



Unstable Bickley Jets in 2D-MHD

EGU 2022

Jonathan Tessier¹, Francis J. Poulin¹, David W. Hughes²

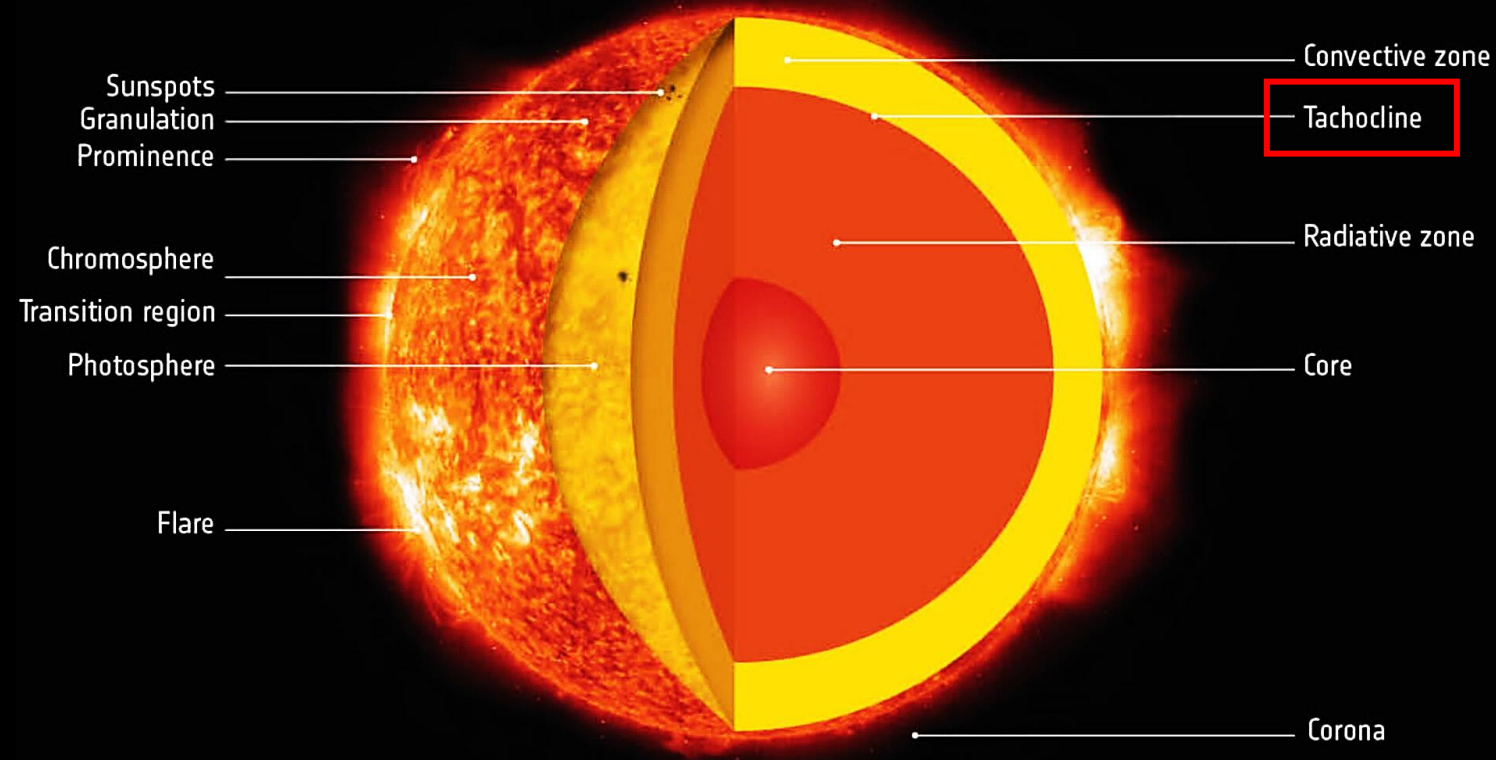
¹University of Waterloo, Applied Mathematics, Waterloo, Canada

²University of Leeds, School of Mathematics, Leeds, England



Motivation: The Solar Tachocline

- Transition from radiative zone to convective zone
- Very thin ~ 2D motion
- High shear & Jets



