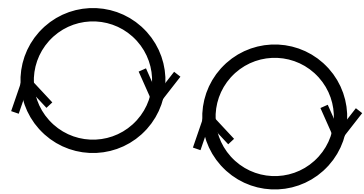
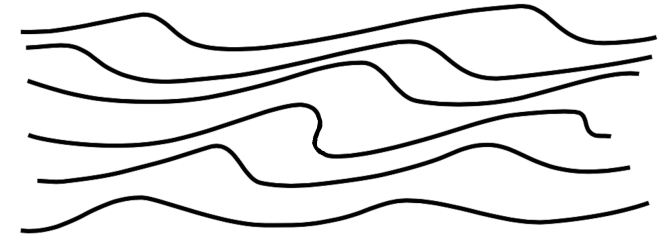
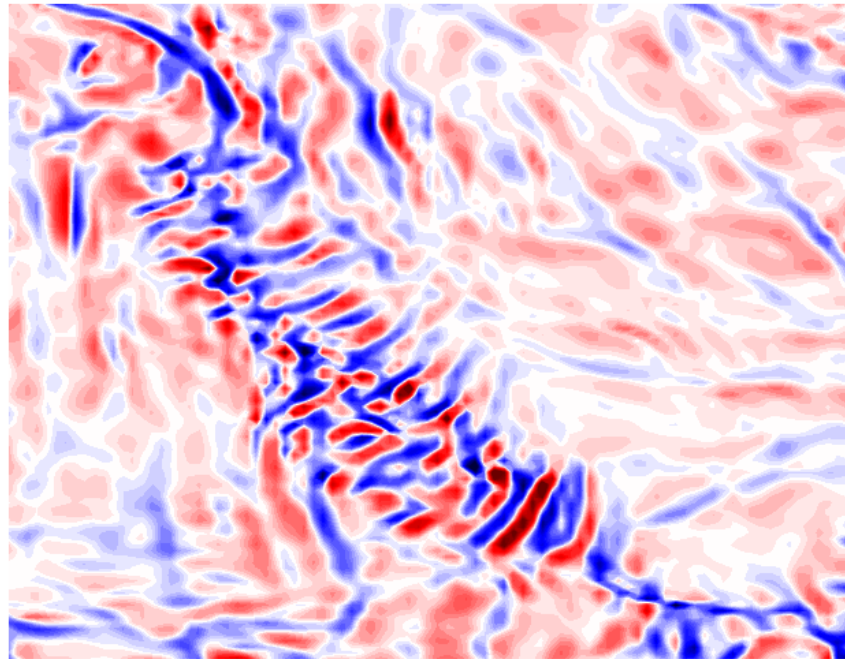


Energy Transfers Between Balanced And Unbalanced Motions In Geophysical Flows



Mesoscale eddies
(balanced)



Internal gravity waves
(unbalanced)

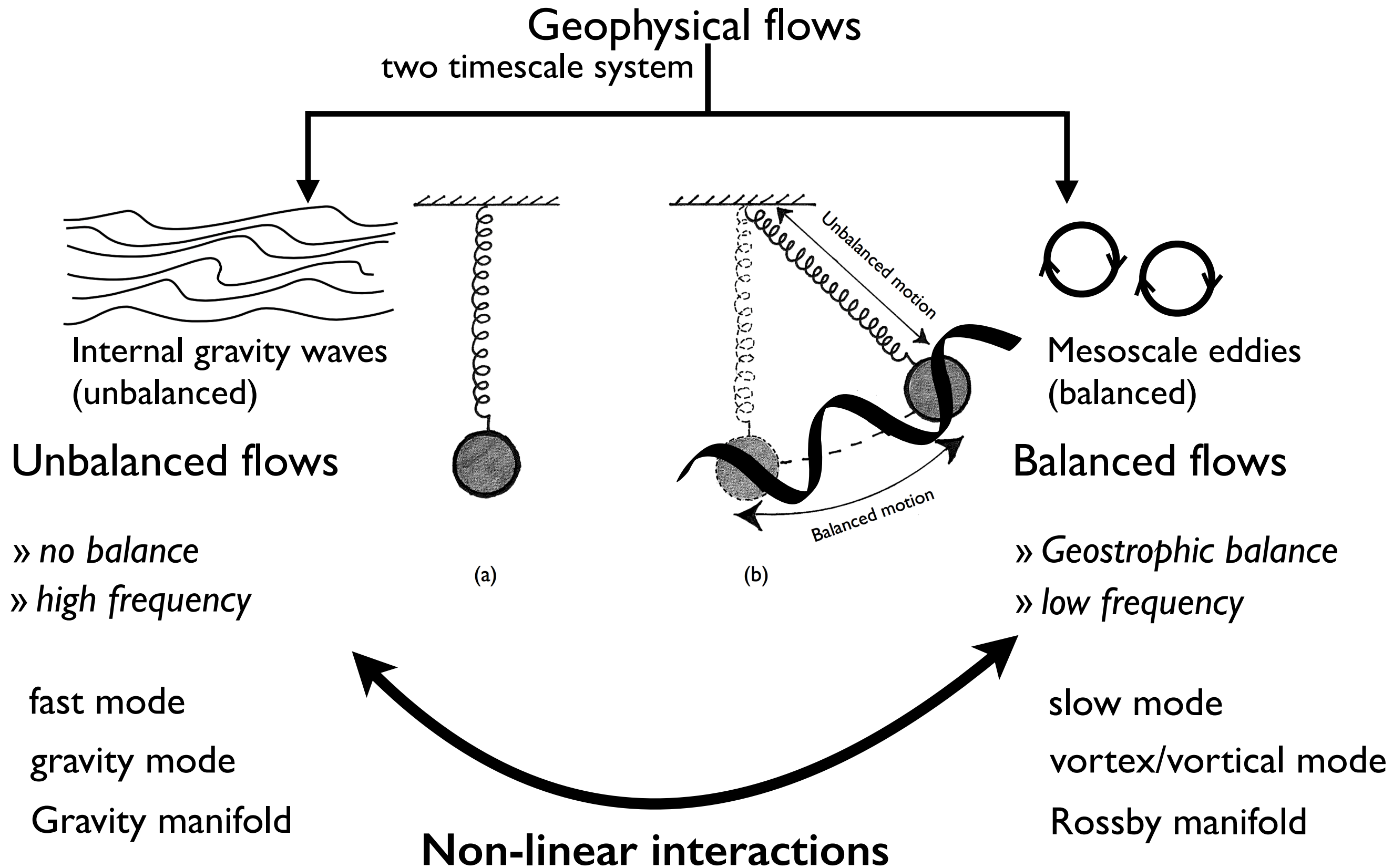
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Session NPI.I Mathematics of Planet Earth

