

Global Ten-Moment Multifluid Simulations of the Solar Wind Interaction with Mercury: From the Planetary Conducting Core to the Dynamic Magnetosphere

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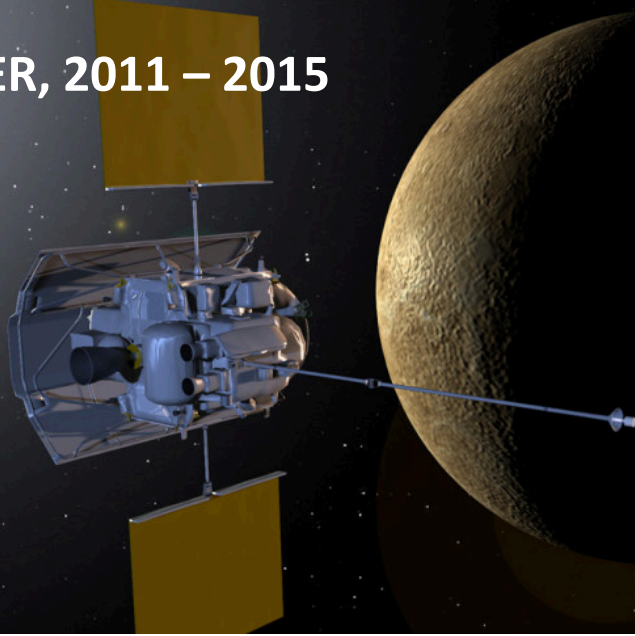
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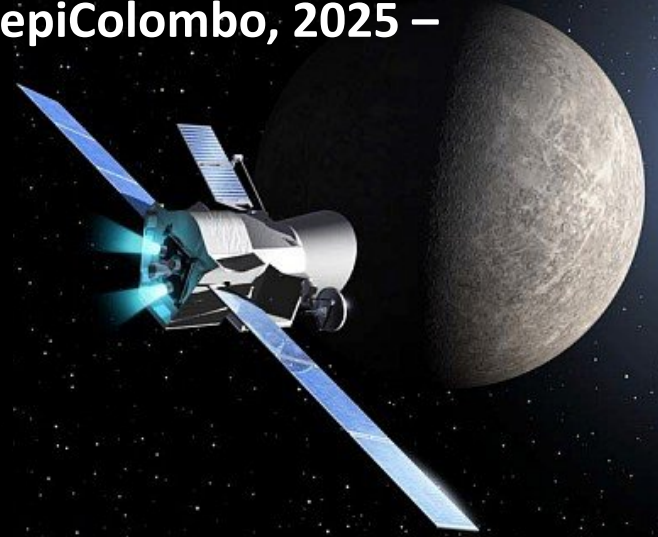
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MESSENGER and BepiColombo at Mercury

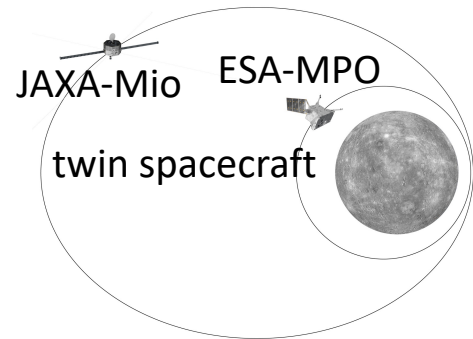
MESSENGER, 2011 – 2015



BepiColombo, 2025 –



BepiColombo is a joint mission of the European Space Agency (ESA) and the Japan Aerospace Exploration Agency (JAXA) to the planet Mercury. It was launched on 20 October 2018 and will arrive at Mercury in late 2025.

