

Electron Injection via Stochastic Shock Drift Acceleration: Theory, Simulation, and Observation

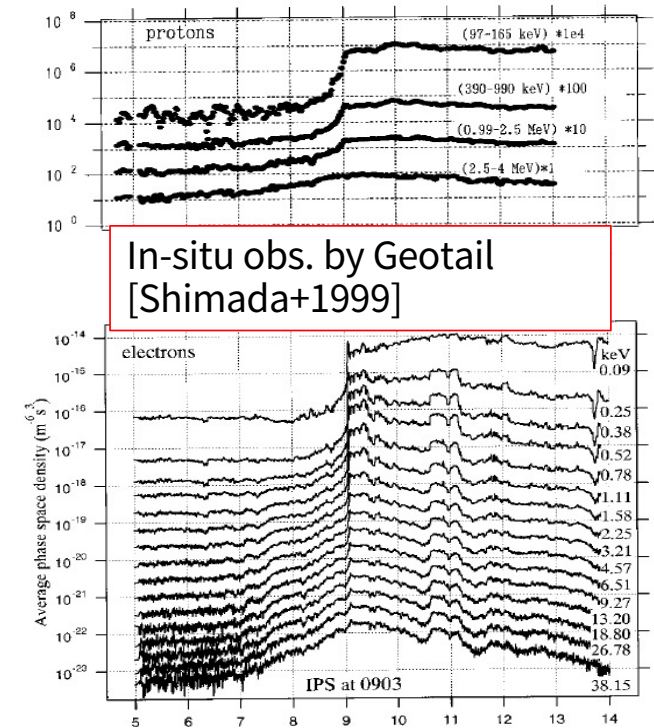
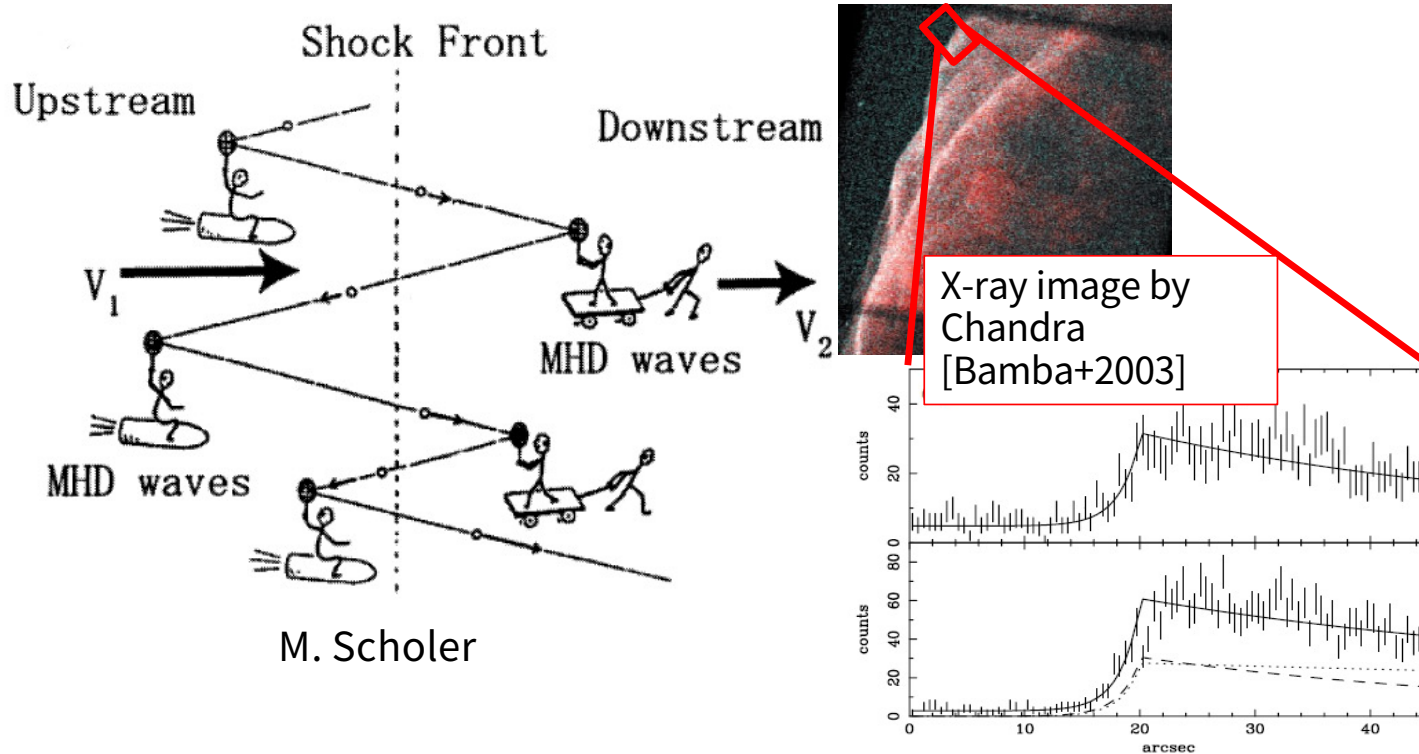
Takanobu Amano (U-Tokyo)

Collaborators

T. Katou, M. Hoshino (U-Tokyo), Y. Matsumoto (Chiba U), M. Oka (UCB), S. Matsukiyo (Kyusyu U) O. Kobzar (Jagiellonian U), J. Niemiec (IFJ/PAN), A. Bohdan M. Pohl (DESY)