Model: 2D MHD

Parameter Regime:

$$R_e = R_m = 10^4$$

$$M = 0 - 0.1$$

$$\partial_t q + \mathbf{u} \cdot \nabla q = M^2 \mathbf{b} \cdot \nabla j + R_e^{-1} \nabla^2 q,$$

$$\partial_t A + \mathbf{u} \cdot \nabla A = R_m^{-1} \nabla^2 A,$$

$$\mathbf{u} = \hat{z} \times \nabla \psi,$$

$$\mathbf{b} = \hat{z} \times \nabla A,$$

$$q = \nabla^2 \psi,$$

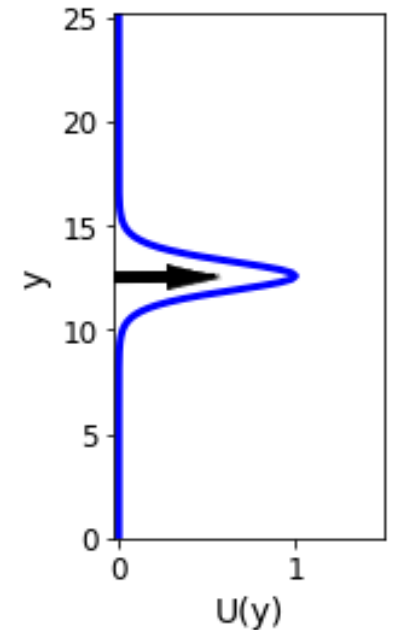
$$j = \nabla^2 A.$$

$$M = \frac{B_0 / \sqrt{\mu_0 \rho}}{U_0}$$

Background state:

