

# Alfvénicity of Velocity and Magnetic Field Increments Observed by Parker Solar Probe

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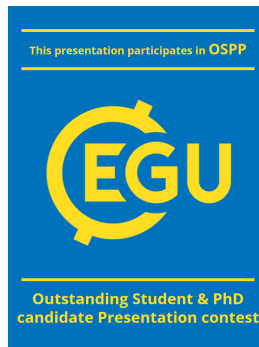
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Abstract  
QR code



## Introduction

- Alfvénicity is the correlation of velocity fluctuations and magnetic fluctuations

$$\boldsymbol{v} = \pm \boldsymbol{b}$$

where magnetic fluctuation  $\boldsymbol{b}$  is measured in Alfvén speed units

- We want to study the physics of Alfvénicity and how Alfvénicity depends on other variables using data from Parker Solar Probe (PSP) when the solar wind was super-Alfvenic or sub-Alfvenic.

# Alfvénicity of increments

Magnetic and velocity fluctuations are defined as follows with  $\tau$  as a time increment

$$\Delta \mathbf{v} = \mathbf{v}(t + \tau) - \mathbf{v}(t)$$

$$\Delta \mathbf{b} = \mathbf{b}(t + \tau) - \mathbf{b}(t)$$

Measures to quantify Alfvénicity (Parashar et al. 2020)

- **Normalized cross helicity ( $\sigma_c$ )**

$$\sigma_c = \frac{2\langle \Delta \mathbf{v} \cdot \Delta \mathbf{b} \rangle}{\langle |\Delta \mathbf{v}|^2 \rangle + \langle |\Delta \mathbf{b}|^2 \rangle}$$

- **Normalized residual energy ( $\sigma_r$ )**

$$\sigma_r = \frac{\langle |\Delta \mathbf{v}|^2 \rangle - \langle |\Delta \mathbf{b}|^2 \rangle}{\langle |\Delta \mathbf{v}|^2 \rangle + \langle |\Delta \mathbf{b}|^2 \rangle}$$

- **Alfvén ratio ( $r_A$ )**

$$r_A = \frac{\langle |\Delta \mathbf{v}|^2 \rangle}{\langle |\Delta \mathbf{b}|^2 \rangle}$$

- **Cosine of the alignment angle between  $\mathbf{v}$  and  $\mathbf{b}$  ( $\cos \theta_{vb}$ )**

$$\cos \theta_{vb} = \frac{\langle \Delta \mathbf{v} \cdot \Delta \mathbf{b} \rangle}{\sqrt{\langle |\Delta \mathbf{v}|^2 \rangle + \langle |\Delta \mathbf{b}|^2 \rangle}}$$

## Frequency in Fourier spectra

To match the frequency that contributes the most to  $\Delta \mathbf{b}$  or  $\Delta \mathbf{v}$ , we assume that  $\mathbf{b}$  and  $\mathbf{v}$  have Kolmogorov spectra and assume the continuous equation for  $\mathbf{b}$  (or  $\mathbf{v}$ )

$$\mathbf{b}(t) = \int \mathbf{b}(\omega) e^{i\omega(t)} d\omega$$

Let  $\Delta \mathbf{b} = \mathbf{b}(t + \tau) - \mathbf{b}(t)$  and

$$G(\omega_0) = \frac{8\pi}{T} \int_0^{\omega_0} \mathbf{b}(\omega) \mathbf{b}(-\omega) (1 - e^{-i\omega\tau}) d\omega \quad \text{for very large } T$$