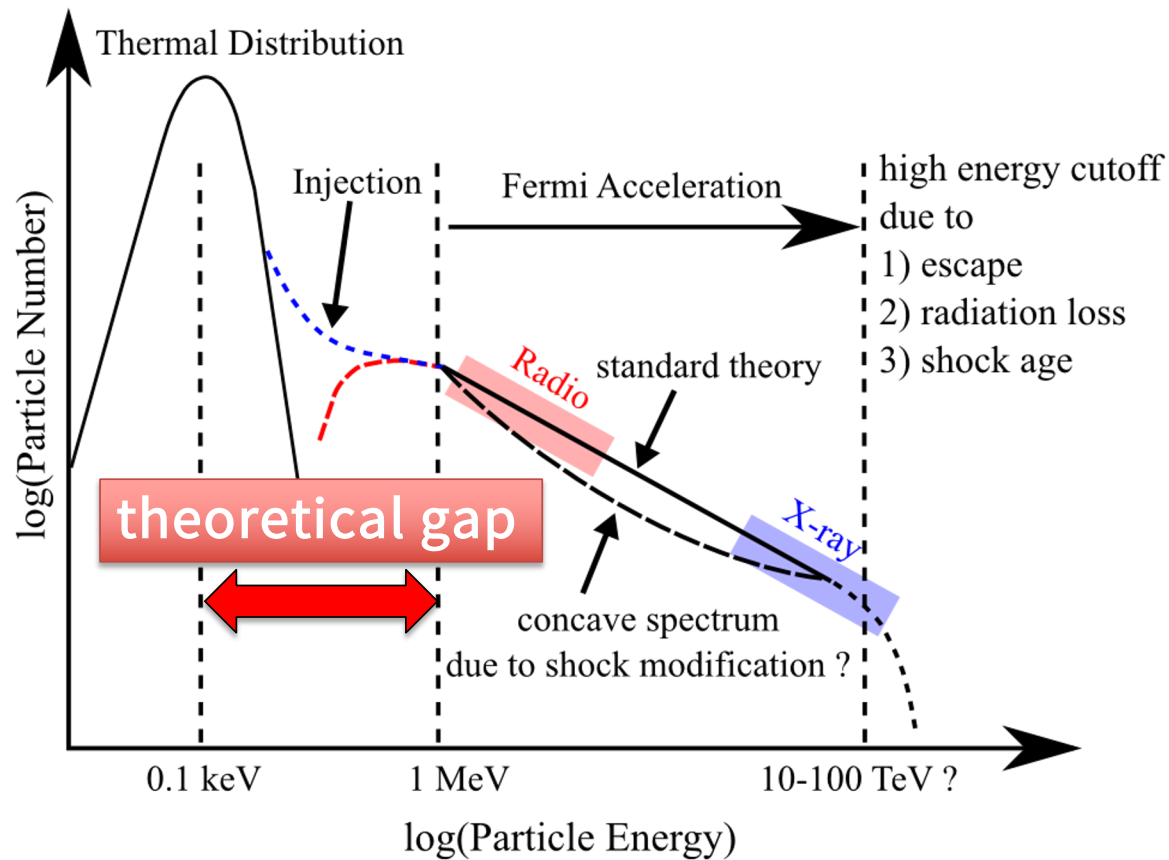
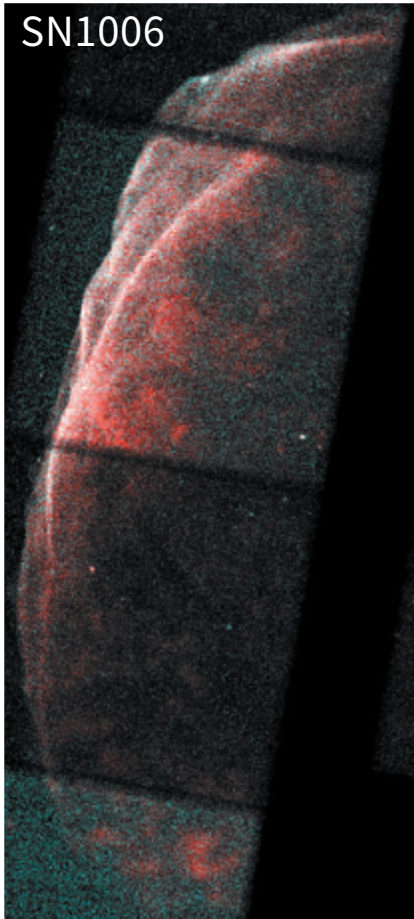
The Standard Shock Acceleration Theory

- Diffusive Shock Acceleration (DSA)
 - A simple yet powerful model that predicts a nearly universal power-law; $N(E) \propto E^{-2}$
[e.g., Bell 1978, Blandford & Ostriker 1978]



The Electron Injection

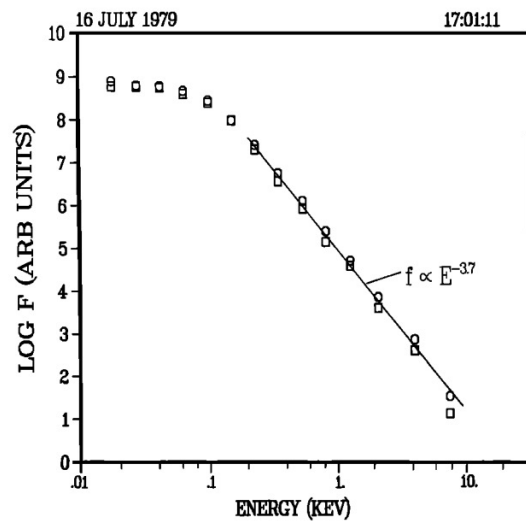
electrons with $< 0.1-1$ MeV cannot be scattered by MHD waves $\omega - kv_{\parallel} = \Omega/\gamma$



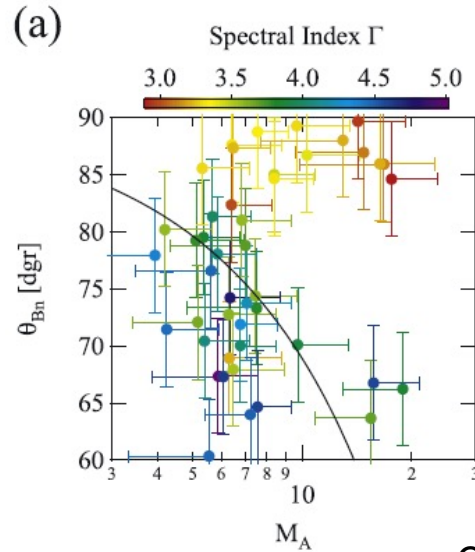
- ✓ Sub-relativistic electrons cannot be accelerated by the standard first-order Fermi mechanism.
- ✓ Substantial energy gain is needed from thermal to relativistic energies by some other mechanisms.
- ✓ Sub-relativistic suprathermal electrons are “invisible” with typical astrophysical observations, while they are observable with in-situ spacecraft measurement.

