

Spatio-temporal patterns of fluid-driven aseismic slip transients: implications for seismic swarms

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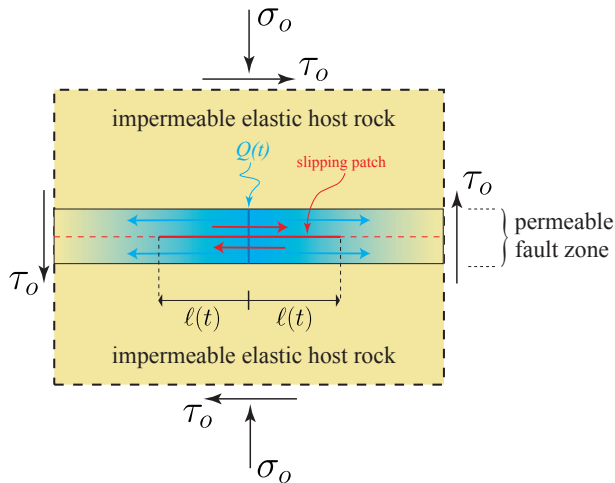
- Common driving mechanisms:
 - Pore pressure diffusion
(e.g., Parotidis *et al.*, *JGR*, 2005)
 - Aseismic slip
(e.g., Lohman & McGuire, *JGR*, 2007)

- Evidence for coupled mechanism:
Pore pressure → Aseismic slip → Seismicity
(e.g., Bourouis & Bernard, *GJI*, 2007; Wei *et al.*, *EPSL*, 2015; Guglielmi *et al.*, *Science*, 2015; Yukutake *et al.*, *JGR*, 2020)

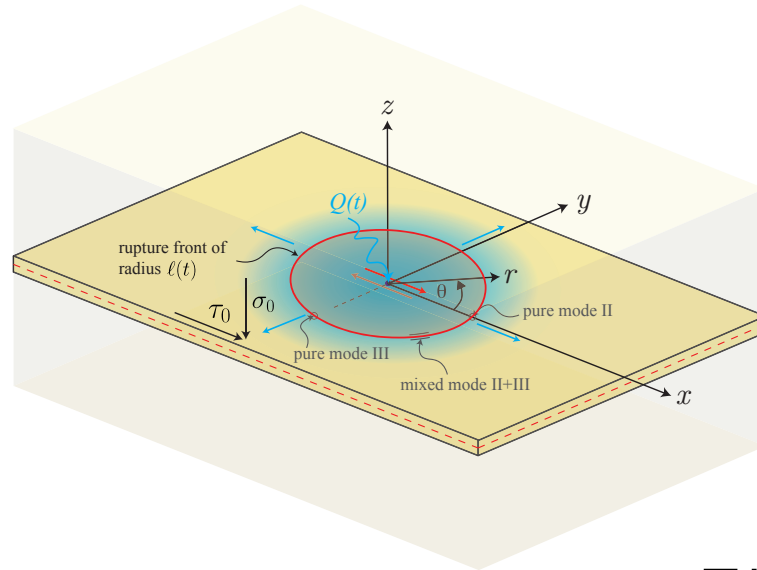
- Spatio-temporal patterns?