$$\langle \Delta \mathbf{b}^2 \rangle = G(\infty)$$

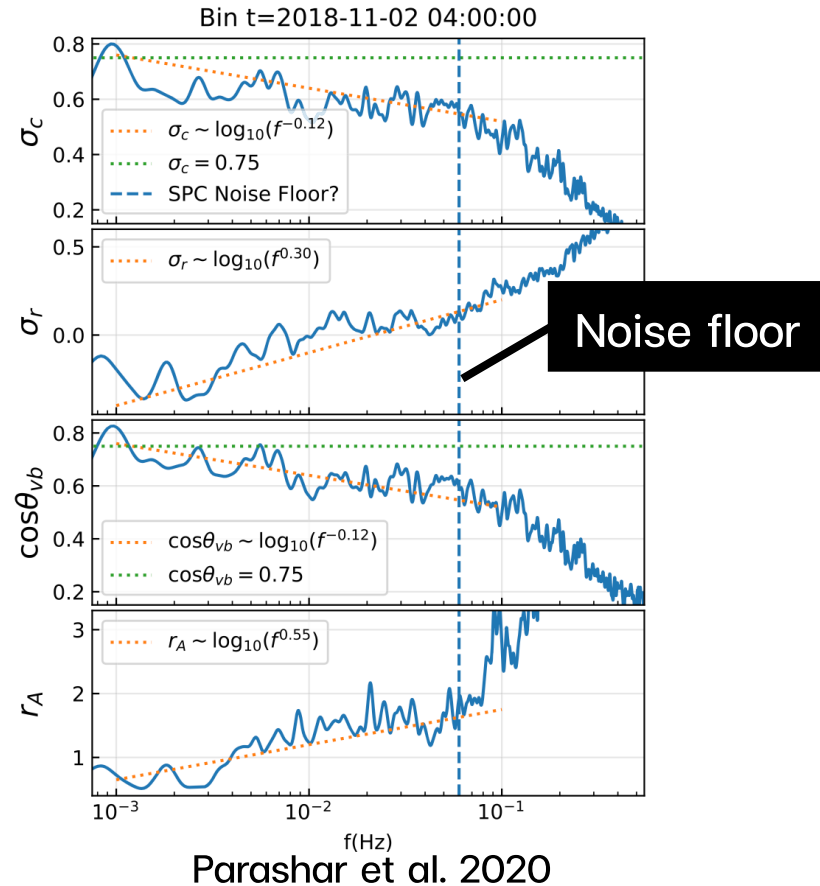
Consider the contribution to  $\langle \Delta \mathbf{b}^2 \rangle$  per logarithm of  $\omega$ :

$$\frac{dG}{d \ln \omega} = \omega \frac{dG}{d\omega} \propto G(2\omega_0) - G(\omega_0/2) \sim \mathbf{b}(\omega) \mathbf{b}(-\omega) (1 - e^{-i\omega\tau}) \omega$$

This contribution is maximized for

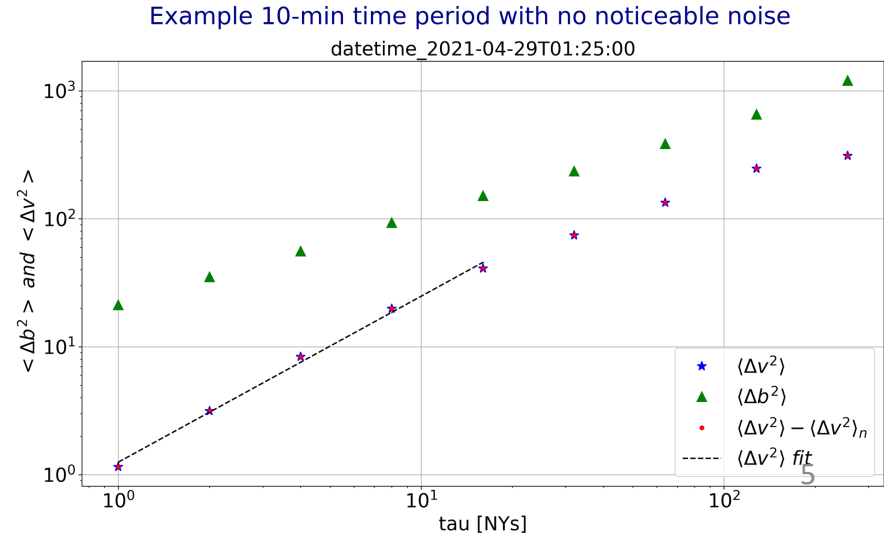
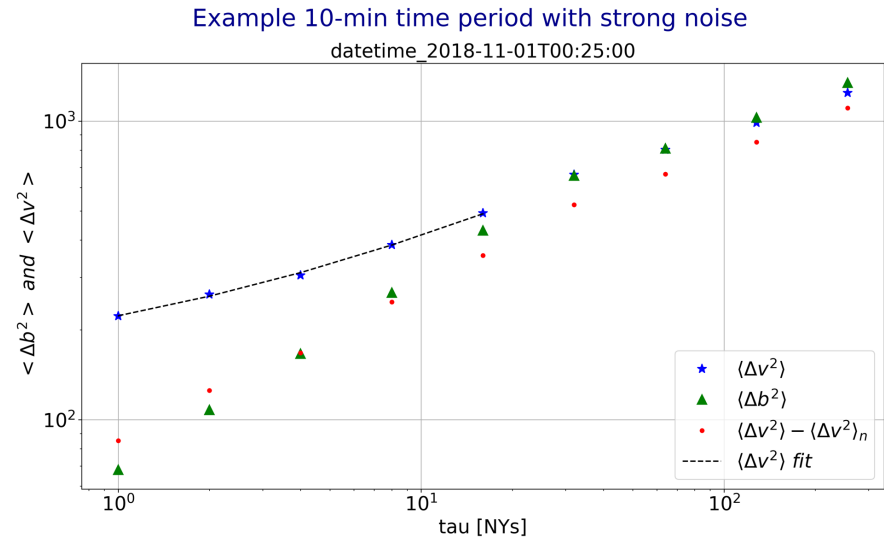
$$\omega \approx \frac{2.65}{\tau}$$