Ref: Levinson 1992, Amano & Hoshino 2010, Shimada+1999, Lario+2003, Dresing+2016

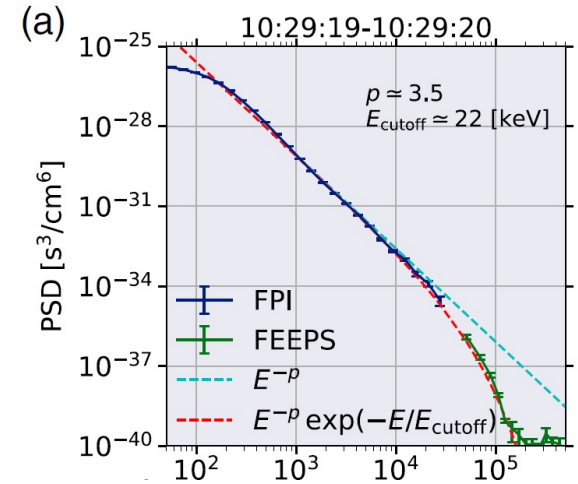
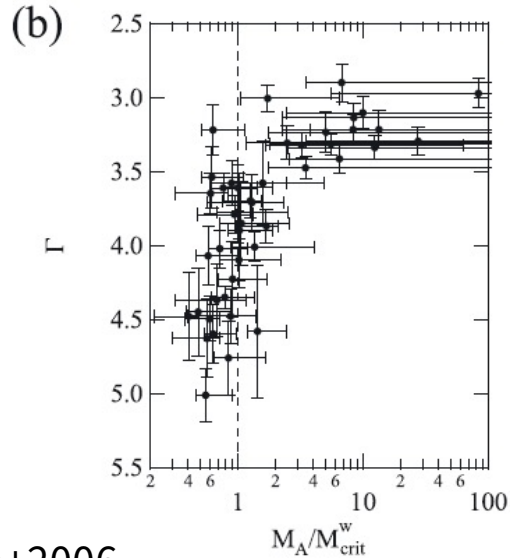
Earth's Bow Shock: Laboratory for Electron Injection



Gosling+1989



Oka+2006



Amano+2020

- Non-thermal electrons with a clear power-law spectrum have been observed occasionally at the bow shock.
- The typical energy range of non-thermal electrons measured at the bow shock is the most important energy range for the injection.

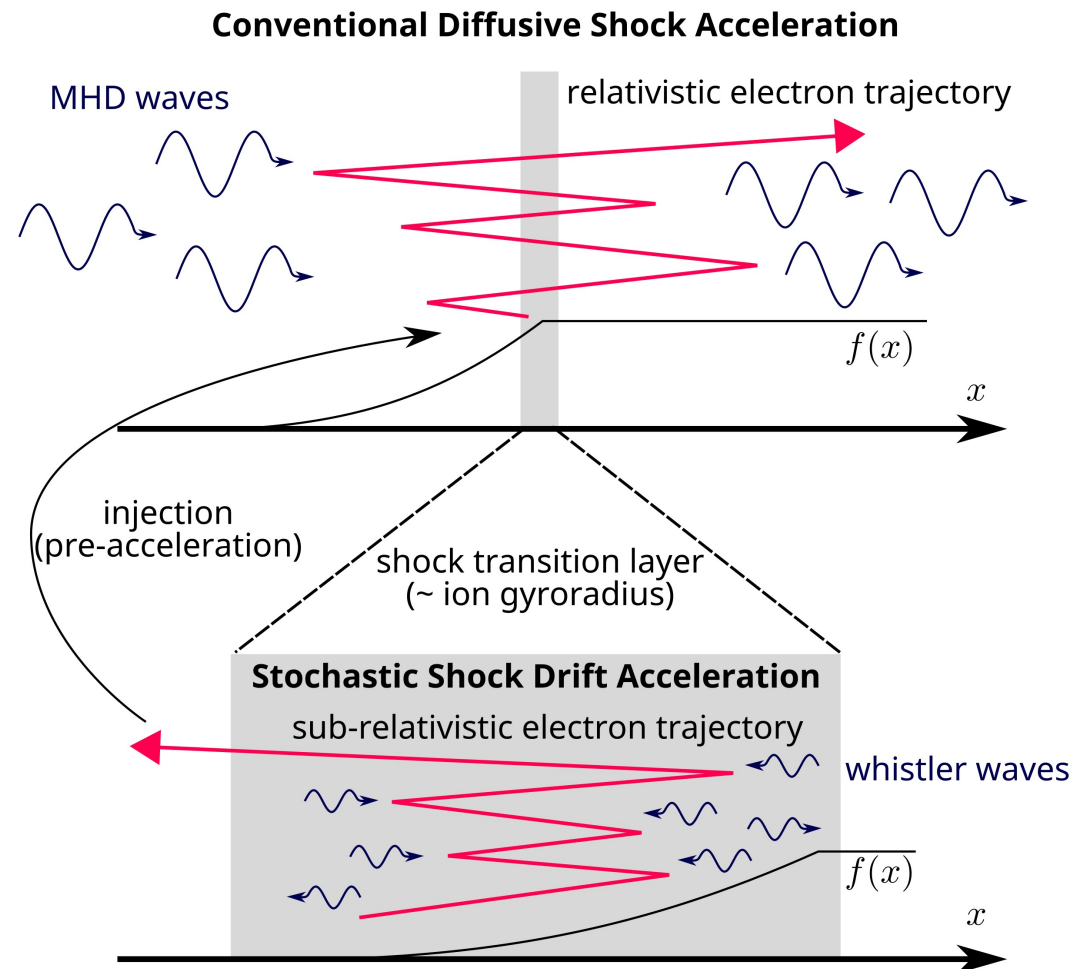
Electron Injection Scenario

DSA (diffusion length $\gg$ shock thickness)

- Diffusive and slow particle acceleration well beyond the shock thickness.
- The canonical power-law: $f(p) \propto p^{-4}$
- It may operate only when SSDA provides sufficiently energetic electrons.

SSDA (diffusion length $\sim$ shock thickness)

- Diffusive and fast particle acceleration within the shock transition layer.
- It results in a steeper power-law for energy-independent diffusion (consistent with observations at the bow shock.)
- Higher-energy electrons will eventually escape toward upstream because of diffusion lengths longer than the shock thickness.