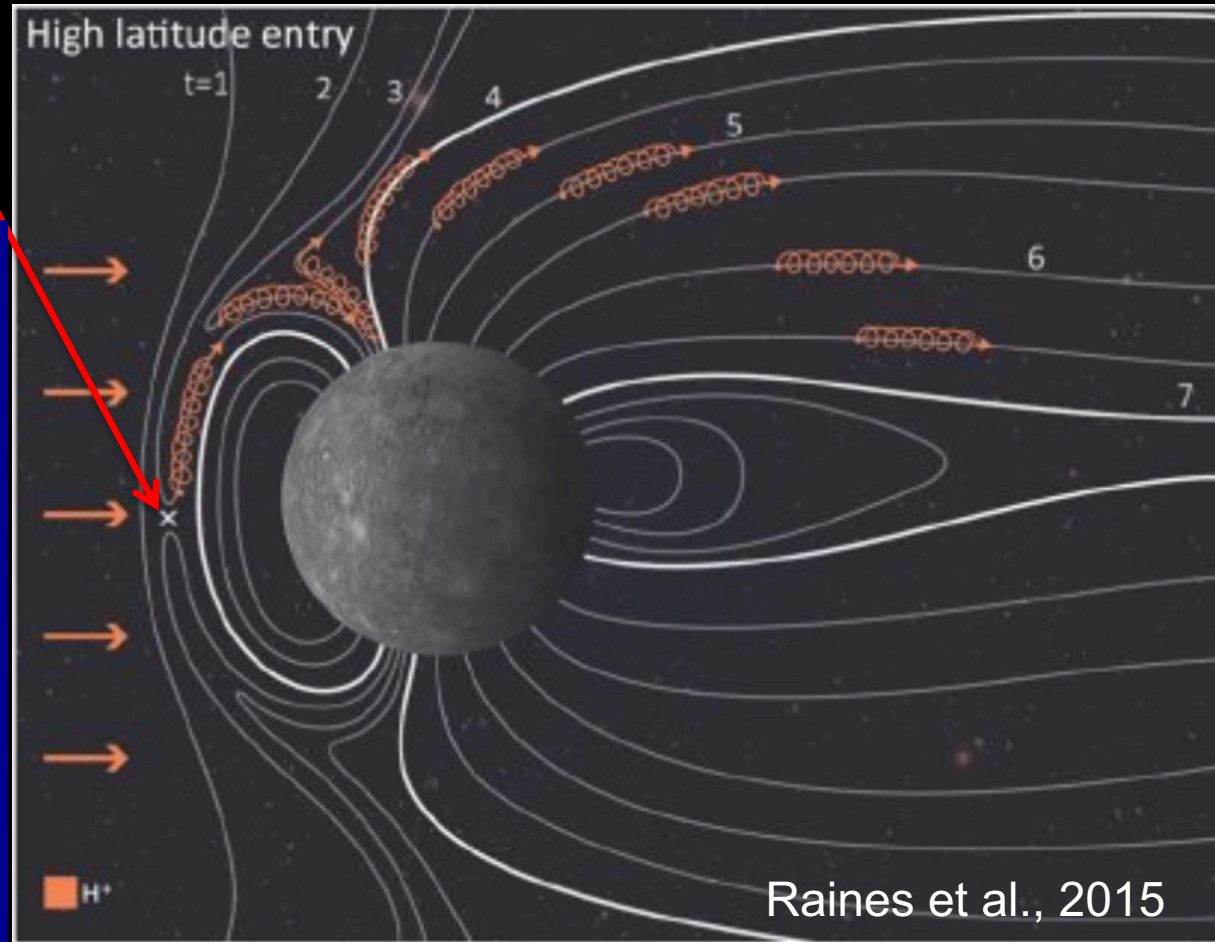


Why Develop a New Model?

Magnetic Reconnection and Solar Wind Plasma Entry in Mercury's Magnetosphere

***Improve kinetic (e.g., collisionless magnetic reconnection) physics in global magnetosphere codes.**

***Note that MHD and hybrid codes cannot capture the collisionless reconnection physics.**



Our Approach to Capture the Collisionless Magnetic Reconnection: Ten-Moment Multifluid Model