# Noise subtraction

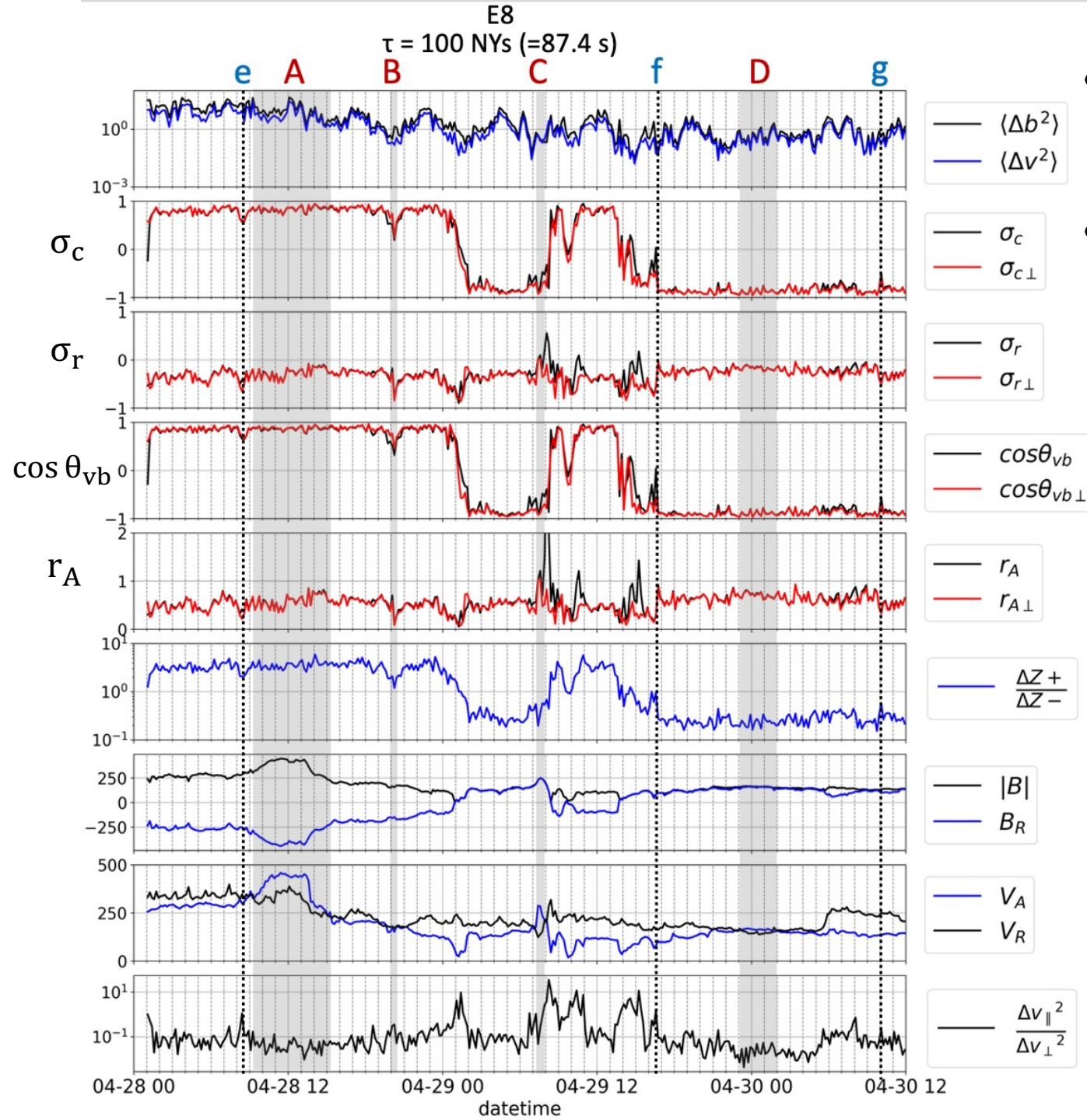


In previous work, we can observe noise at high frequencies. Without the noise, we can estimate Alfvénicity at small scales or high frequencies.

- We can eliminate the noise by assuming that the power spectrum of the  $\langle |\Delta v|^2 \rangle$  is power-law in the inertial range without noise.