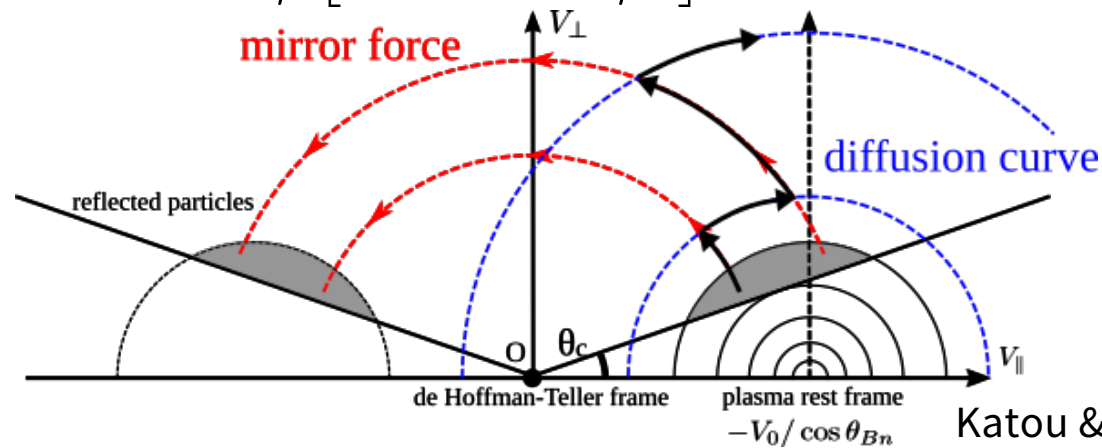
Amano & Hoshino (2022, ApJ), Amano+(2020, PRL), Katou & Amano (2019, ApJ), Kobzar+(2021, ApJ), Matsumoto+(2017, PRL)

Stochastic SDA (SSDA)

The transport of electrons in the de Hoffmann-Teller frame ($\mathbf{u} = u_{\parallel} \mathbf{b}$) may be governed by

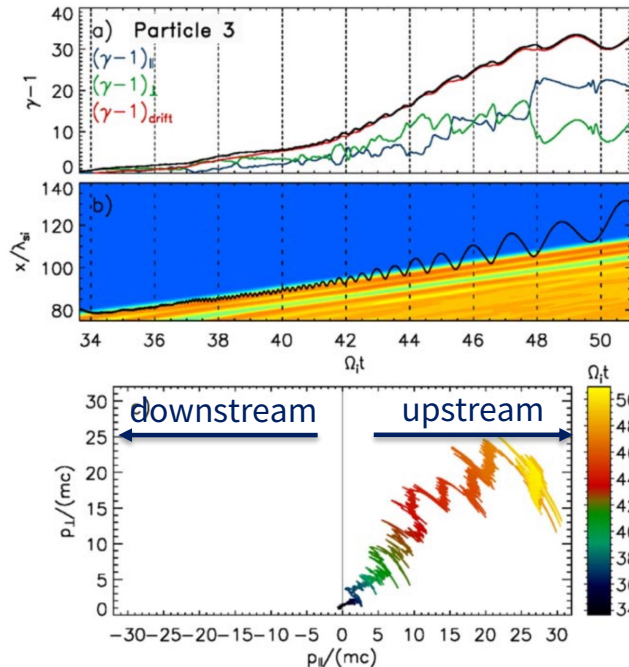
$$\begin{aligned}
 & \frac{\partial}{\partial t} f + (v\mu + u_{\parallel}) \frac{\partial}{\partial s} f \quad \text{negligible at Qperp shock (major term for DSA)} \\
 & + \left(\frac{1 - \mu^2}{2} u_{\parallel} \frac{\partial \ln B}{\partial s} - \mu^2 \frac{\partial u_{\parallel}}{\partial s} \right) v \frac{\partial f}{\partial v} \\
 & - \frac{1 - \mu^2}{2} \left((v\mu + u_{\parallel}) \frac{\partial \ln B}{\partial s} + 2\mu \frac{\partial u_{\parallel}}{\partial s} \right) \frac{\partial f}{\partial \mu} \\
 & = \frac{\partial}{\partial \mu} \left[(1 - \mu^2) D_{\mu\mu} \frac{\partial}{\partial \mu} f \right] + Q.
 \end{aligned}$$

mirror force (indicated by red arrows pointing to the boxed terms in the equation)



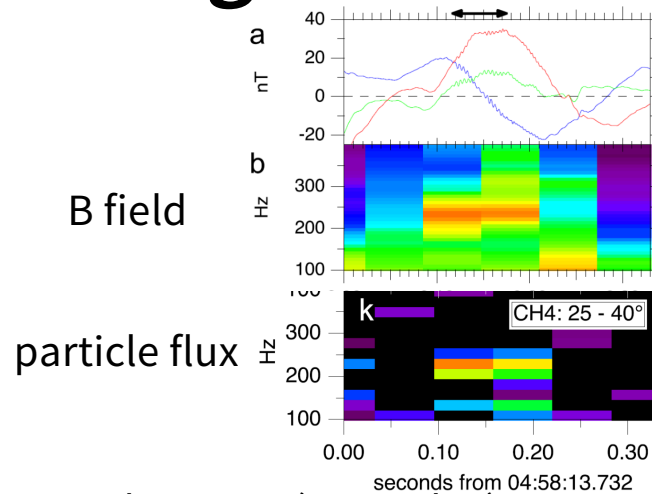
Katou & Amano (2019)

SSDA Signatures



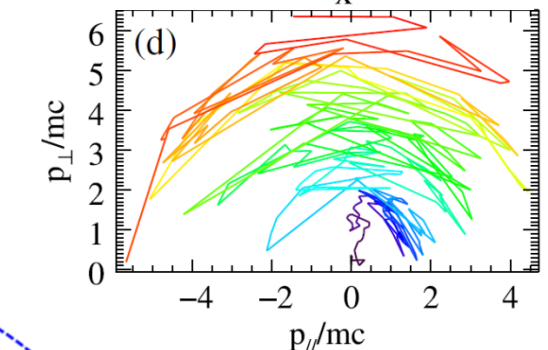
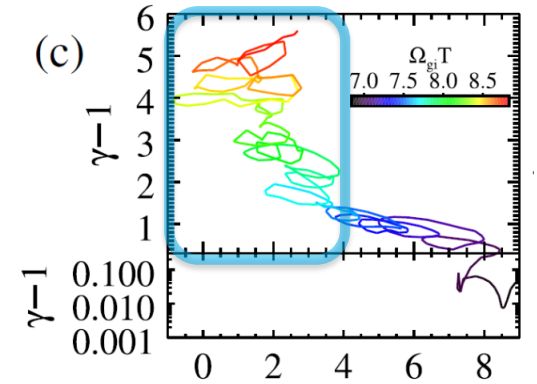
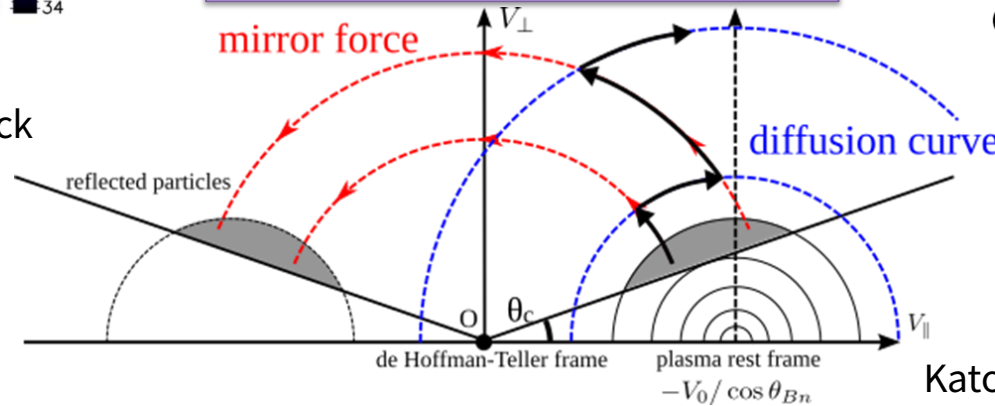
Kobzar+2021