EGU 2020, May 4th 2020

Balanced and unbalanced motions



Different dynamical regimes

Rossby number

$$Ro = \frac{\text{flow frequency}}{\text{frequency of rotation}}$$

- ▶ Weakly stratified
 - ▶ e.g. mixed layer

- ▶ Strongly stratified
 - ▶ e.g. ocean interior

$Ro < 1$

$Ro \ll 1$

Ageostrophic

Quasi-geostrophic

Balanced and unbalanced motions are:

- ▶ **Strongly coupled**
- ▶ Timescale separation is complex

- ▶ **Weakly coupled**
- ▶ Timescales well-separated

Different dynamical regimes in the Ocean

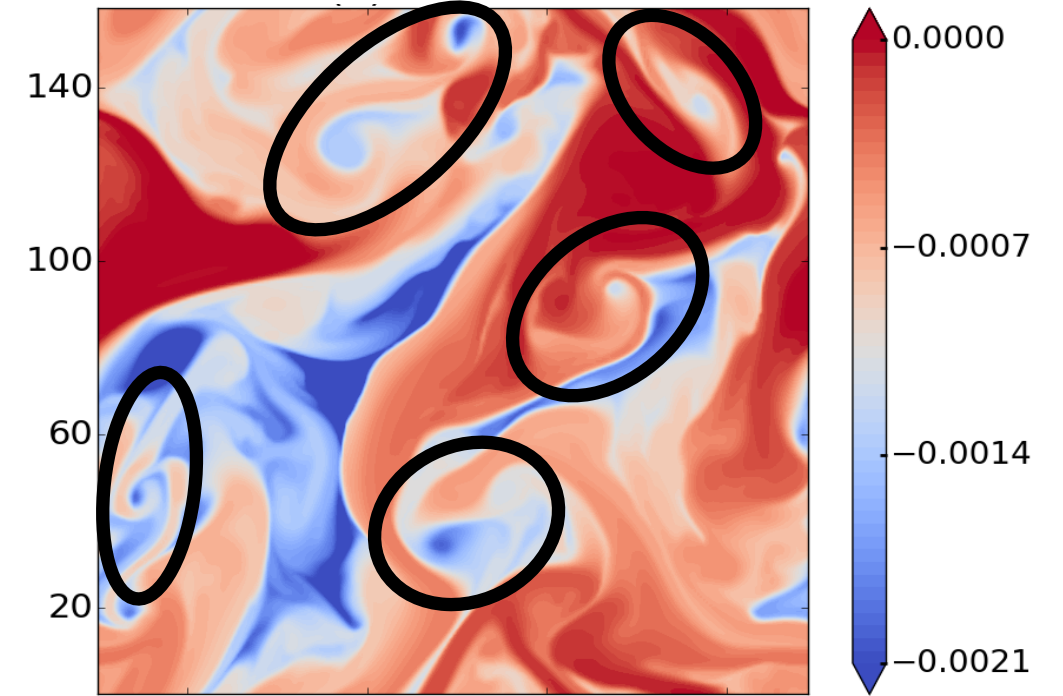
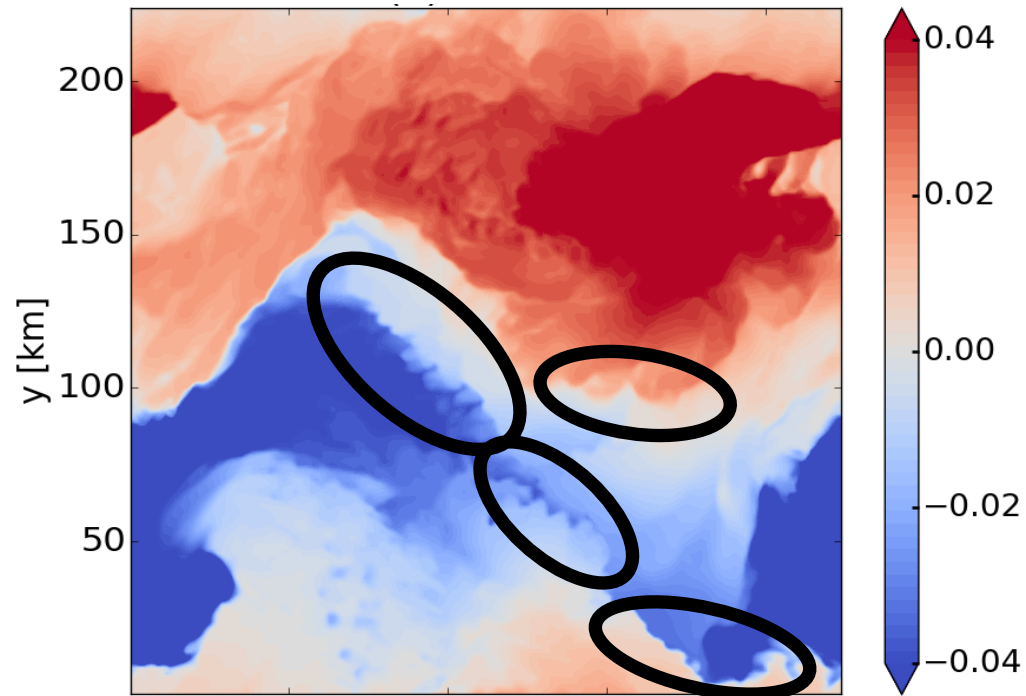
$Ro < 1$

(Weakly stratified)

