



Abstract QR code



Analysis of Turbulence Energy Transfer at an Interplanetary Shock Observed by MMS

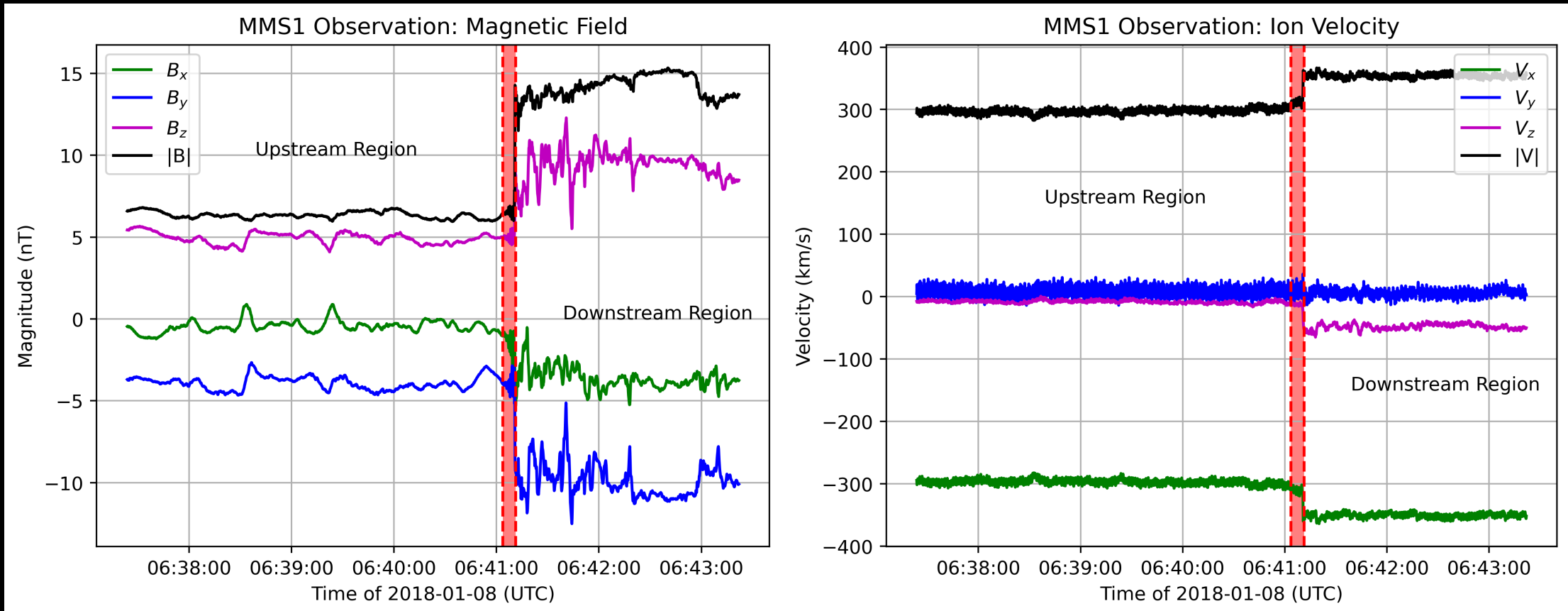
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Nawin Ngampoopun (1,2); David J Ruffolo (1); Riddhi Bandyopadhyay (3); William H Matthaeus (4);

(1) Department of Physics, Faculty of Science, Mahidol University, Thailand; (2) Mullard Space Science Laboratory, University College London, UK; (3) Department of Astrophysical Sciences, Princeton University, USA; (4) Department of Physics and Astronomy and Bartol Research Institute, University of Delaware, USA