2D high-beta and low Mach no. shock



Oka+2017 (MMS obs.)

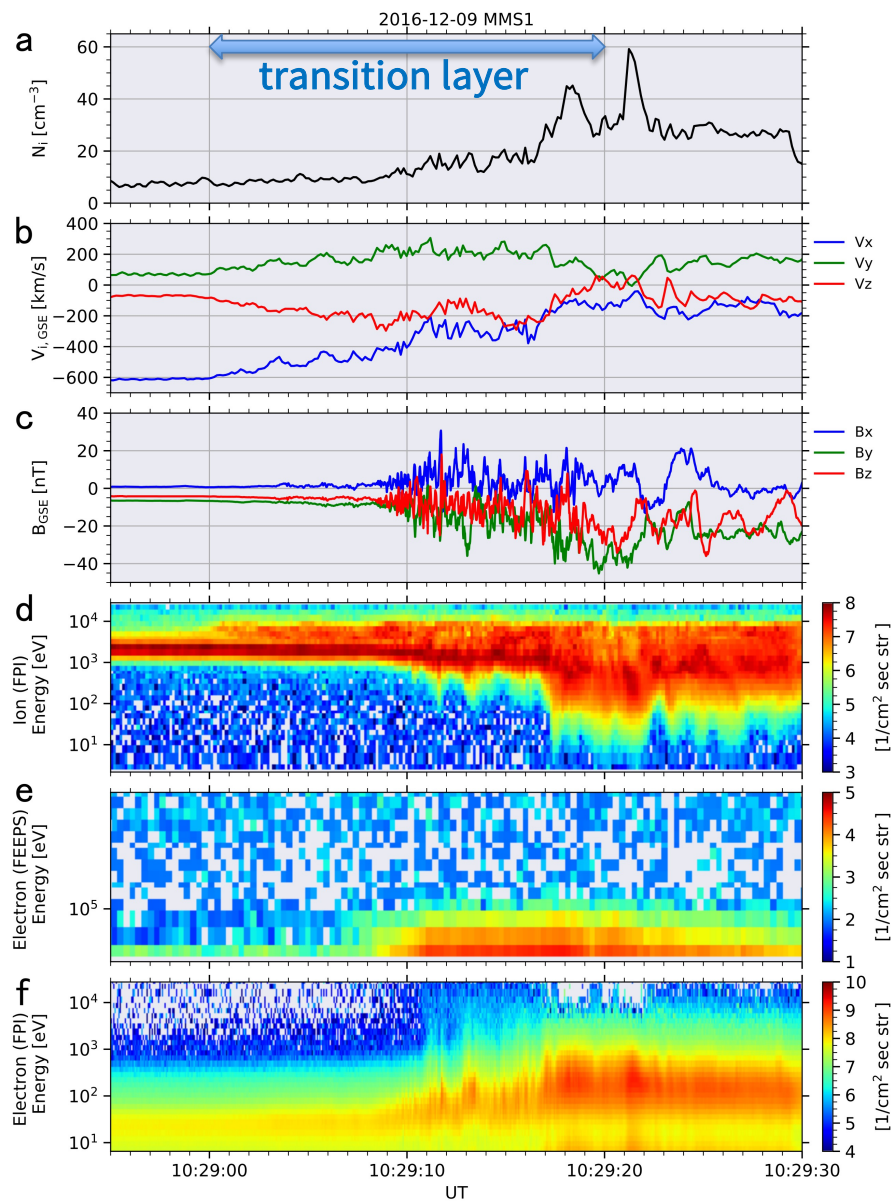
Stochastic Shock Drift Acceleration
mirror reflection + scattering



