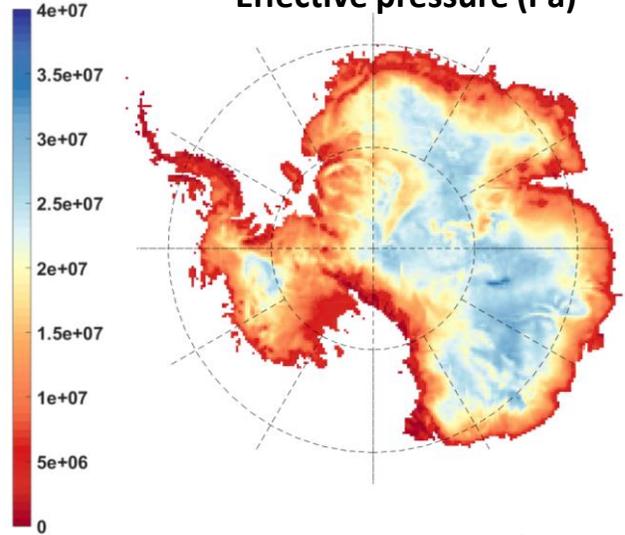
$p_w$  = subglacial water pressure

$g$  = gravity acceleration

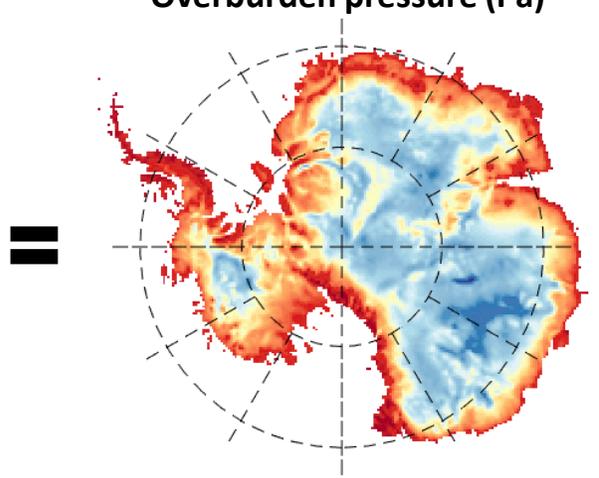
$H$  = ice thickness

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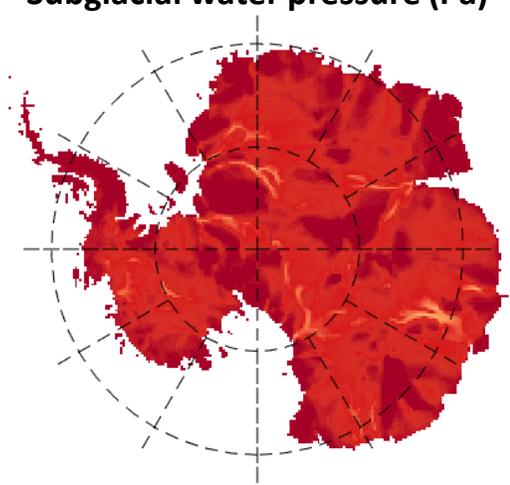
Effective pressure (Pa)



Overburden pressure (Pa)



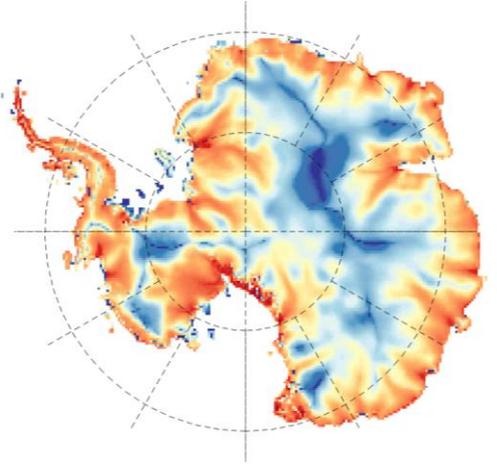
Subglacial water pressure (Pa)



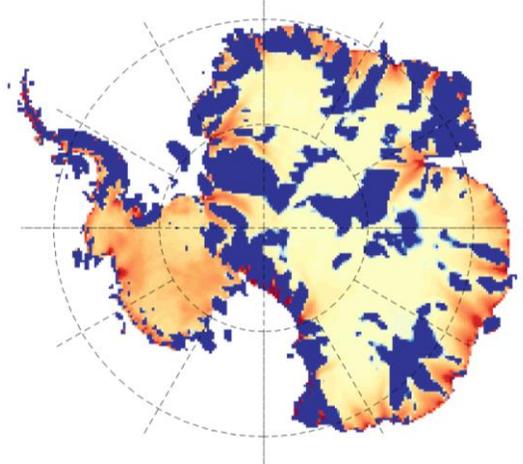
$$p_w = 0.95 \cdot \rho g h \cdot (W/W_{\text{max}})$$