- Solve 10-moment equations for *all* electron and ion groups

one density eqn: $\frac{\partial n_s}{\partial t} + \nabla \cdot (n_s \mathbf{v}_s) = 0$

three momentum eqns: $m_s \frac{\partial (n_s \mathbf{v}_s)}{\partial t} + \nabla \cdot \mathcal{P}_s = n_s q_s (\mathbf{E} + \mathbf{v}_s \times \mathbf{B})$

six pressure *tensor* eqns: $\frac{\partial \mathcal{P}_{ij}}{\partial t} + \frac{\partial \mathcal{Q}_{ijk}}{\partial x_k} = n q v_{[i} E_{j]} + \frac{q}{m} \epsilon_{[ikl} \mathcal{P}_{kj]} B_l$

plus *full* Maxwell equations $c^2 \frac{\partial \mathbf{E}}{\partial t} + \mu_0 \mathbf{J} = \nabla \times \mathbf{B}$, $\frac{\partial \mathbf{B}}{\partial t} = -\nabla \times \mathbf{E}$

- *Closure* to approximate heat-flux

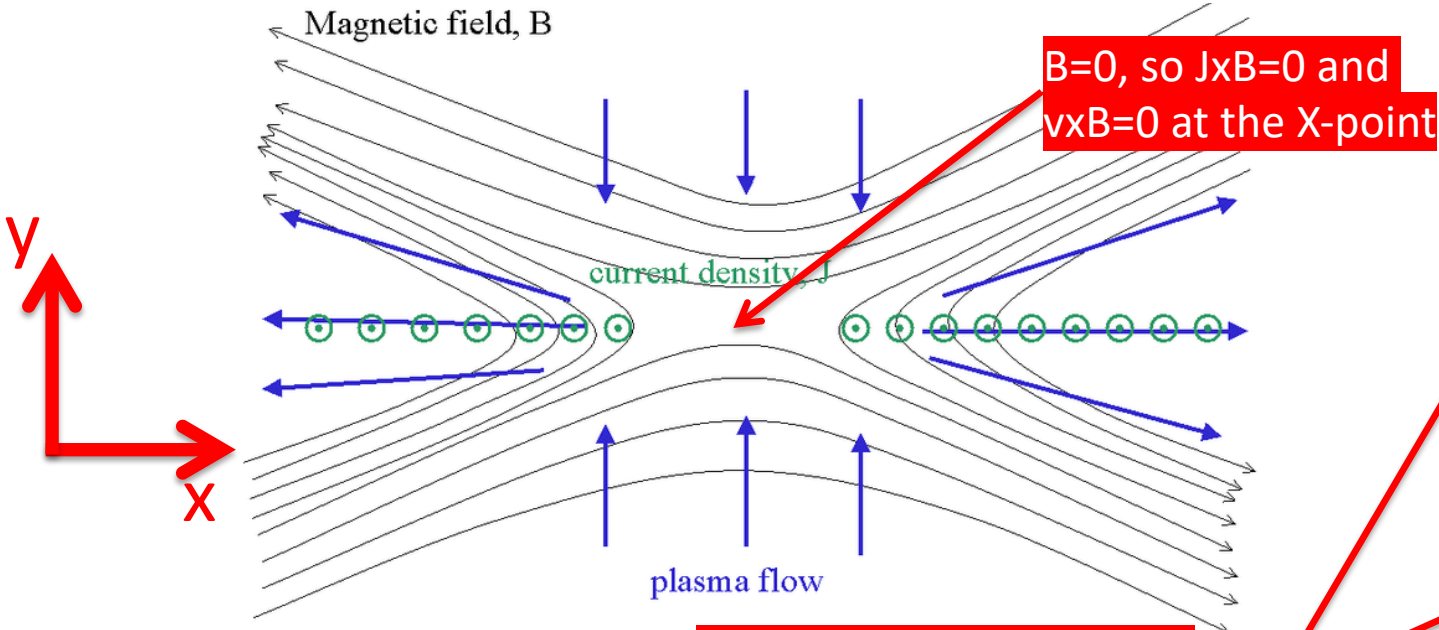
$$\partial_m Q_{ijm} = v_t |k_0| (P_{ij} - p \delta_{ij})$$

- For magnetic reconnection, a reasonable choice is $k_{s0} = 1/d_{s0}$



Why Need Electron Pressure Tensor?