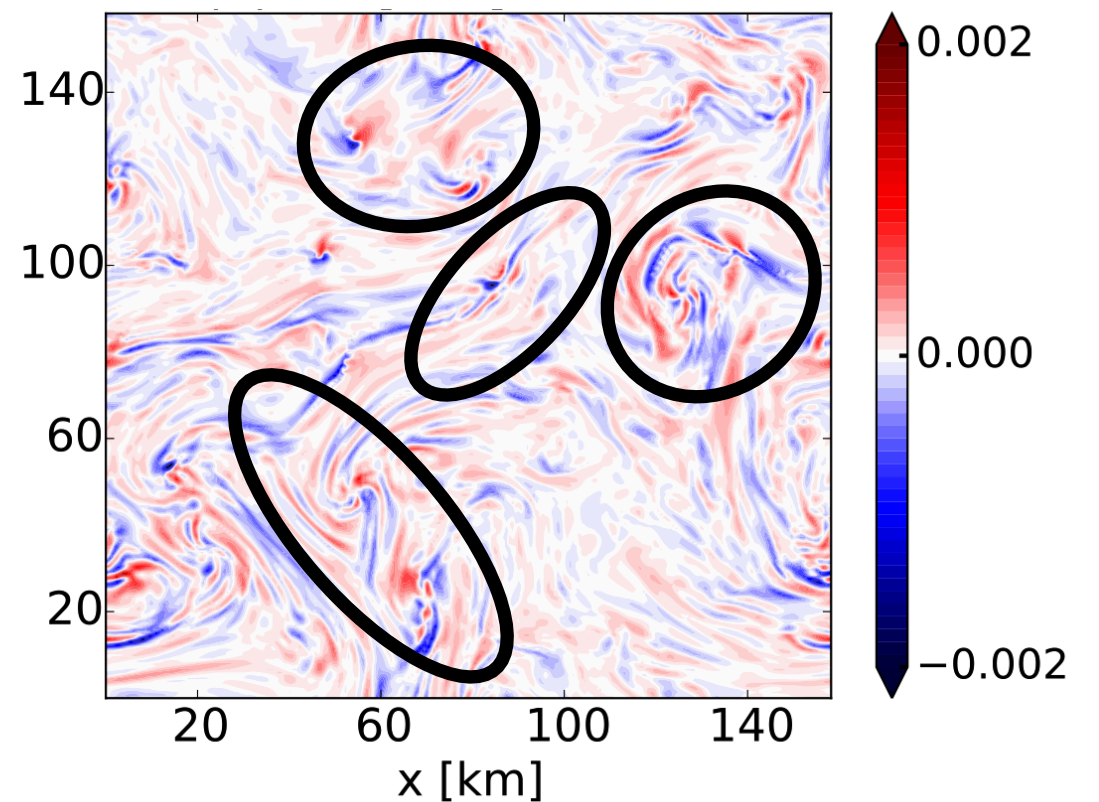
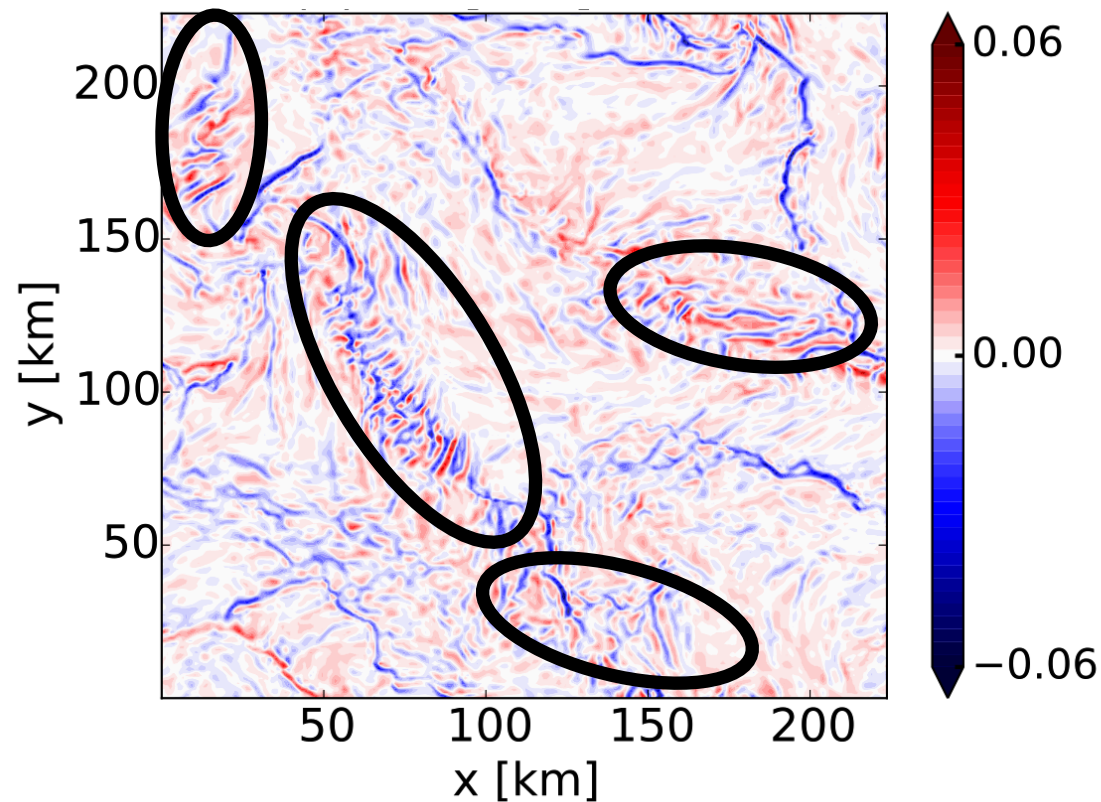
$Ro \ll 1$

(Strongly stratified)

Buoyancy



Vertical velocity



Flow decomposition methods

Non-linear normal mode initialization (NNMI)

- » **Machenhauer (1977)**: Initially developed for Numerical Weather Prediction
 - » The idea is to initially suppress gravity waves to minimize excitation
 - » **Leith (1980)**: Quasi-geostrophic balanced state (1st order in Ro), first iteration
- » **Warn et. al (1995)**: higher order in Ro (nth order)