- 2D rupture: Mode II or III shear crack



- 3D rupture: Mixed-mode (II+III) circular crack

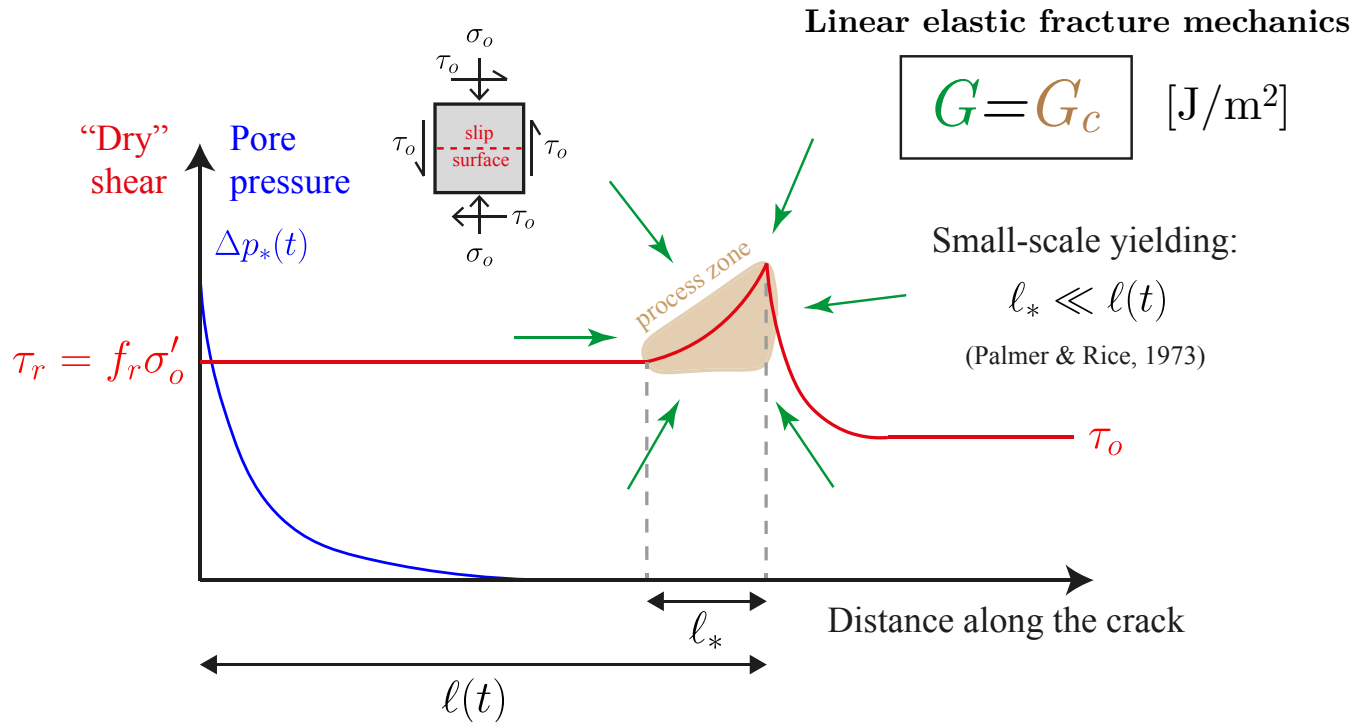


Assumptions:

Constant hydraulic properties, self-similar pore pressure diffusion



Energy balance at the rupture front



Assumptions:

Constant frictional fracture energy G_c , constant residual friction



Ultimate regime

$$G = G_c \Rightarrow$$

Fracture energy term is ultimately **negligible** when:

$$t \rightarrow \infty \quad \text{and} \quad \ell \rightarrow \infty$$

$$\underbrace{f_r \Delta p_*(t) \sqrt{\ell(t)} \int_0^1 \frac{\Pi(\ell(t) \eta / \sqrt{4\alpha t})}{\sqrt{1-\eta^2}} \eta^\gamma d\eta}_{\text{Fluid injection}} = \underbrace{\delta(f_r \sigma'_o - \tau_o) \sqrt{\ell(t)}}_{\text{Pre-stress \& residual friction}} + \underbrace{\frac{\sqrt{\pi}}{2} K_c^*}_{\text{Fracture energy}}$$



Equation for $\ell(t)$:

$$\int_0^1 \frac{\Pi\left(\frac{\ell(t)}{\sqrt{4\alpha t}}\eta\right)}{\sqrt{1-\eta^2}} \eta^\gamma d\eta = \delta \frac{(f_r \sigma'_0 - \tau_0)}{f_r \Delta p_*(t)}$$

$$\Rightarrow F(\ell, t; p_1, \dots) = 0$$

Family of solutions

Table 1: Temporal patterns of $\ell(t)$ for different modes of propagation and types of fluid source.