Matsumoto+2017

3D Weibel-dominated shock

Katou & Amano 2019

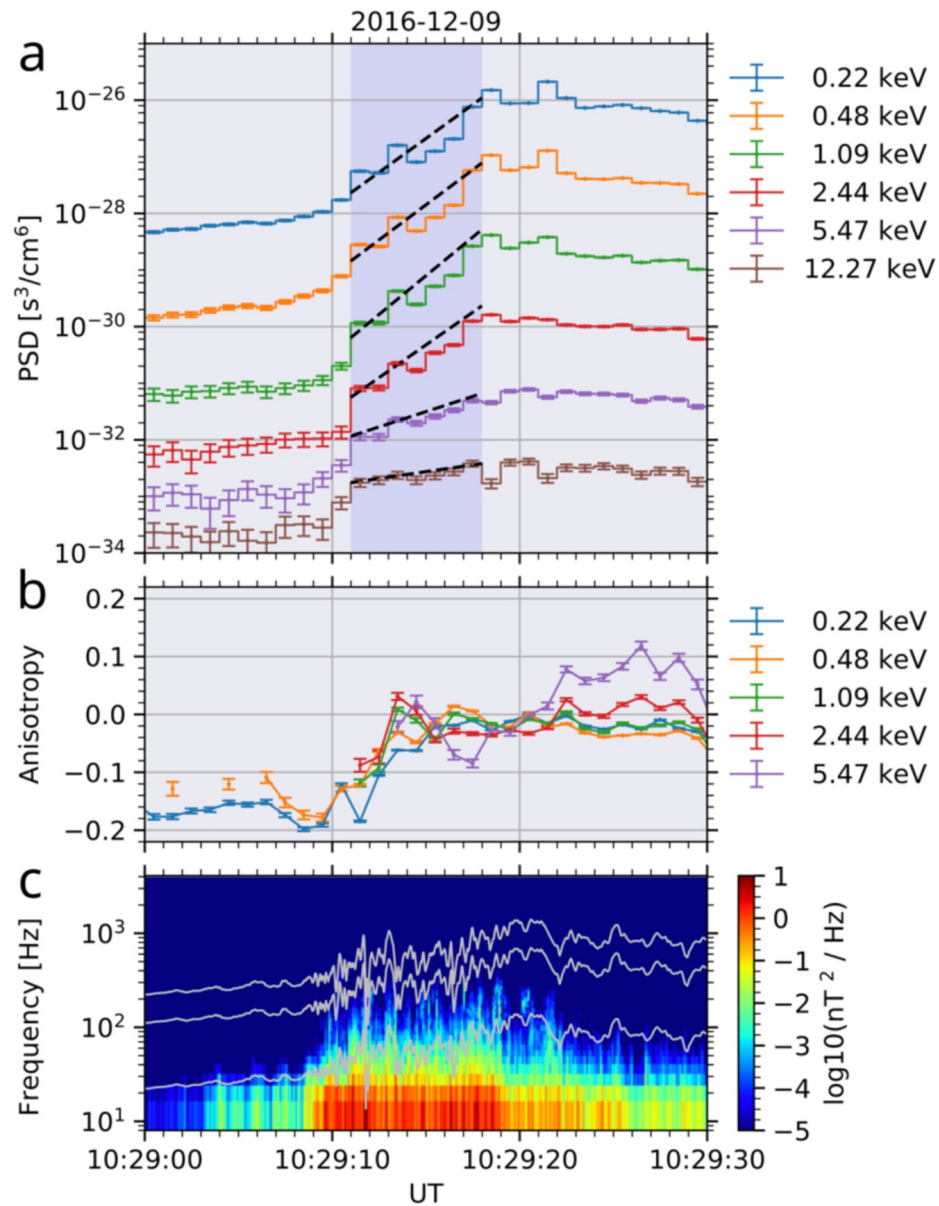


Bow Shock Crossing measured by MMS spacecraft

- $V_{\text{sw}} \sim 600 \text{ km/s}$
- $\theta_{\text{Bn}} \sim 85$ (quasi-perp.)
- $M_A \sim 8.9$ (high Mach num.)

Substantial flux enhancements for high energy ($>1 \text{ keV}$) electrons. FEEPS also detected electrons up to $\sim 100 \text{ keV}$. Unusual for bow shock crossings.

Evidence for the proposed electron injection scenario by MMS observation at Earth's bow shock [Amano+2020, PRL]



Exponential increase of particle intensity

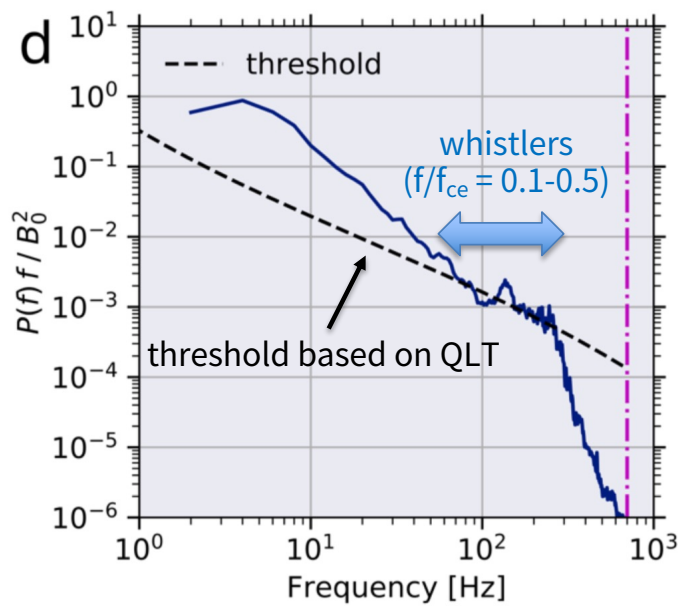
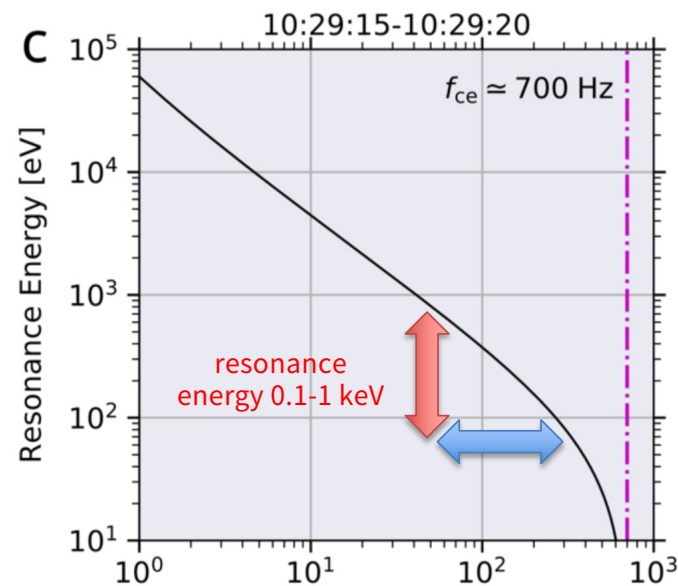
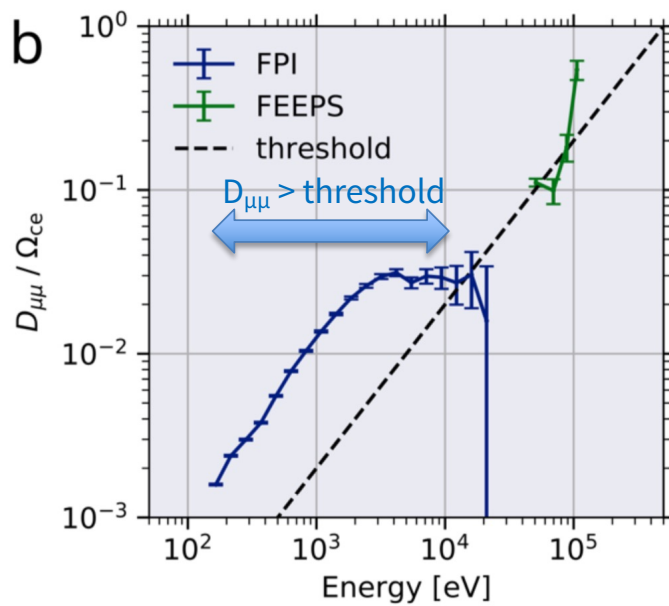
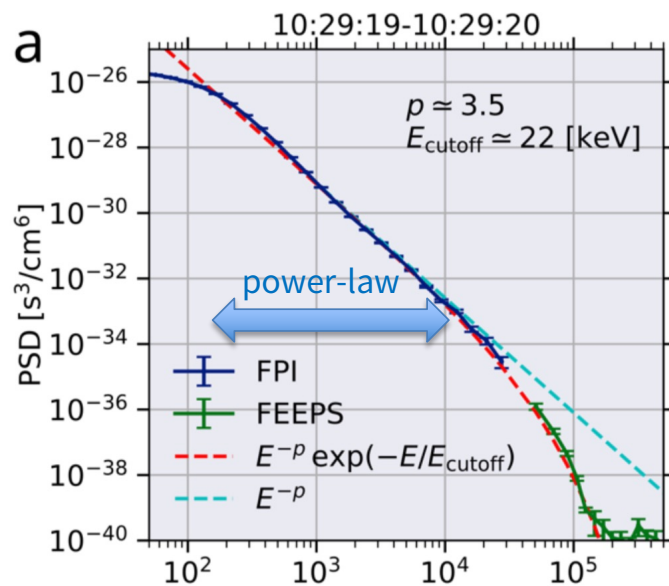
Nearly isotropic pitch-angle distribution