$$U(y) = \text{sech}^2(y)$$

$$B(y) = 1$$

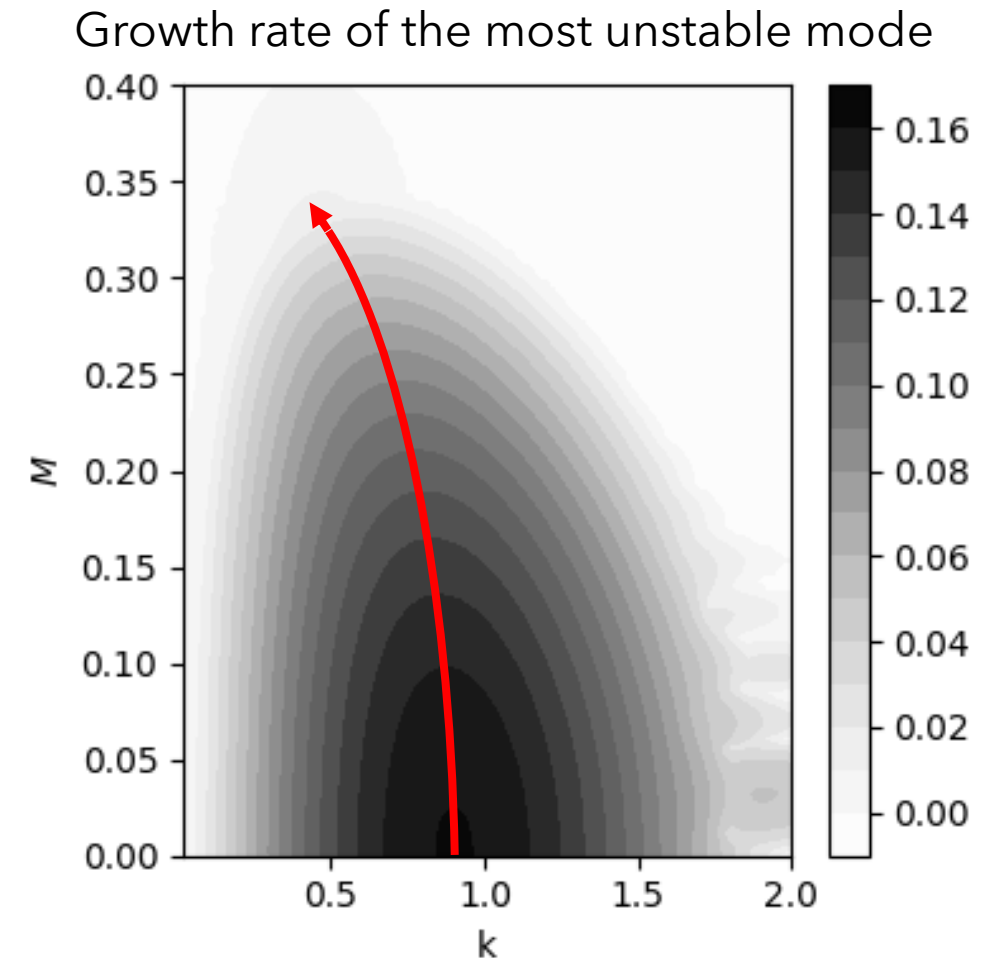


Linear Theory: The Bickley Jet in 2D MHD

Normal decomposition

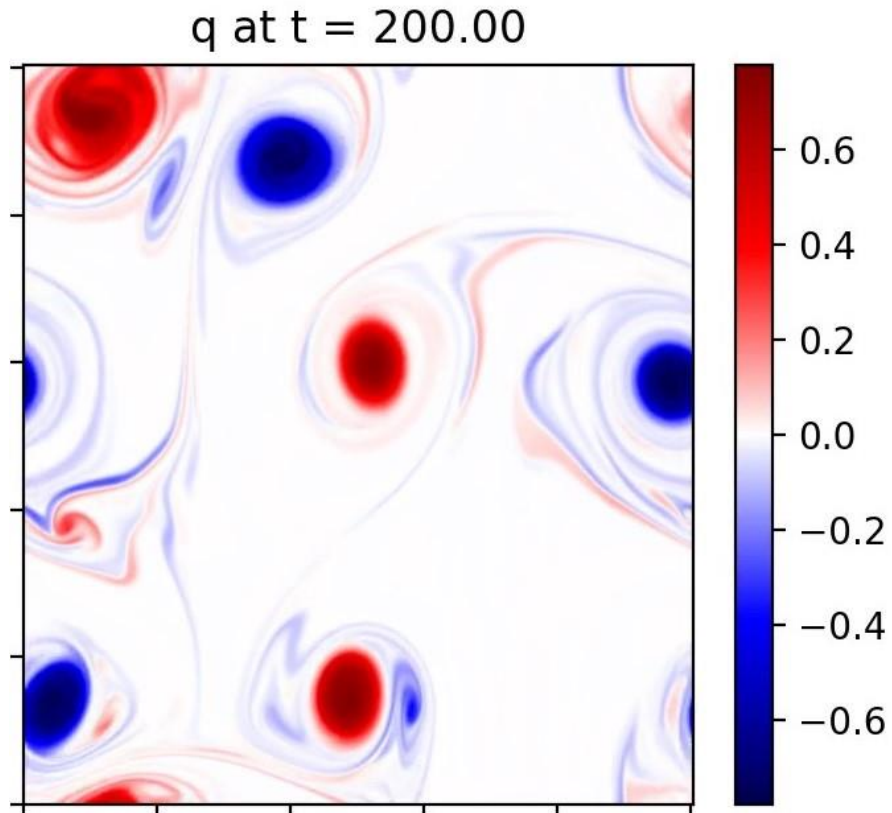
$$[q', A'] = [\hat{q}(y), \hat{A}(y)] \exp(i(kx - \omega t))$$