Optimal balance

- » Masur and Oliver, 2020, JGAFD : Optimal potential vorticity (OPV) balance based on Viúdez and Dritschel (2004) for a quasi-geostrophic balanced state
 - » Iterative procedure
 - » Ramp time to match a target potential vorticity

Modal decomposition

Single layer model (scaled): $\partial_t \mathbf{u} + \frac{1}{Ro} \mathbf{u} + \nabla h = -Ro \mathbf{u} \cdot \nabla \mathbf{u} \quad \partial_t h + c^2 \nabla \cdot \mathbf{u} = -Ro \nabla \cdot h \mathbf{u}$

Fourier space:

$$\partial_t \hat{\mathbf{z}} - \underbrace{i\mathbf{A} \cdot \hat{\mathbf{z}}}_{\text{Linear}} = \underbrace{Ro \hat{n}}_{\text{Non-linear}}$$

vector $\hat{\mathbf{z}}(\mathbf{k}) = (\hat{u}, \hat{v}, \hat{h})^T$

$$\mathbf{A} = \begin{pmatrix} 0 & -i & -k_x \\ i & 0 & -k_y \\ -c^2 k_x & -c^2 k_y & 0 \end{pmatrix}$$

Modal decomposition

Single layer model (scaled): $\partial_t \mathbf{u} + \frac{\mathbf{u}}{f} + \nabla h = -Ro \mathbf{u} \cdot \nabla \mathbf{u} \quad \partial_t h + c^2 \nabla \cdot \mathbf{u} = -Ro \nabla \cdot h \mathbf{u}$

Fourier space:

vector $\hat{\mathbf{z}}(\mathbf{k}) = (\hat{u}, \hat{v}, \hat{h})^T$

$$\partial_t \hat{\mathbf{z}} - \underbrace{i\mathbf{A} \cdot \hat{\mathbf{z}}}_{\text{Linear}} = \underbrace{Ro \hat{\mathbf{n}}}_{\text{Non-linear}}$$

$$\mathbf{A} = \begin{pmatrix} 0 & -i & -k_x \\ i & 0 & -k_y \\ -c^2 k_x & -c^2 k_y & 0 \end{pmatrix}$$

Balanced mode

Unbalanced mode

from C-grid
discrete
operators

Eigenvalues: $\omega^0 = 0$

$$\omega^\pm = \pm \sqrt{1 + c^2 k^2}$$

Eigenvectors: $\mathbf{q}^0, \mathbf{p}^0$

$$\mathbf{q}^\pm, \mathbf{p}^\pm$$

Projection: mode amplitude $g^s = \mathbf{p}^s \cdot \hat{\mathbf{z}}$ with $s = 0, \pm$

$$\partial_t g^s - i\omega^s g^s = Ro \mathbf{p}^s \cdot \hat{\mathbf{n}} = -iRo I^s(g^0, g^\pm)$$

$$I^s(g^0, g^\pm) = I^s(g^0, 0) + I^s(0, g^\pm) + K^s(g^0, g^\pm)$$

Modal decomposition

» **Modal representation:** $\partial_t g^s - i\omega^s g^s = Ro p^s \cdot \hat{n} = -iRo I^s(g^0, g^\pm)$
 $(Ro \partial_T + \partial_{t*}) g^s - i\omega^s g^s = -iRo (I^s(g^0, 0) + I^s(0, g^\pm) + K^s(g^0, g^\pm))$

- » **Weak interaction assumption:** weakly growing waves $g^\pm = Ro f_1^\pm + Ro^2 f_2^\pm + \dots$
- » expansion in Ro as e.g. in Warn (1996), Kafiabad and Bartello (2017)
- » introduce fast and slow time scale with $T = Ro t^*$ and $\partial_t = Ro \partial_T + \partial_{t*}$
- » slow mode g^0 varies on T only, while fast mode g^\pm has two time scales t^* and T

Modal decomposition

» **Modal representation:** $\partial_t g^s - i\omega^s g^s = Ro p^s \cdot \hat{n} = -iRo I^s(g^0, g^\pm)$
 $(Ro \partial_T + \partial_{t*}) g^s - i\omega^s g^s = -iRo (I^s(g^0, 0) + I^s(0, g^\pm) + K^s(g^0, g^\pm))$

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» expansion in Ro as e.g. in Warn (1996), Kafiabad and Bartello (2017)

» introduce fast and slow time scale with $T = Ro t^*$ and $\partial_t = Ro \partial_T + \partial_{t^*}$

» slow mode g^0 varies on T only, while fast mode g^\pm has two time scales t^* and T

» **SLOW MODE $s=0$**

for increasing order in Ro :

» **FAST MODE $s=\pm$**

$$\partial_T g^0 = -il^s(g^0, 0)$$

$$\partial_T g^0 = -il^s(g^0, f_1^\pm) + il^s(0, f_1^\pm)$$

$$\partial_T g^0 = -il^s(g^0, f_2^\pm) + il^s(0, f_2^\pm) - il^0(0, f_1^\pm)$$

$$\partial_{t^*} f_1^\pm - i\omega^\pm f_1^\pm = -il^\pm(g^0, 0)$$

$$\partial_T f_1^\pm + \partial_{t^*} f_2^\pm - i\omega^\pm f_2^\pm = -iK^\pm(g^0, f_1^\pm)$$

$$\partial_T f_2^\pm + \partial_{t^*} f_3^\pm - i\omega^\pm f_3^\pm = -il^\pm(0, f_1^\pm) - iK^\pm(g^0, f_2^\pm)$$

» suppress any wave generation by $\partial_{t^*} f_n^\pm = 0 \rightarrow$ 'slaved' modes f_n