Enhanced wave power (in particular, high-frequency whistlers)

Smoking-Gun Evidence ?

- Simultaneous appearance of energetic electron profiles, weak anisotropy, enhanced wave power are all consistent with the theory. However, the agreement is only qualitative.
- The observed power-law index is roughly consistent with the theory. But it does not necessarily identify the mechanism.
- The theory predicts that the high-energy cutoff of the spectrum is determined by the single parameter $D_{\mu\mu}$

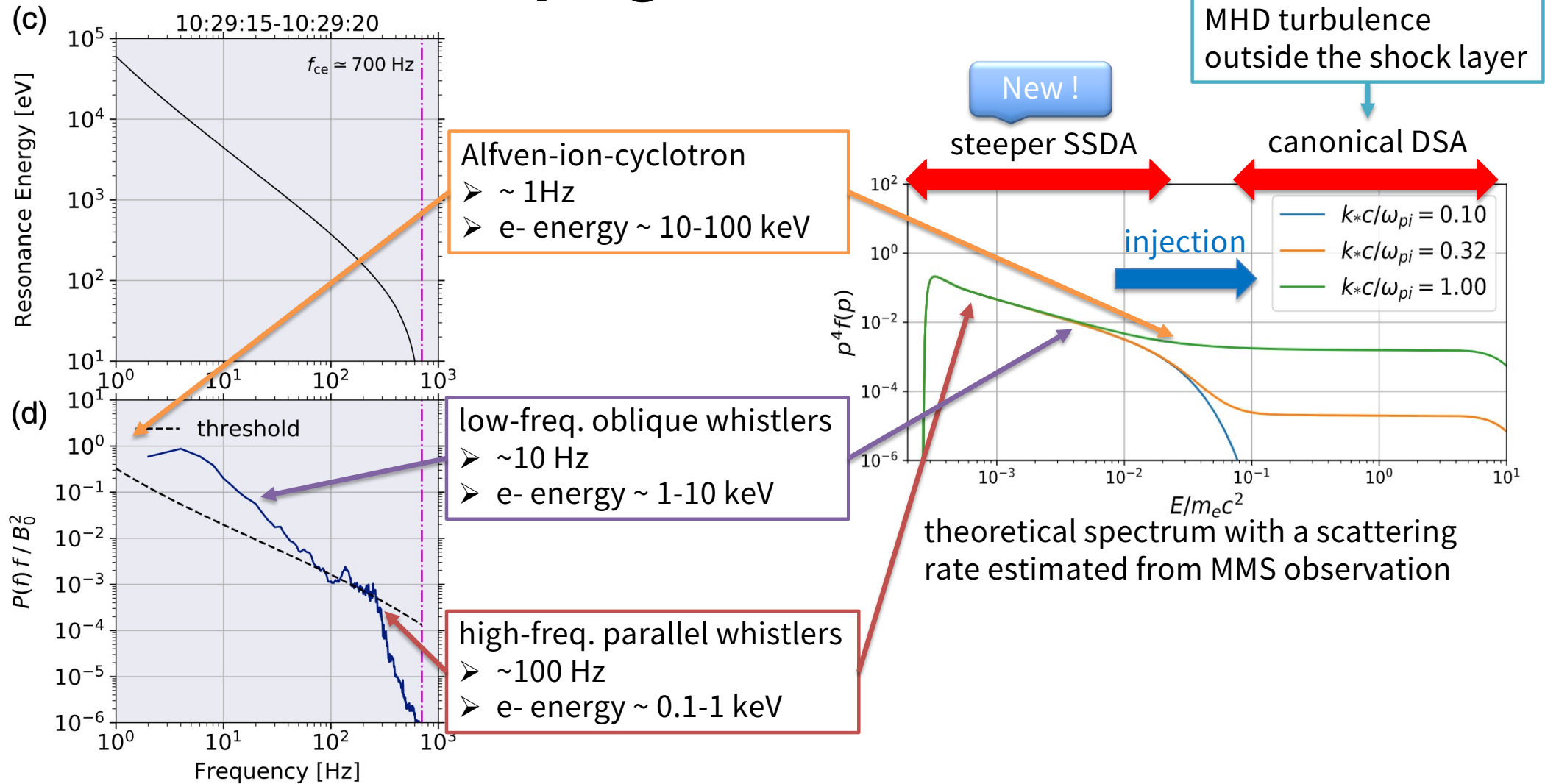
$$E_{\text{cutoff}} \sim E_{\text{sh}} \left(\frac{m_i}{m_e} \right) \left(\frac{D_{\mu\mu}}{\Omega_{\text{ce}}} \right) = \frac{1}{2} m_i \left(\frac{u_0}{\cos \theta_{\text{Bn}}} \right)^2 \left(\frac{D_{\mu\mu}}{\Omega_{\text{ce}}} \right)$$