0.95 = fixed fraction of the  $p_o$   
 $W$  = stored liquid water thickness  
 $W_{\text{max}}$  = max saturated till thickness (fixed at 2m) which has an impact on the till weakening by pressurized water

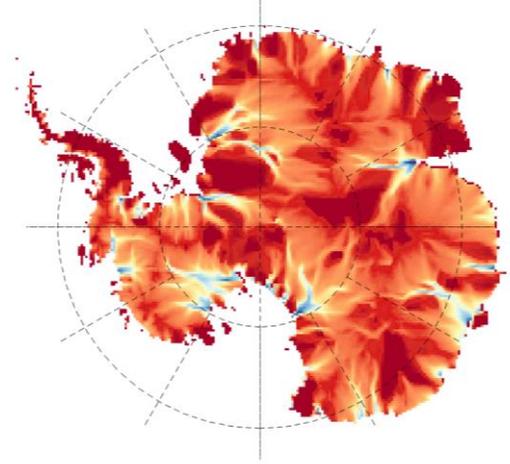
Basal sliding (m/a)



Basal melting (m/a)



Subglacial water thickness (m)



**Subglacial water pressure is a function of the subglacial water thickness.**

The figures show different parameters after an optimization in a power-law basal sliding with  $m = 2$ .

$$N_{\text{eff}} = N_0 * (\delta p_0 / N_0)^s * 10^{(e_0 / C_c)(1-s)}$$

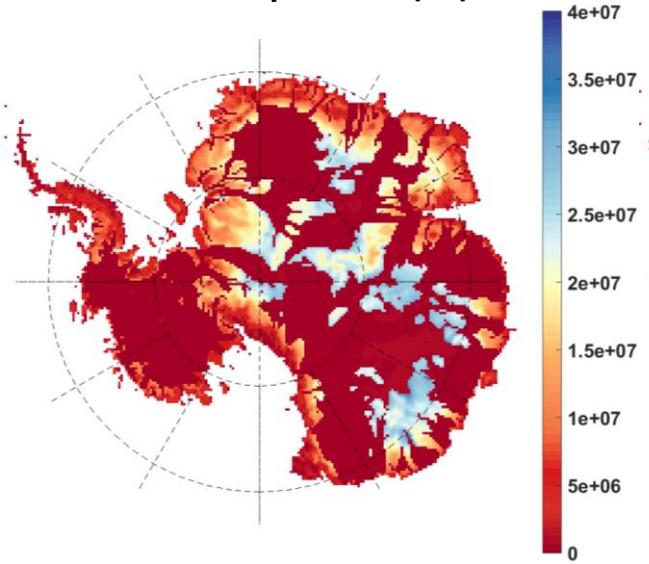
OPTION 3 (Bueler and van Pelt, 2015)

$e_0$  = void ratio at a reference effective pressure  $N_0$   
 $\delta = 0.02$   
 $C_c$  = coefficient of compressibility of the till

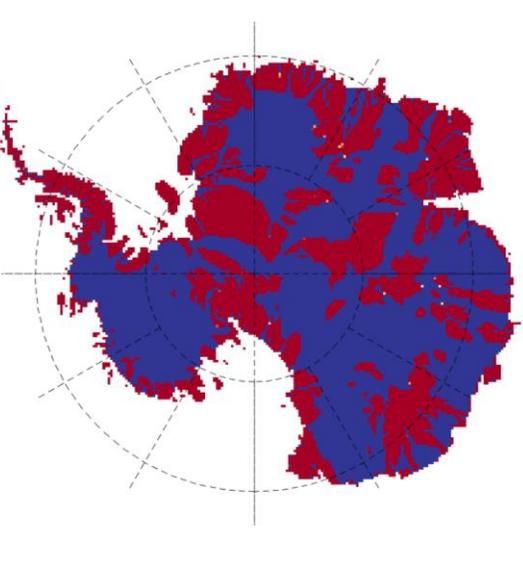
$s = W_{\text{til}} / W_{\text{max}}$   
 $W_{\text{til}}$  = active till layer thickness  
 $W_{\text{max}}$  = max saturated till thickness  
 with boundaries  $N = \min \{p_0, N\}$