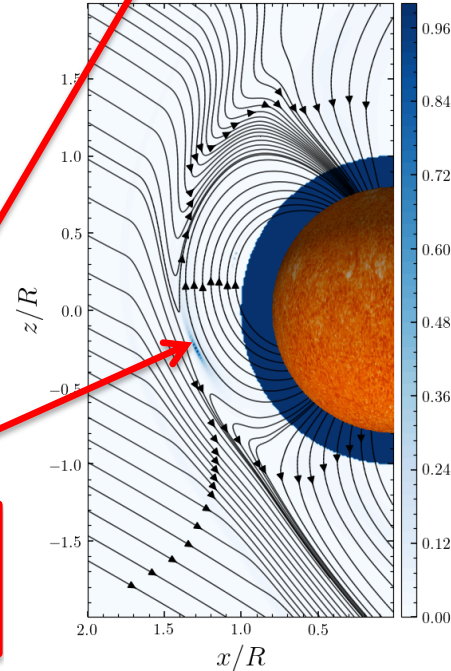
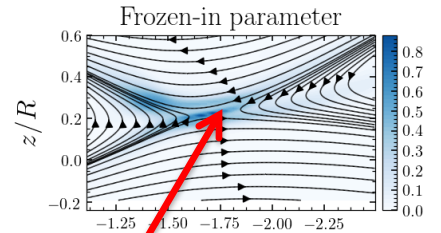
$$E_z = -(\partial_x P_{xz,e} + \partial_y P_{yz,e})/n_e|e|$$