| Type of fluid source | 2D - Mode II or III | 3D - Mixed mode II+III |
|--------------------------|--|---|
| Constant injection rate | $\ell(t) \propto (t - \sqrt{t})$ at early times (1) $\ell(t) \propto t$ at large times (1) | $\ell(t) \propto \sqrt{t}$ (1) |
| Linearly increasing rate | $\ell(t) \propto (t^{3/2} - 1)/t$ at early times (2) $\ell(t) \propto t^2$ at large times (2) | $\ell(t) \propto \sqrt{t}e^{-t}$ at early times (2) $\ell(t) \propto t$ at large times (2) |
| Constant pressure | $\ell(t) \propto \sqrt{t}$ (3) | $\ell(t) \propto t^\beta$ (1)(4) |

Notes: Derived in (1) Sáez, Lecampion, Bhattacharya & Viesca, *JMPS*, 2022.

(2) Sáez & Lecampion (in preparation). (3) Bhattacharya & Viesca, *Science*, 2019.

(4) β between 0 and 1/2.

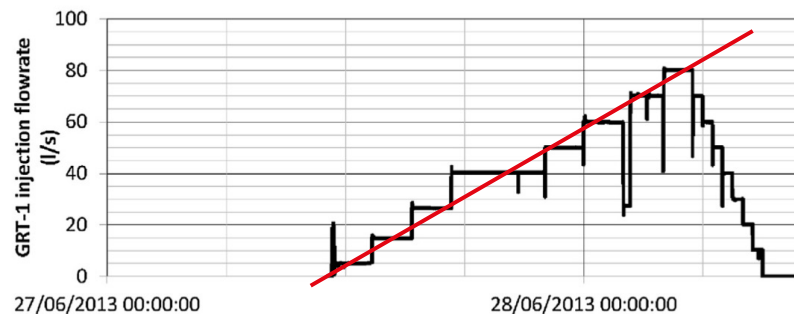


Example of application – Borehole fluid injection

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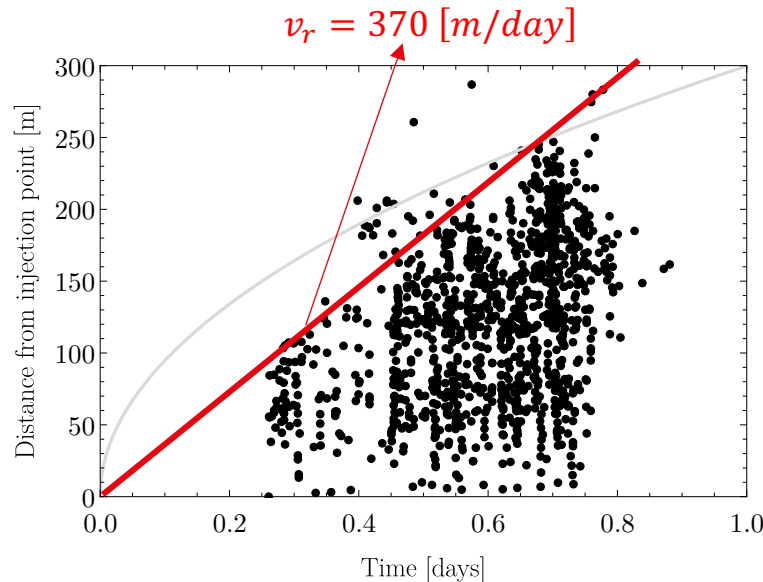
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Baujard et al., *Geothermics*, 2017


$$\ell(t) = v_r t, \text{ with } v_r = \sqrt{\frac{f_r \alpha \dot{Q} \eta / 4 \pi k w}{f_r \sigma'_0 - \tau_0}}$$