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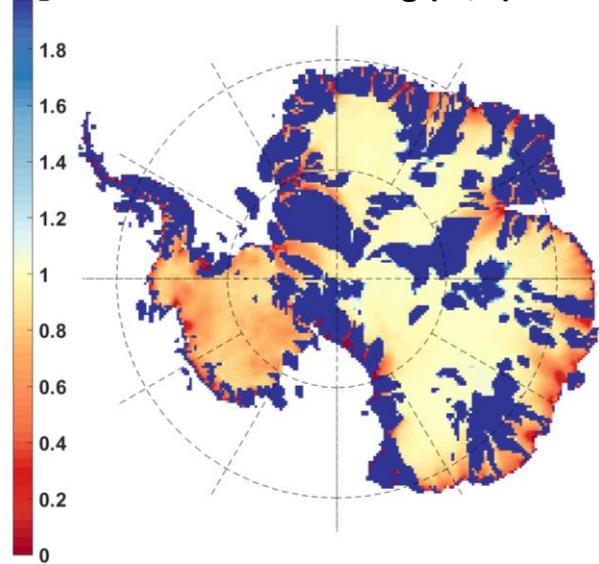
Effective pressure (Pa)



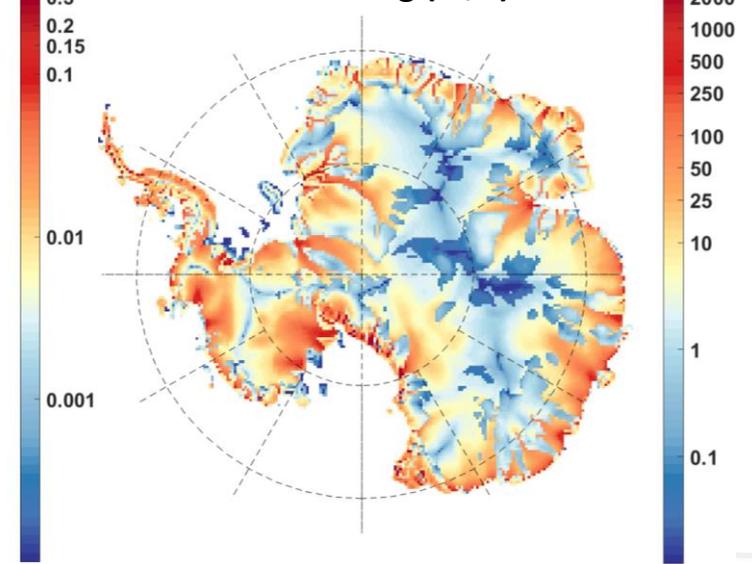
« Active » layer of the till (m)



Basal melting (m/a)



Basal sliding (m/a)



→  $N_{\text{eff}}$  in the case of a deformable bed composed by a permeable till

→ Hydrological model of subglacial water drainage within an active layer of the till :  $W_{\text{til}}$

The figures show different parameters after an optimization in a power-law basal sliding with  $m = 2$ .



$$A(\phi) = A_0 * \exp^{-(\phi/m\phi_0)}$$

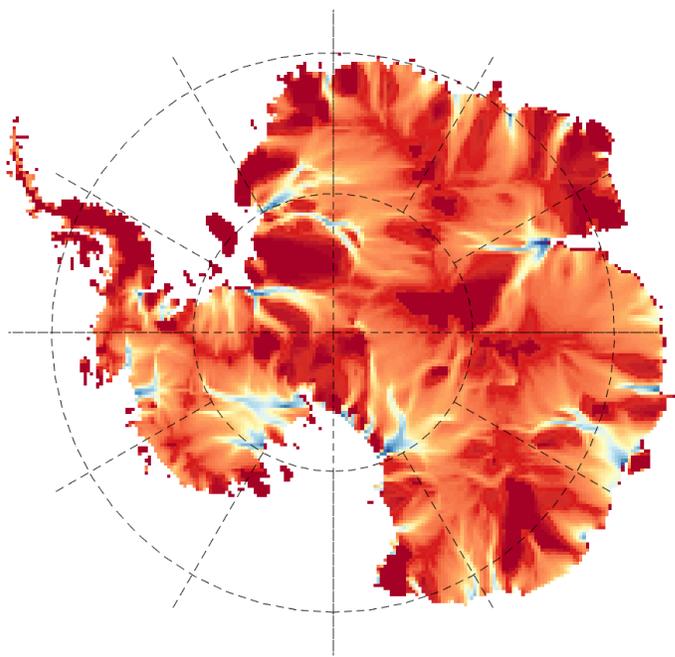
OPTION 4 (Goeller et al., 2013)