» **Machenauer(1977)**

» **QG balanced state**

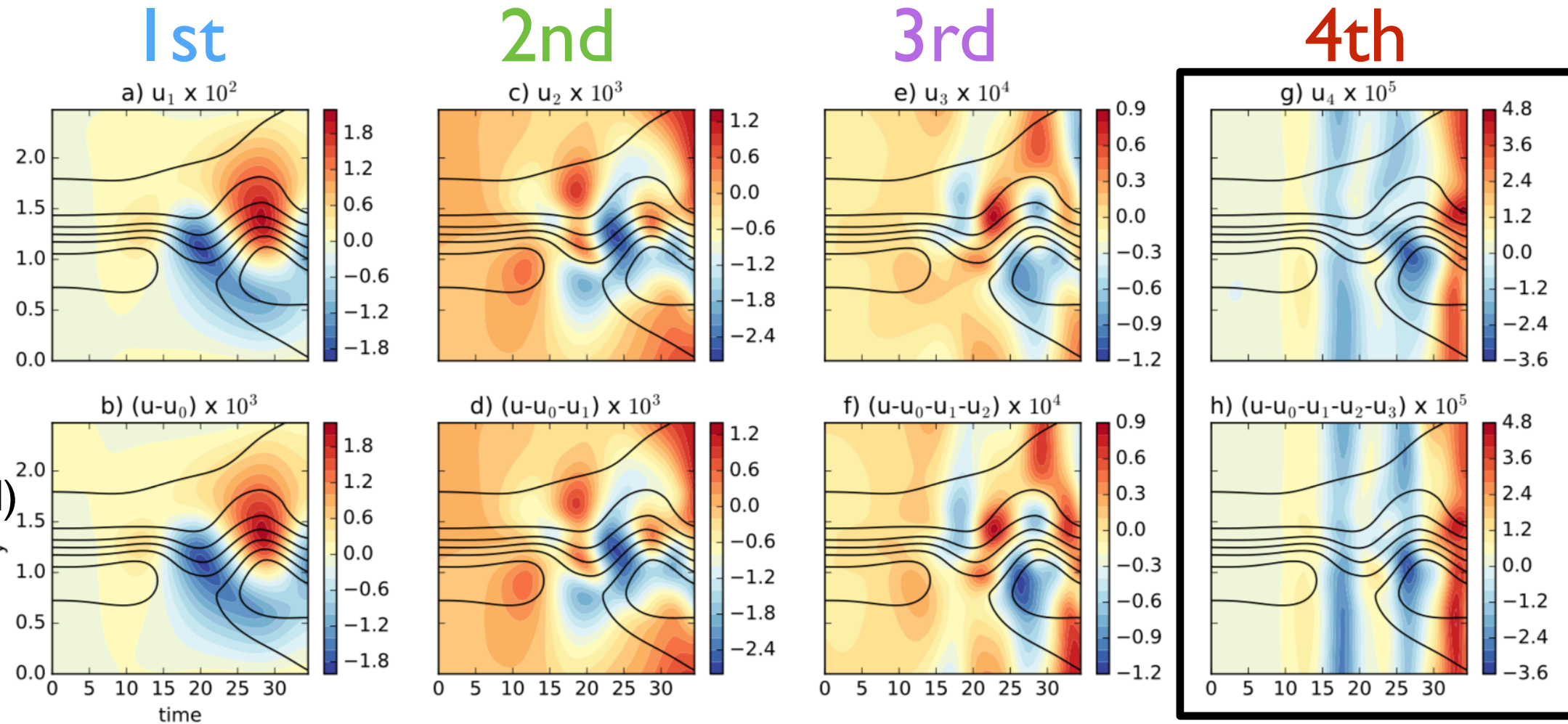
$$f_1^\pm = I^\pm(g^0, 0)/\omega^\pm, f_2^\pm = (K^\pm(g^0, f_1^\pm) - i\partial_T f_1^\pm)/\omega^\pm, \dots$$

» **first order slaved mode**

Wave emission at higher orders: single layer model

Ro = 0.1

Balanced velocity



- » double periodic domain 10x5 (dimensionless)
- » initialized with an unstable zonal jet

- » zonal jet meanders and dissolves into eddies
- » range of $Ro = 0.02$ to 0.3 , i.e. from mesoscale to sub-mesoscale conditions

» wave signal seen: is not related to spontaneous emission by shear instability of the balanced flow

» rather, the wave signal could be related to local Rossby numbers with $Ro_{local} > 1$: symmetric instability becomes possible

(apparent) Wave emission at higher orders: primitive equation model

1st

2nd

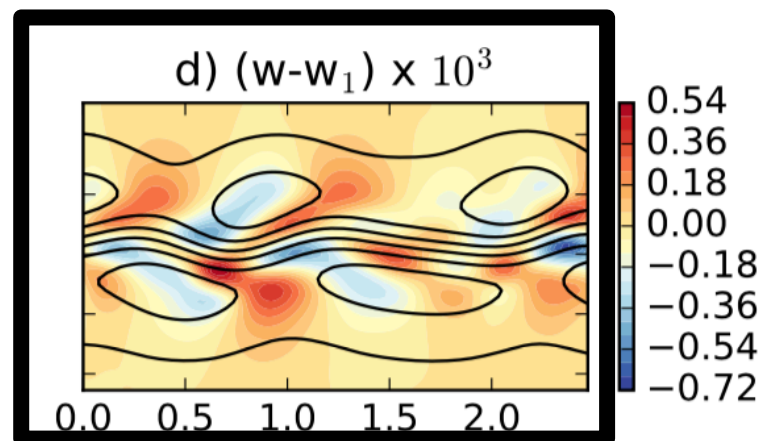
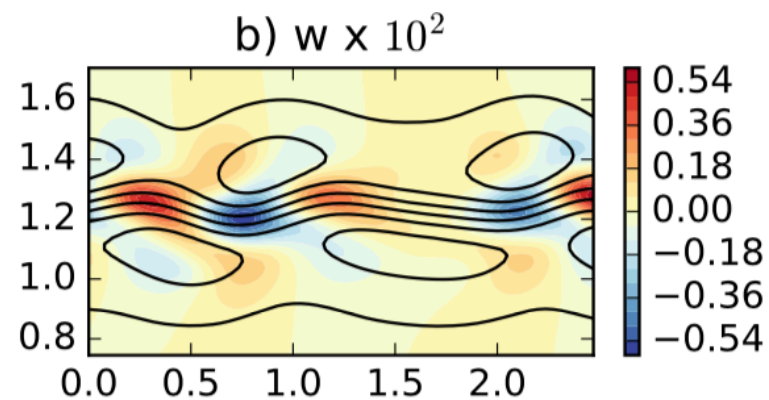
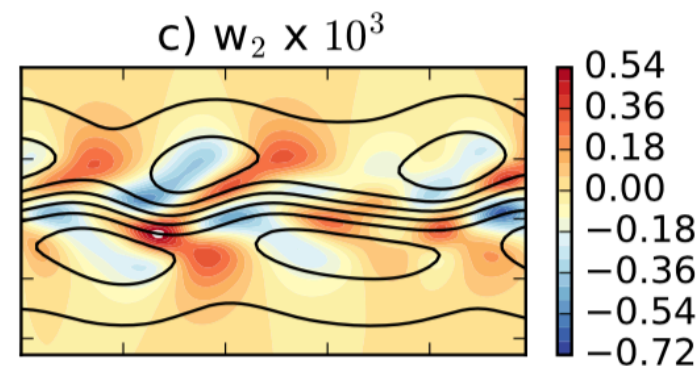
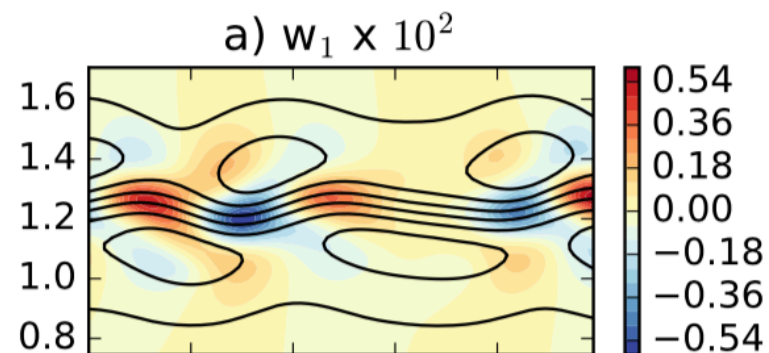
Apparent spontaneous emission
or 1st or 2nd order slaved modes?