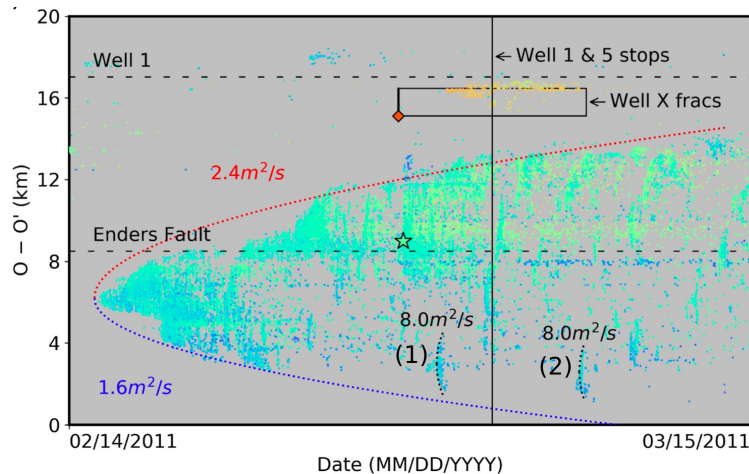


Data from Lengliné et al., *GJI*, 2017



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$$\ell(t) = \sqrt{\alpha_{slip} t}$$

$$\alpha_{slip} = \lambda^2 \alpha_h$$

$$\text{With } \lambda = \sqrt{\frac{1}{8\pi} \frac{f_r Q \eta / k w}{f_r \sigma'_0 - \tau_0}} \gg 1$$

(for critically-stressed faults)

Reconciling high diffusivities!



- Family of **physics-based solutions** for fluid-driven **aseismic slip** fronts $\ell(t)$
- Many **different** spatio-temporal **patterns** may occur depending on **fluid source** and **mode of sliding**, among others.
- For instance,
 - We used a seismic swarm of **known fluid source** to show under which **conditions** a swarm may be expected to migrate **linearly with time** if dominated by aseismic slip.
 - We speculate that some **diffusive seismic swarms** (\sqrt{t}) may be driven by **diffusive aseismic slip** in 3D, which **reconciles** relatively **high diffusivities** inferred in many cases ($\alpha_{slip} = \lambda^2 \alpha_h$, with $\lambda \gg 1$).
- Finally, we expect our solutions to provide a **simple means to interpret observations** of seismic swarms, in the framework of coupled fluid flow and aseismic slip processes.



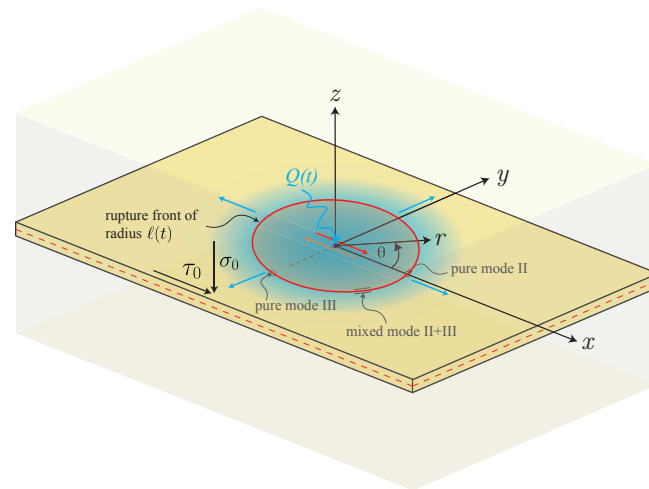
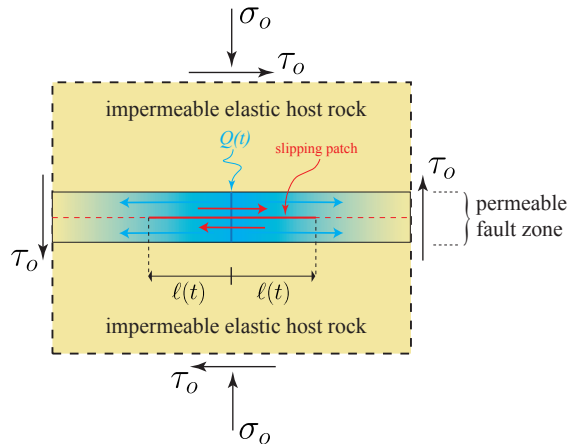


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