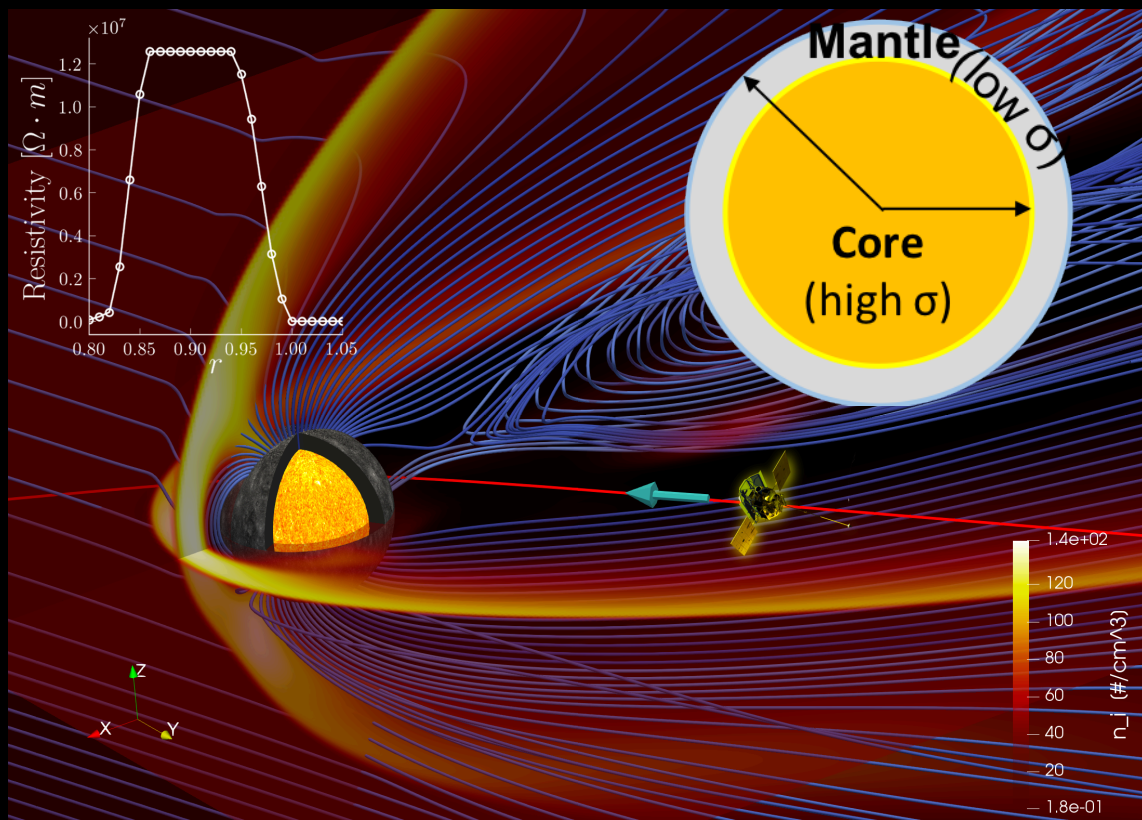
Generalized Ohm's Law

Non-ideal Electric Field

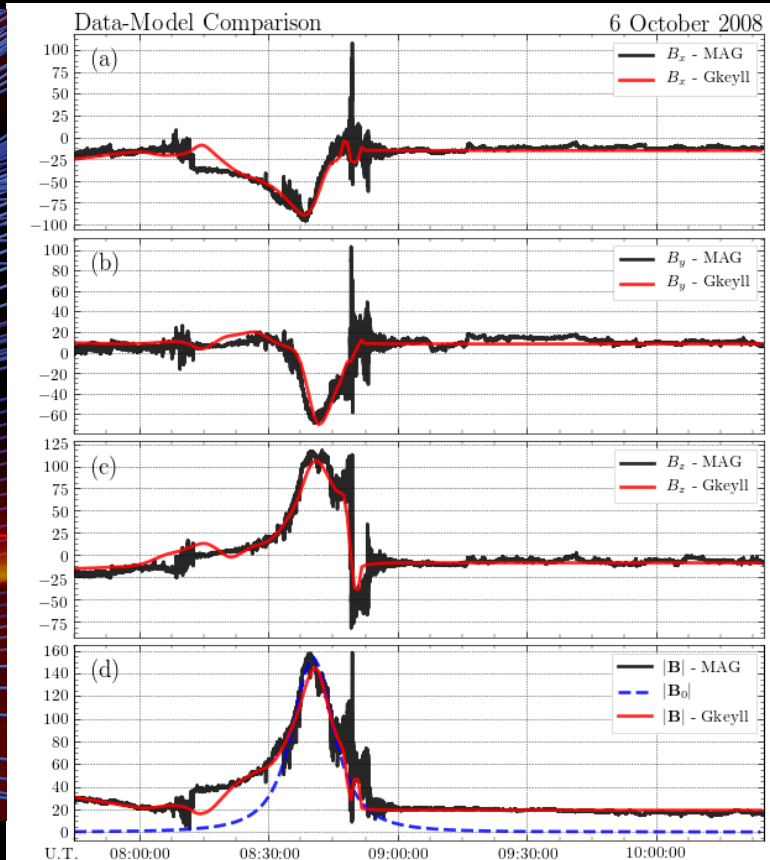
$$\mathbf{E} + \mathbf{v} \times \mathbf{B} = \underbrace{\frac{\eta \mathbf{J}}{0}} + \frac{\mathbf{J} \times \mathbf{B}}{n|e|} - \frac{\nabla \cdot \mathbf{P}_e}{n|e|} + \frac{m_e}{n|e|^2} \left[\frac{\partial \mathbf{J}}{\partial t} + \nabla \cdot \left(\mathbf{v} \mathbf{J} + \mathbf{J} \mathbf{v} - \frac{\mathbf{J} \mathbf{J}}{n|e|} \right) \right]$$