Balanced velocity

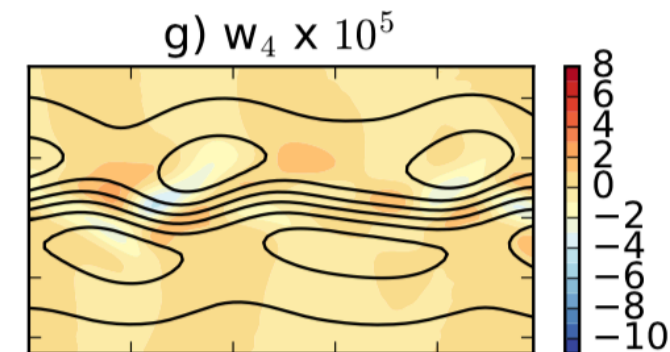
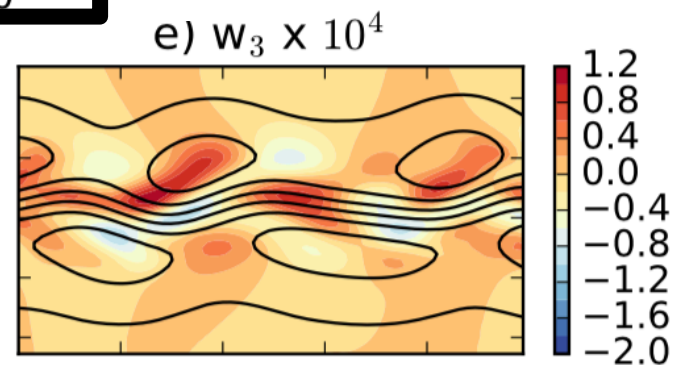
Residual (unbalanced)
velocity

3rd

4th

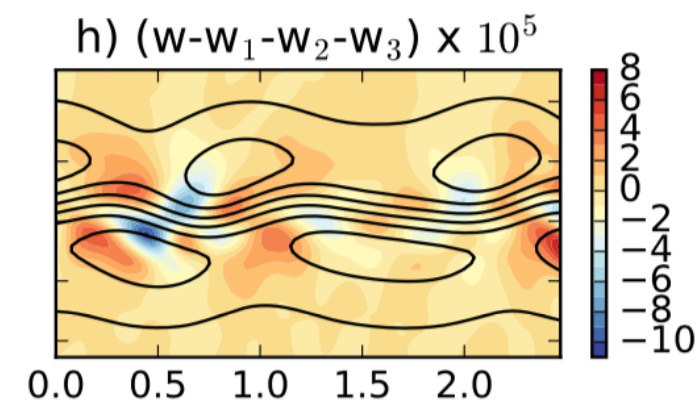
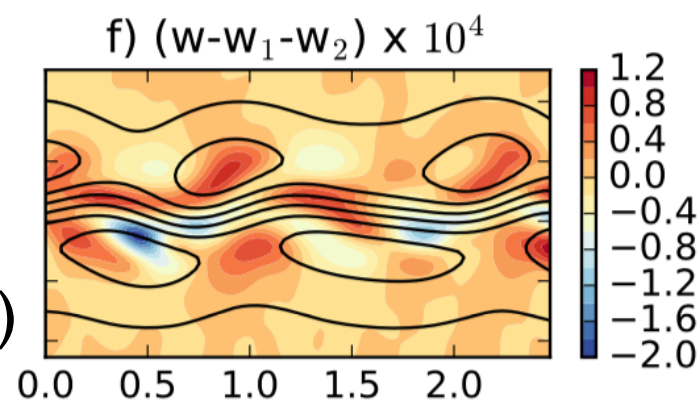


Balanced velocity



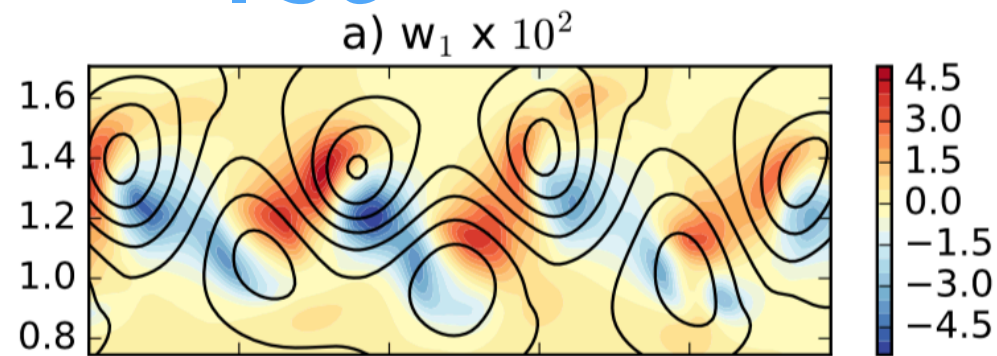
Weak wave signal is seen
only at 4th order.