$\Phi$  = subglacial water flux

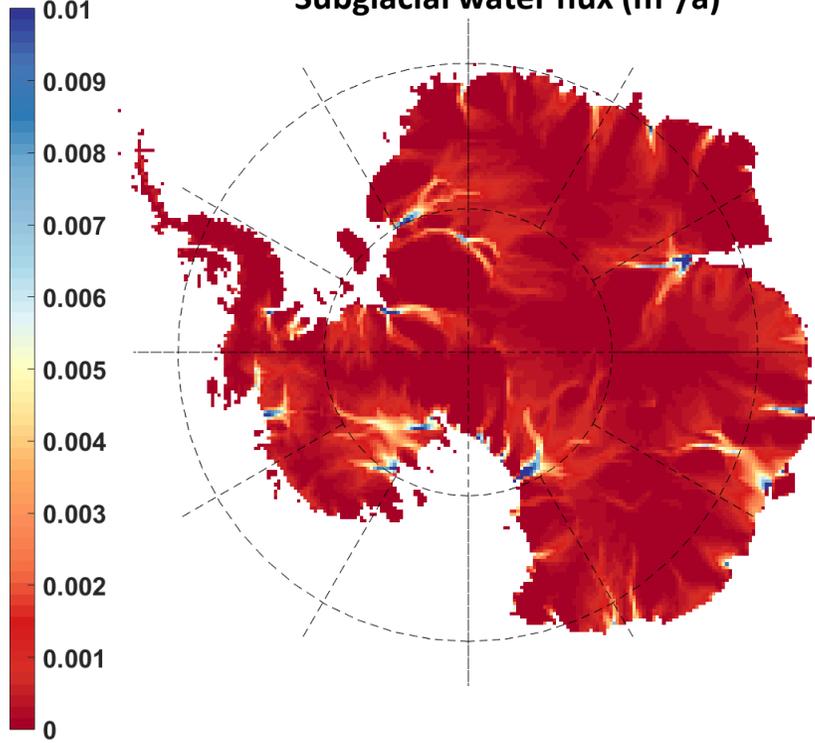
**A** = basal sliding factor

m = sliding law exponent

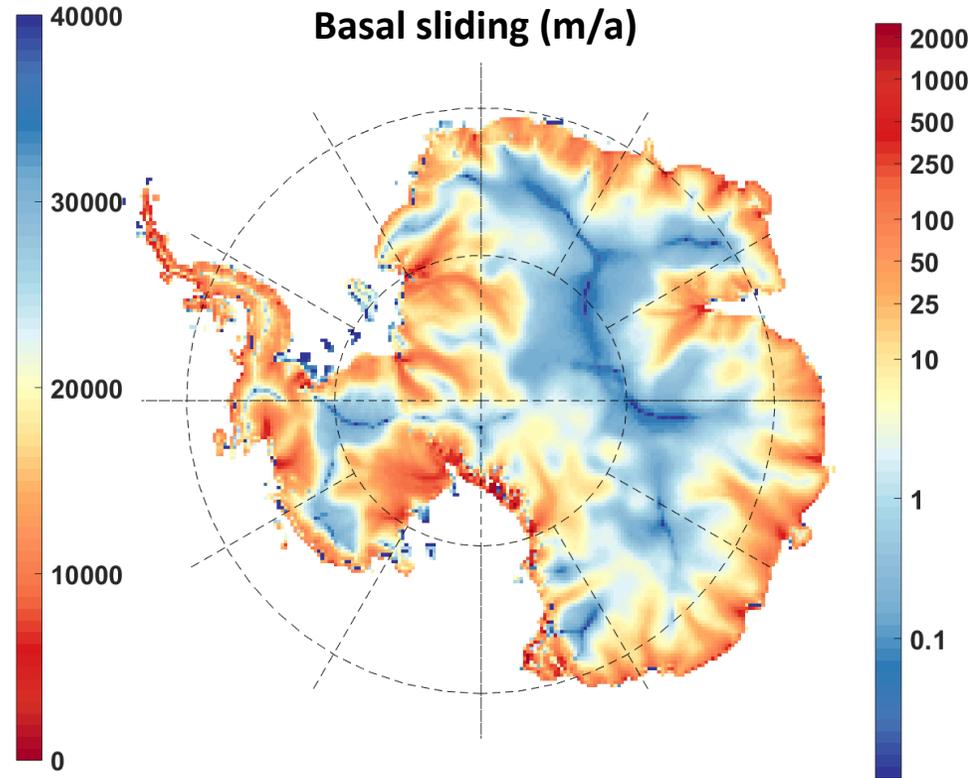
Subglacial water thickness (m)



Subglacial water flux (m<sup>3</sup>/a)



Basal sliding (m/a)



The figures show different parameters after an optimization in a power-law basal sliding with  $m = 2$ .