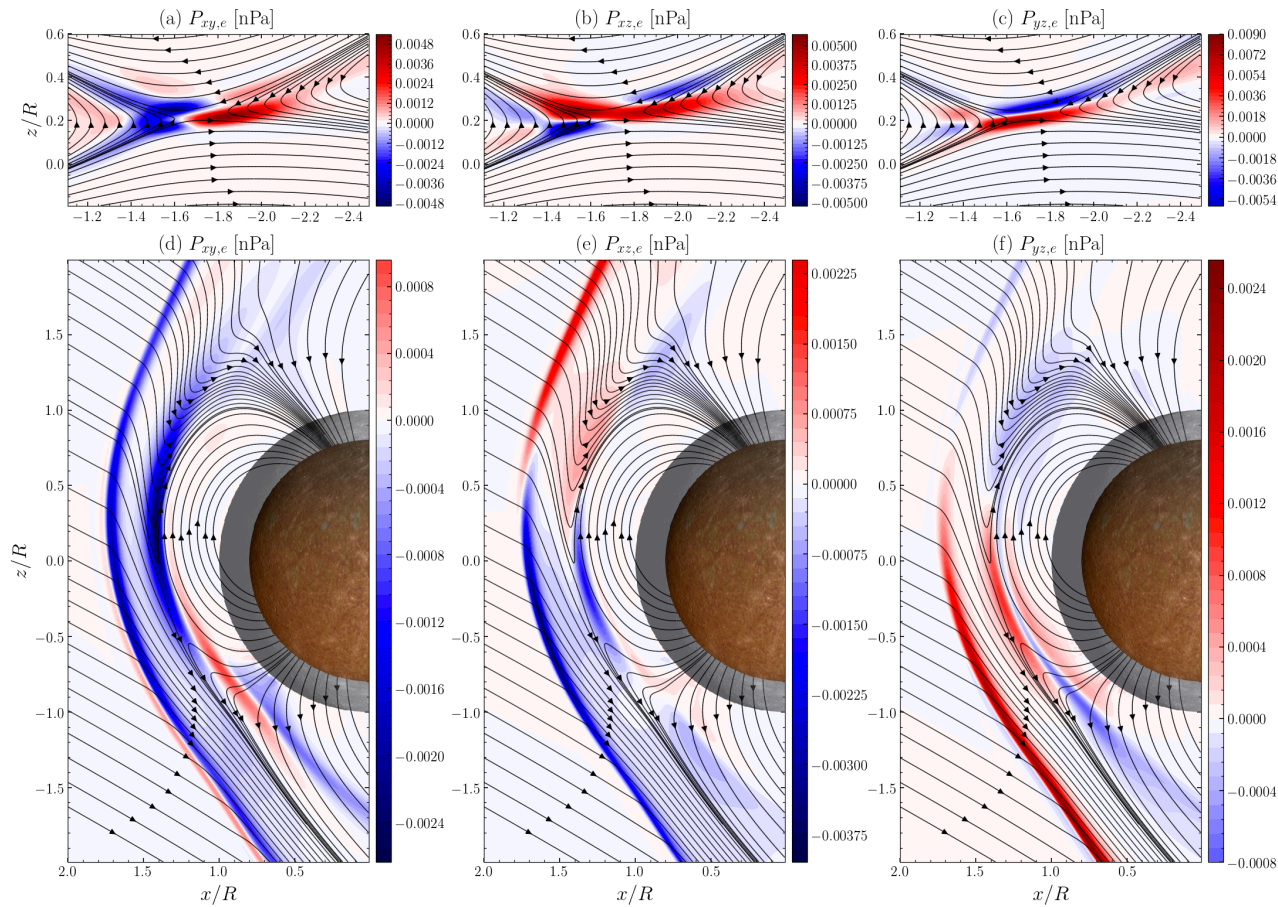
Ten-Moment Multifluid Simulation of Mercury's Magnetosphere



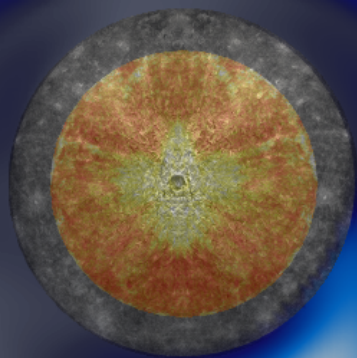
Smallest resolution: $dx=dy=dz=0.01R_M!$



Electron Pressure Tensor at Dayside and Nightside Reconnection Sites



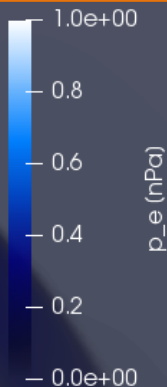
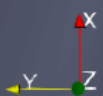
Asymmetry in Mercury's Magnetotail