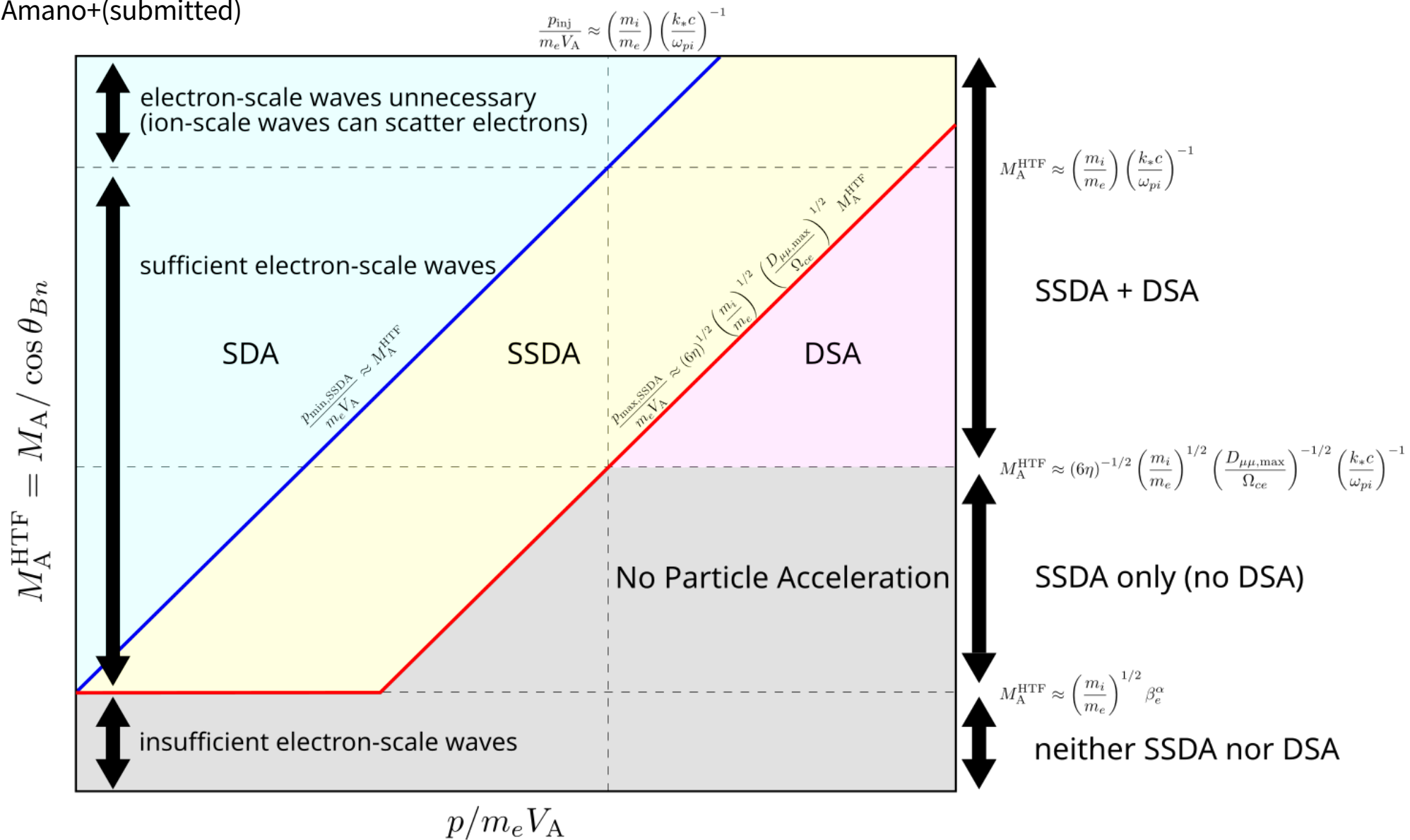
Amano+(2020, PRL)

Ref: Amano+2020 (PRL), Amano & Hoshino (2022), Amano+(submitted)

Unifying SSDA and DSA



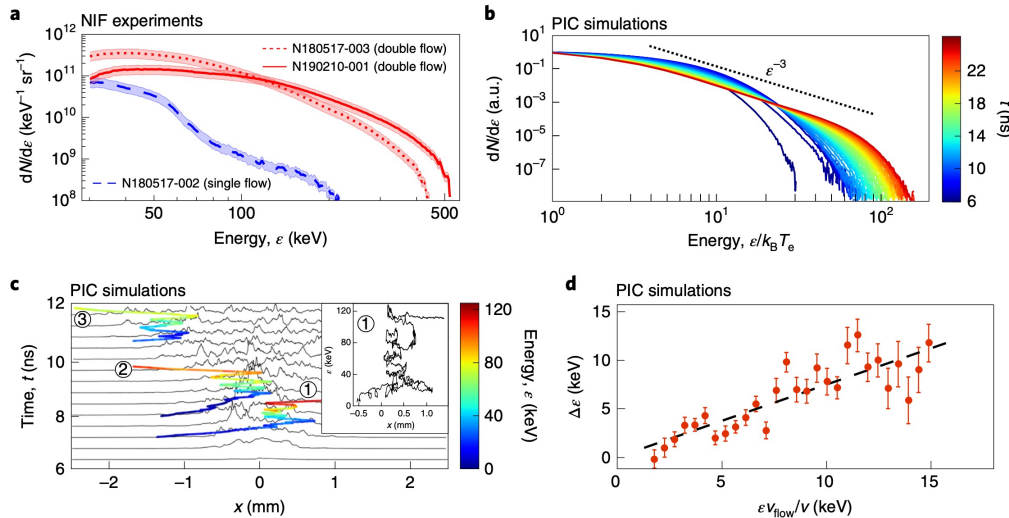
Ref: Amano+(submitted)



Conclusions

- We believe that, at present, Stochastic Shock Drift Acceleration (SSDA) is the most promising mechanism for electron injection at shocks.
- The theory predicts a steeper-than-DSA spectrum at sub-relativistic energies, which will be connected smoothly to the harder DSA at the transition energy of 0.1-1 MeV, if efficient electron injection is realized.
- It may be tested not only with in-situ spacecraft observations but also future astrophysical observations as well as laboratory laser experiments.

Fiuza+2020 (Laser Experiments @ NIF)



Tanaka+2018 (NuSTAR obs for W49B)