Residual (unbalanced)
velocity



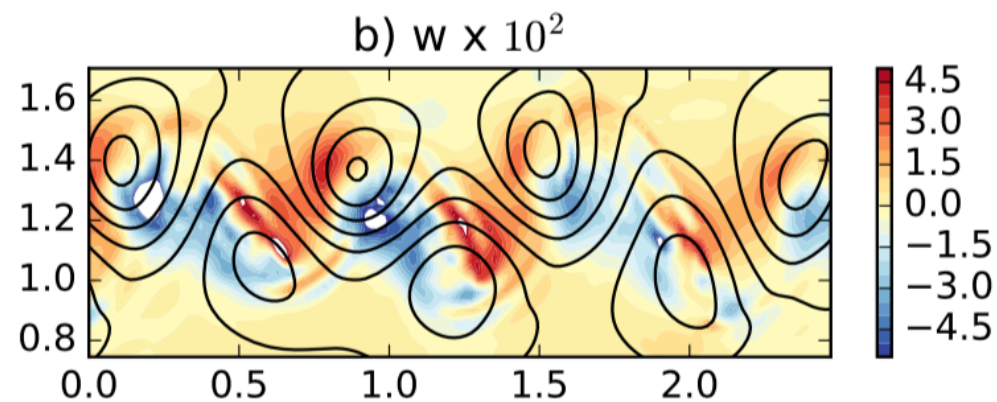
Convectively generated waves at higher orders: primitive equation model

1st



Balanced velocity

Stronger wave signal seen already at 2nd order.

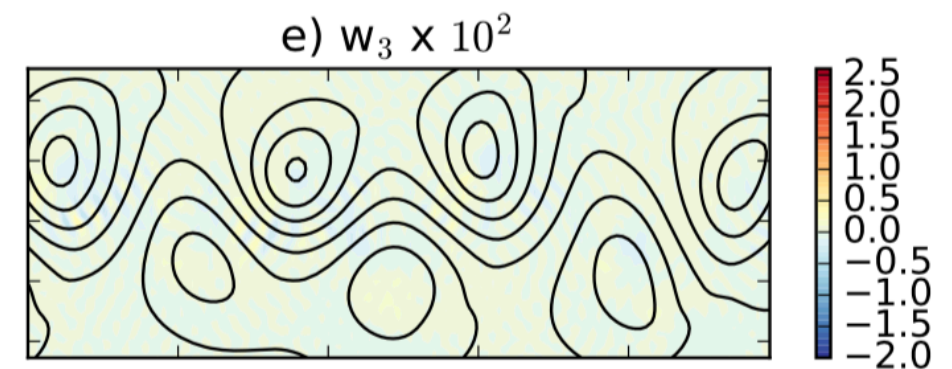
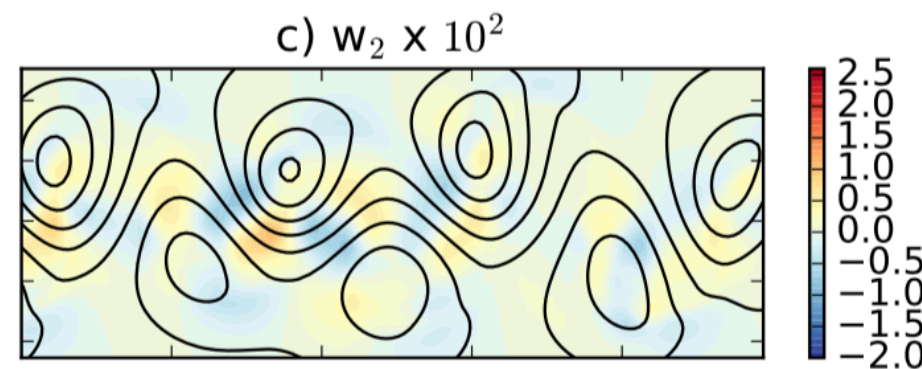


Residual (unbalanced) velocity

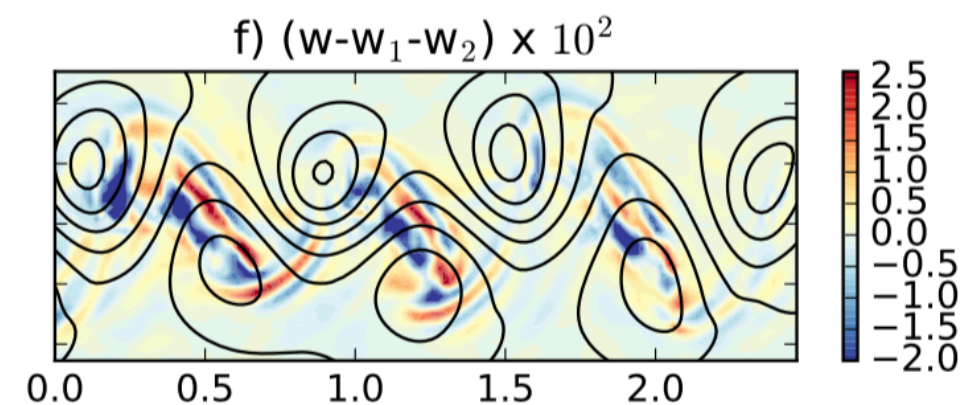
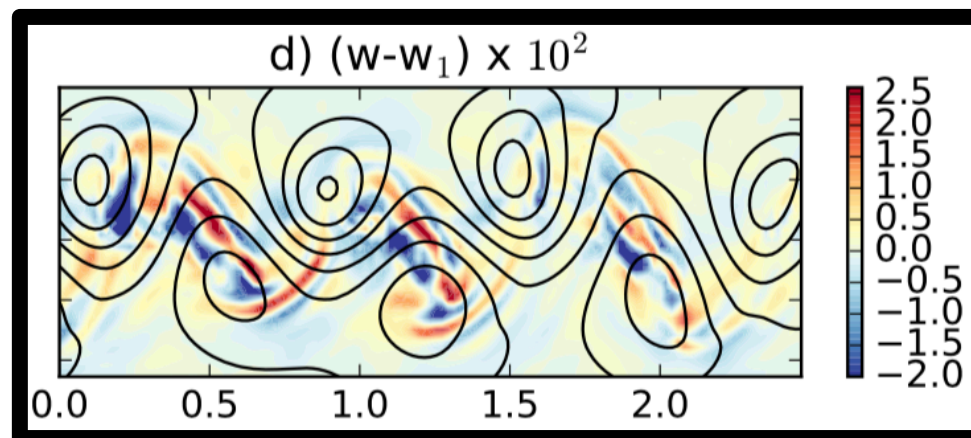
2nd