$$\|q'\| \propto \exp(\Im(\omega)t)$$

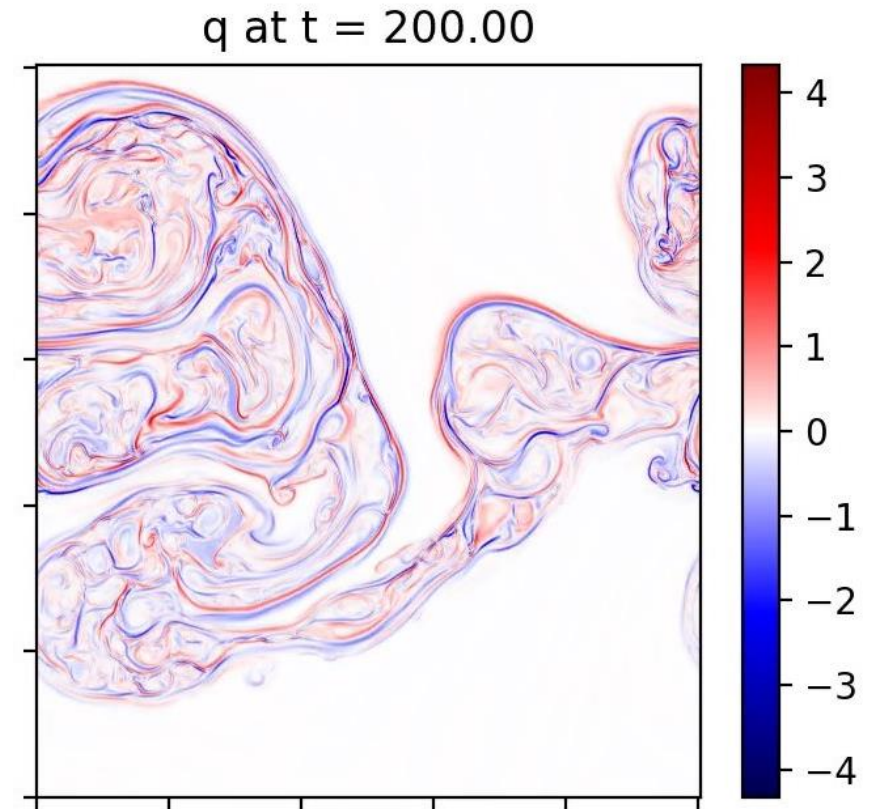


Nonlinear Evolution: The Bickley Jet

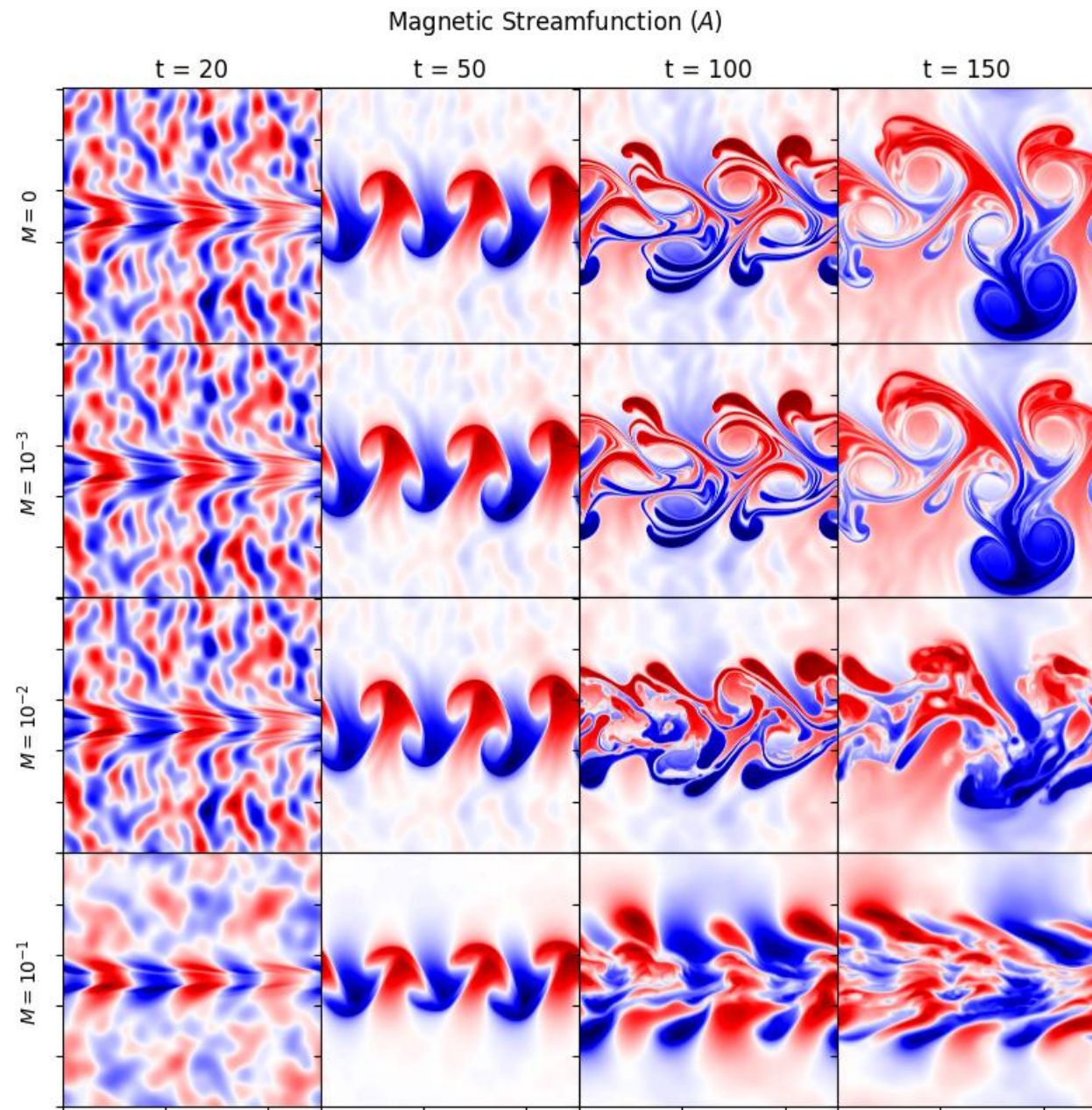
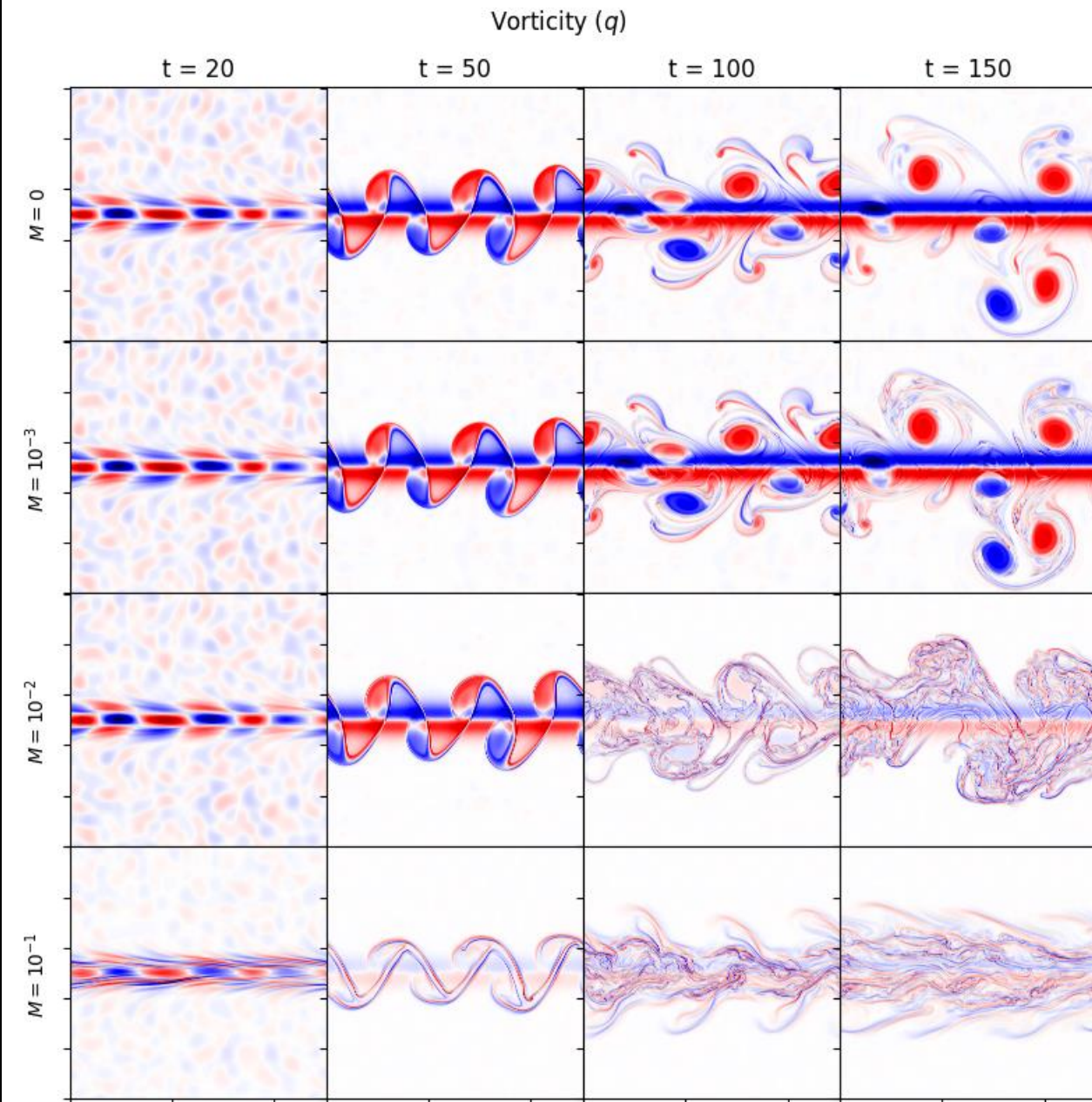
$M = 0$ (Hydrodynamic)



$M = 0.01$ (MHD, $M^2 R_m = 1$)



More cases



Summary

Increasing M :

- Stabilizes the Bickley jet in a uniform parallel field.
- Moves the largest growing mode to larger length scales (in x).
- Generates small scale filaments of vorticity.

References

- Gilman, P. A. (2000). Magnetohydrodynamic shallow water equations for the solar tachocline. APJ 544, 1, L79-L82.
- Hughes, D. W., Rosner, R., and Weiss, N. (2007). The Solar Tachocline. Cambridge University Press.
- Mak, J. & Griffiths, S.D. & Hughes, D.W. (2016). Shear flow instabilities in shallow-water magnetohydrodynamics. JFM, vol. 788, pp. 767-796
- Mak, J. & Griffiths, S.D. & Hughes, D.W. (2017). Vortex disruption by magnetohydrodynamic feedback. PRF. 2. 10.1103/PhysRevFluids.2.113701