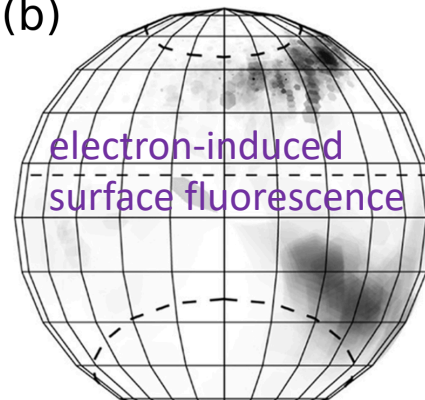
Electron and Ion Motions due to Drifts

Yellow: Electrons

Green: Ions

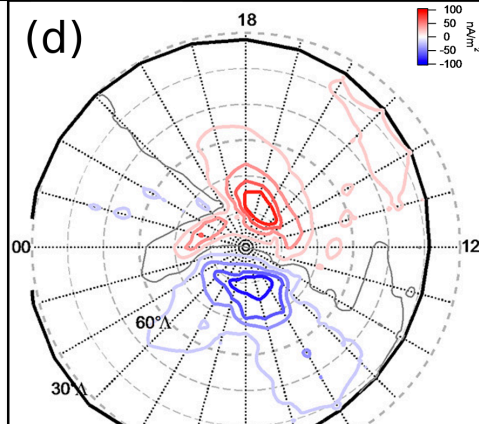


(b)



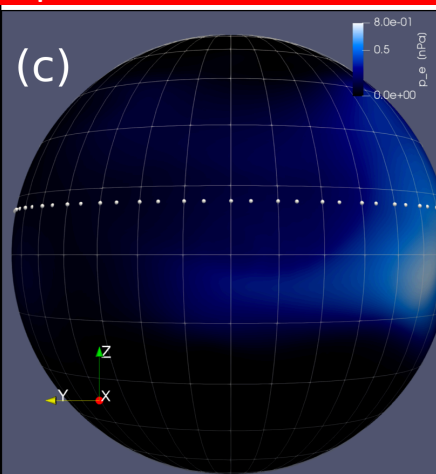
X-Ray Spectrometer Observations

(d)

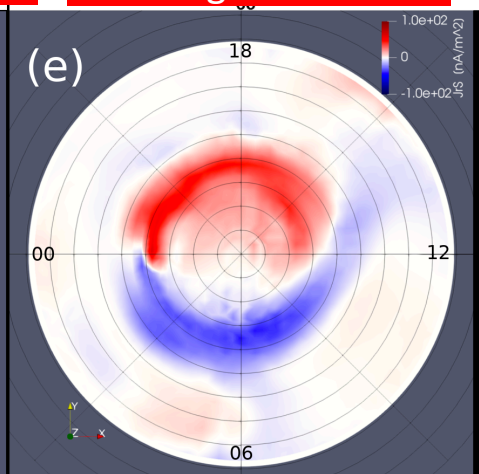


Field-Aligned Currents

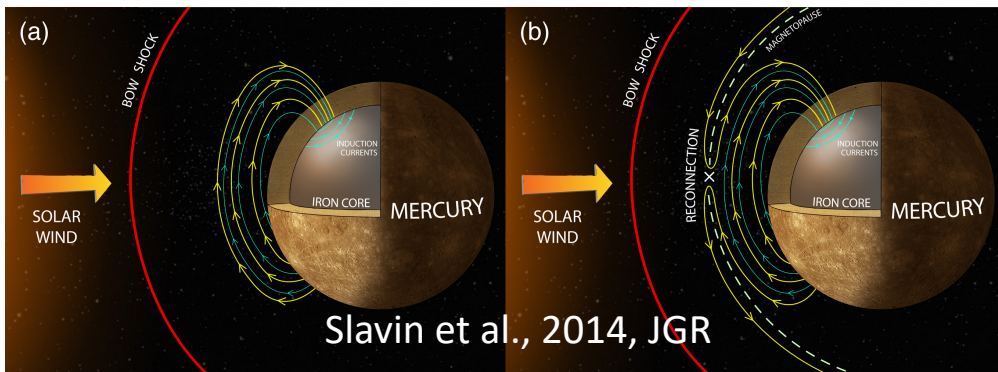
(c)



(e)



Mercury's Responses to an Extreme Event and Plasmoid Formation in Mercury's Magnetotail



The induction response arising from the electromagnetically-coupled interior plays an important role in solar wind-Mercury interaction.