# EXPERIMENTAL SET-UP

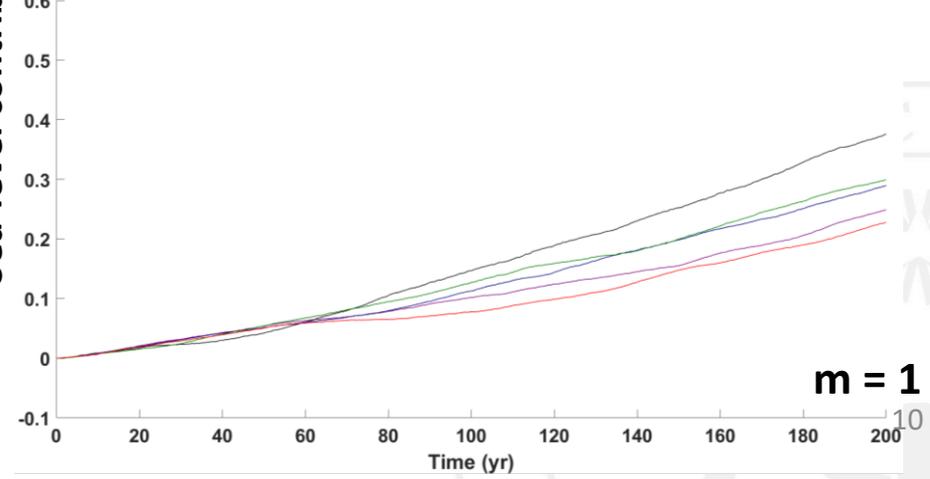
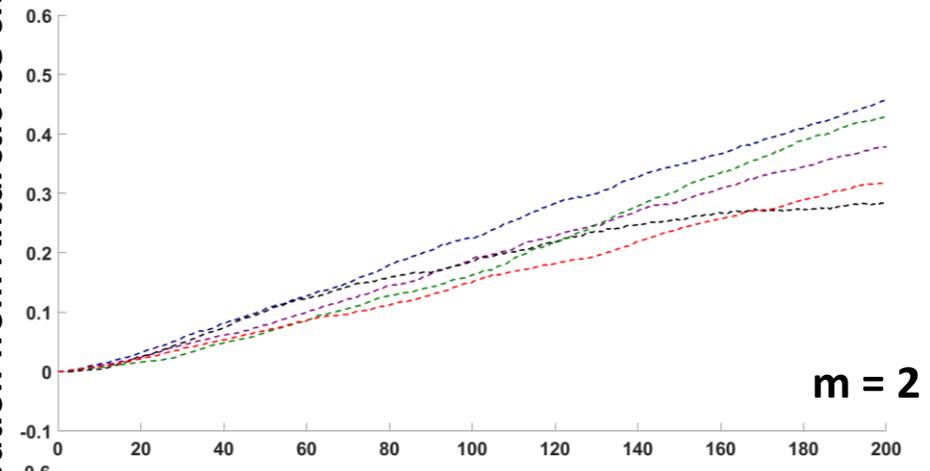
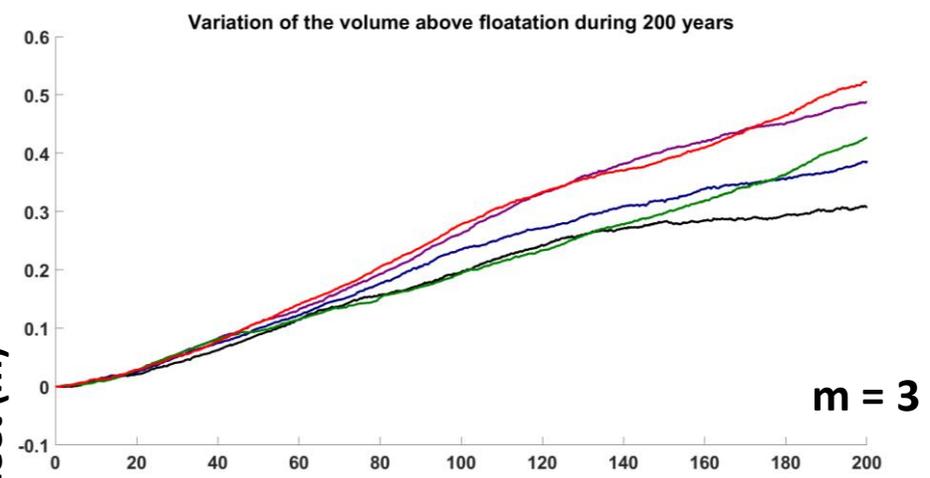
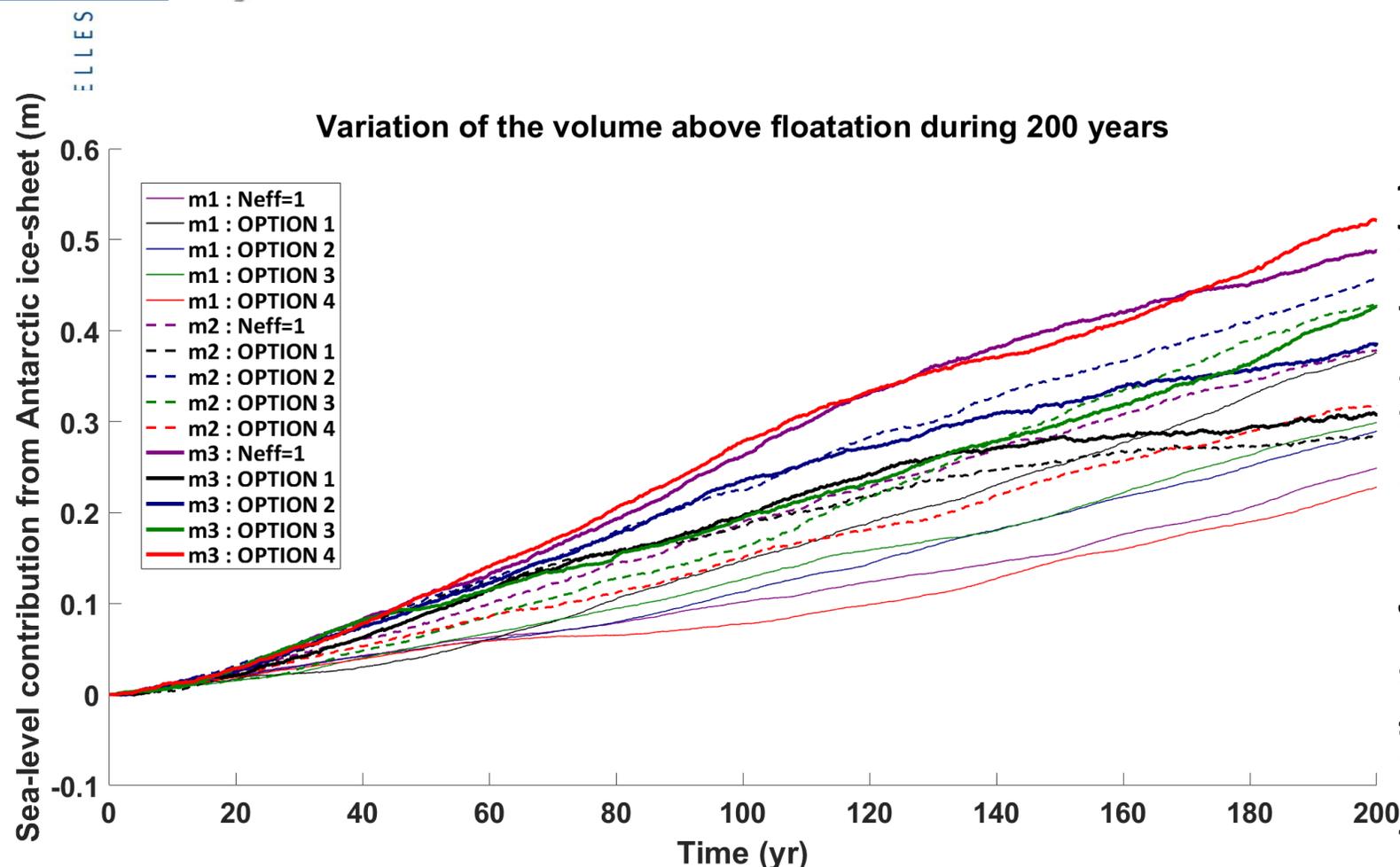
For each of our hydrology-sliding coupling options:

- Forward **200-yr** simulations, starting from present day configuration and forced by subglacial shelf-melt using the **PICO model** (Reese et al., 2018a) enhanced by factor 2 and 4.
  - Model initialised for present-day conditions through **optimization of basal coefficients** (Pollard and DeConto 2012b, Pattyn, 2017);
  - **Input data: present day** ice-sheet geometry (Bedmachine), surface mass balance (RACMO2), surface temperature (RACMO2) and geothermal heat flux (Shapiro and Ritzwoller);
  - **25 km** resolution.