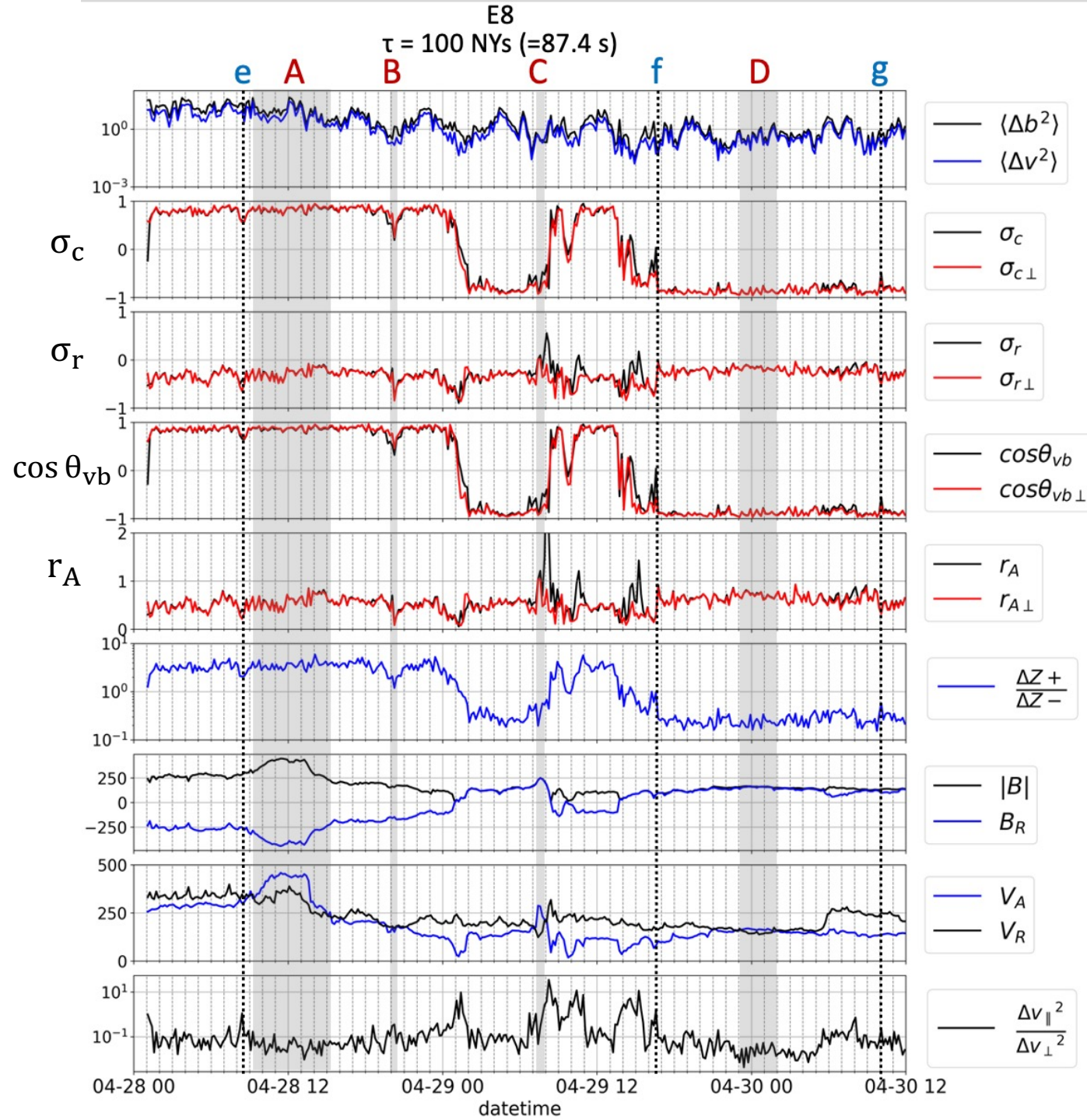
# Time dependence of Alfvénicity near 8<sup>th</sup> perihelion (10-min averages)



- Regions **A**, **B**, **C** and **D** are sub-Alfvénic.
- Regions **A** and **D** have typical high Alfvénicity for  $\tau \sim 14.7 - 87.4$  s, but the region **B** has low Alfvénicity with low velocity fluctuation as the spacecraft possibly passes a twisted magnetic field inside a flux rope.



# Time dependence of Alfvénicity near 8<sup>th</sup> perihelion (10-min averages)



- Region **C** has "ideal Alfvénicity" with  $r_A \approx 1$  and  $\sigma_r \approx 0$  for balanced magnetic and velocity fluctuations for  $\tau \sim 14.7 - 87.4$  s, unlike typical solar wind.
- e**, **f** and **g** observed sudden drop of Alfvénicity at the high ratio of parallel to perpendicular velocity increments.