3rd

Balanced velocity

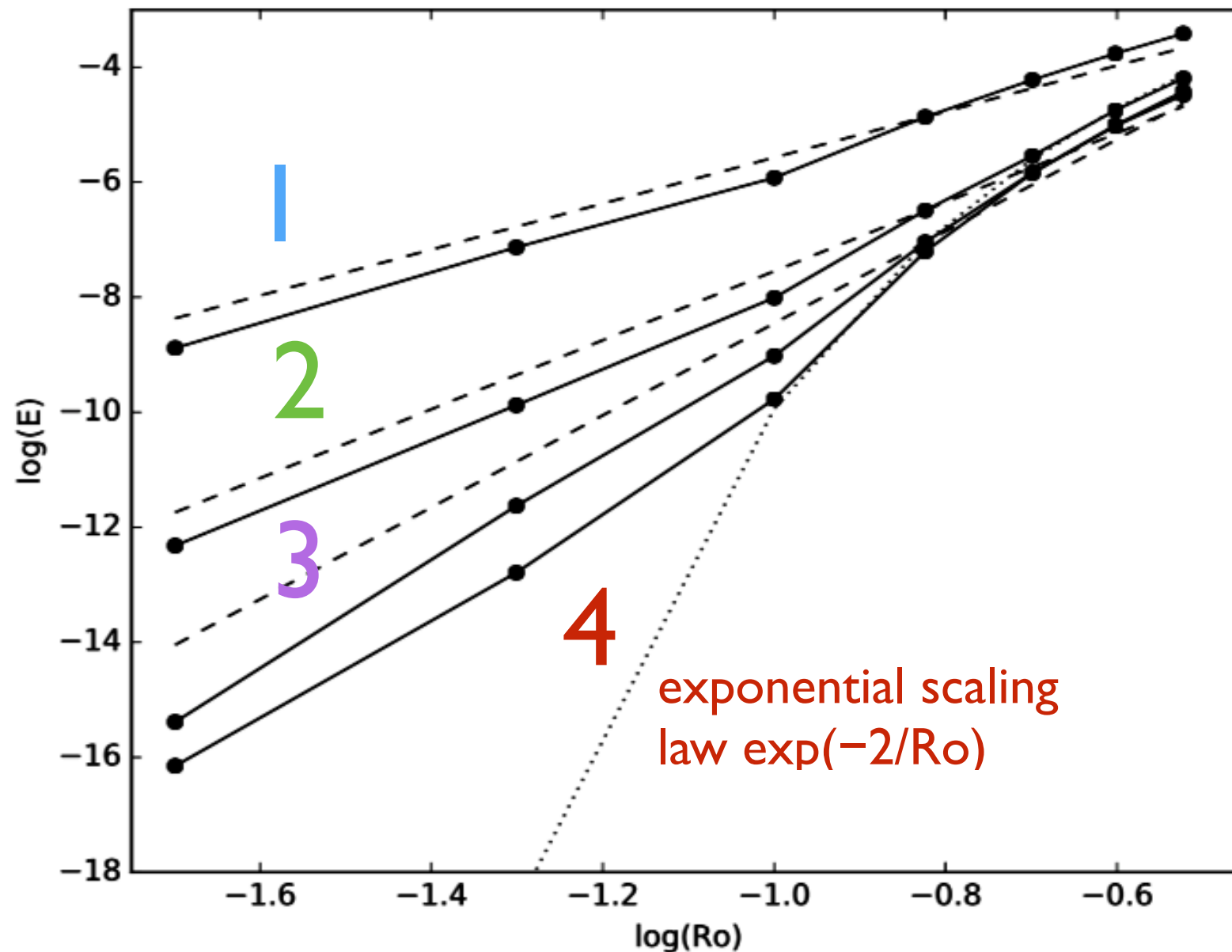


Residual (unbalanced) velocity



Wave emission at higher orders: Ro scaling

Wave generation scales exponentially at higher orders for large Rossby number, Ro .



Total residual wave energy normalized with the total energy integrated over the model domain.

Dashed lines: different power laws

Eden, Chouksey, and Olbers, 2019: Gravity wave emission by shear instability, JPO

Summary

- ❑ The non-linear decomposition of balanced and unbalanced motions is achieved and implemented in model in different dynamical regimes up to fourth order.
 - ❑ Machenhauer (1977) and Warn et al. (1995)
 - ❑ Chouksey et. al (2018) JPO, Eden et. al (2018) JPO, Eden et. al (2019) JPO
- ❑ Balanced state is diagnosed in a single layer and primitive equation model using higher order Ro expansion.
- ❑ The use of C-grid discrete operators is important for the obtained balanced state.
- ❑ **Convective instability generates gravity waves rather than spontaneous emission.**
- ❑ Gravity wave generation scales exponentially at higher orders for large Ro .
- ❑ Spontaneous wave emission by shear instability is negligible.

balanced or unbalanced?!



It's as much numerics, as it's the realm of philosophy.