# RESULTS: POWER-LAW SLIDING

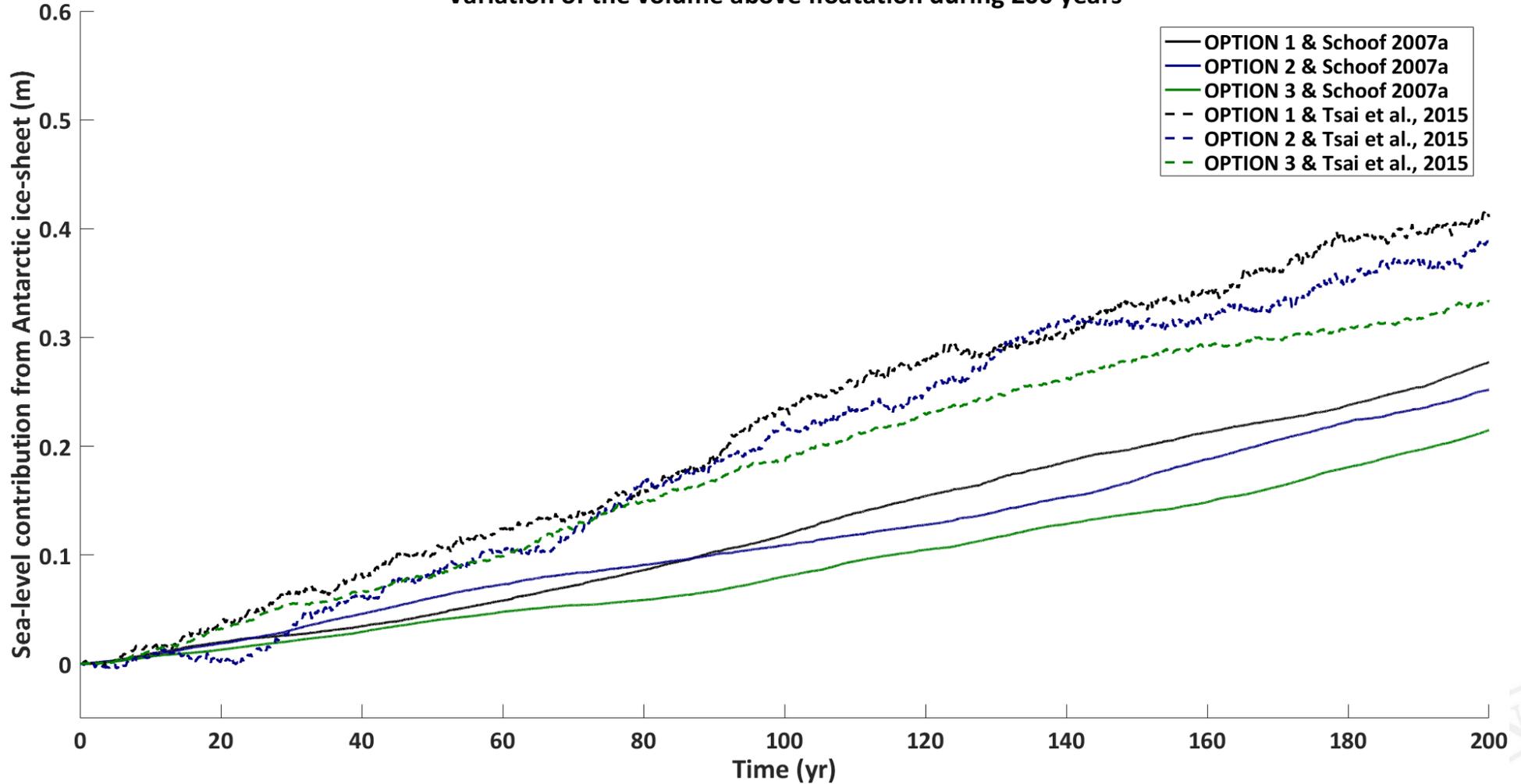


- The variation of the **volume above floatation** is a good indicator of **ice-sheet sensitivity**
- To avoid the signal of different optimizations, we substracted the forcing experiment enhanced by factor 2 to the forcing experiment enhanced by factor 4
- The **purple curve** represents an Neff=1, wich means **no effects of subglacial hydrology**

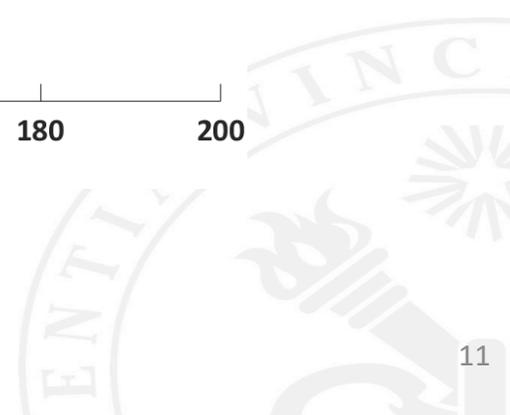


# RESULTS : COULOMB FRICTION LAW SLIDING

Variation of the volume above floatation during 200 years



- The variation of the **volume above floatation** is a good indicator of **ice-sheet sensitivity**
- To avoid the signal of different optimizations, we subtracted the forcing experiment enhanced by factor 2 to the forcing experiment enhanced by factor 4



# CONCLUSION

Recent studies show that the **ice sheet sensitivity** increases with the **sliding law power** or the use of a **Coulomb friction law**.

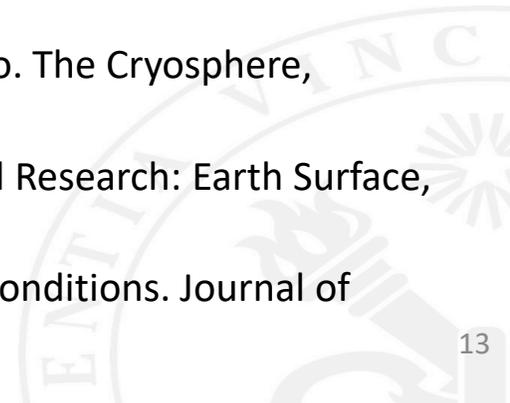
**Coupling with subglacial hydrology** exhibits a more **complex relationship**.





# References

- **E. Bueler and J. Brown.** Shallow shelf approximation as a “sliding law” in a thermomechanically coupled ice sheet model. *Journal of Geophysical Research: Earth Surface*, 114(F3),2009. ISSN 2156-2202. doi: 10.1029/2008JF001179. F03008.
- **E. Bueler and W. van Pelt.** Mass-conserving subglacial hydrology in the parallel ice sheet model version 0.6. *Geoscientific Model Development*, 8(6):1613–1635, 2015. doi: 10.5194/gmd-8-1613-2015.
- **G. E. Flowers.** Modelling water flow under glaciers and ice sheets. *Proceedings of the Royal Society A: Mathematical, Physical and Engineering Sciences*,471(2176), 20140907, 2015.
- **S. Goeller, M. Thoma, K. Grosfeld, and H. Miller.** A balanced water layer concept for subglacial hydrology in large-scale ice sheet models. *The Cryosphere*, 7(4), 1095-1106, 2013.
- **A. Le Brocq, A. Payne, M. Siegert, and R. Alley.** A subglacial water-flow model for west antarctica. *Journal of Glaciology*, 55(193):879–888, 2009. doi: 10.3189/002214309790152564.
- **F. Pattyn.** Sea-level response to melting of Antarctic ice shelves on multi-centennial timescales with the fast Elementary Thermomechanical Ice Sheet model (f. ETISh v1. 0). *Cryosphere*, 11(4), 2017.
- **D. Pollard and R. M. DeConto.** A simple inverse method for the distribution of basal sliding coefficients under ice sheets, applied to antarctica. *The Cryosphere*, 6(5):953–971, 2012b. doi: 10.5194/tc-6-953-2012.
- **R. Reese, T. Albrecht, M. Mengel, X. Asay-Davis, and R. Winkelmann.** Antarctic sub-shelf melt rates via pico. *The Cryosphere*, 12(6):1969–1985, 2018a. doi: 10.5194/tc-12-1969-2018.
- **C. Schoof.** Ice sheet grounding line dynamics: Steady states, stability, and hysteresis. *Journal of Geophysical Research: Earth Surface*, 112(F3):n/a–n/a, 2007a. ISSN 2156-2202. doi:10.1029/2006JF000664. F03S28.
- **V. C. Tsai, A. L. Stewart, and A. F. Thompson.** Marine ice-sheet profiles and stability under coulomb basal conditions. *Journal of Glaciology*, 61(226):205–215, 2015. doi: doi:10.3189/2015JoG14J221.



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