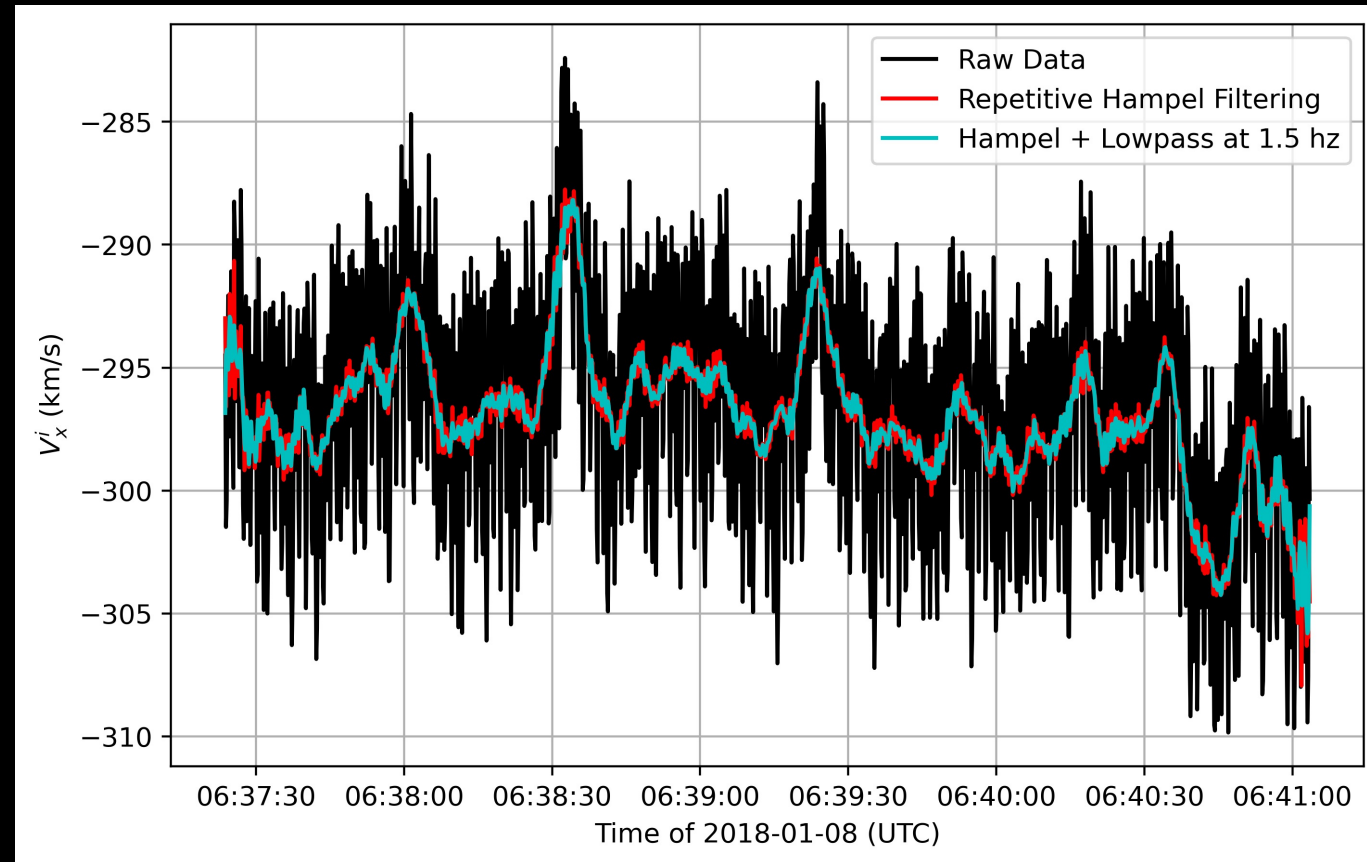
Introduction

On 2018/01/08 MMS observed an interplanetary shock driven by a corotating interaction region (CIR) on 06.41 UTC



Cleaning up FPI measurements of Solar Wind

- Issues of FPI (Fast plasma investigator): It is not designed for solar wind measurement
 - Periodic noise corresponds to spin tone period: ≈ 19.67 s and its harmonics
 - Large amplitude fluctuations
- Apply Hampel Filtering-Based Technique to remove artificial noise
 - Based on Bandyopadhyay et al. 2018



Energy Transfer Rate (ϵ) calculation

- Assuming Taylor's frozen in condition
- Inertial Scale : Kolmogorov-Yaglom Law (Third-Order Law)

$$Y^{\pm}(\tau) = -\frac{4}{3}\epsilon^{\pm}\langle V_{SW} \rangle \tau$$

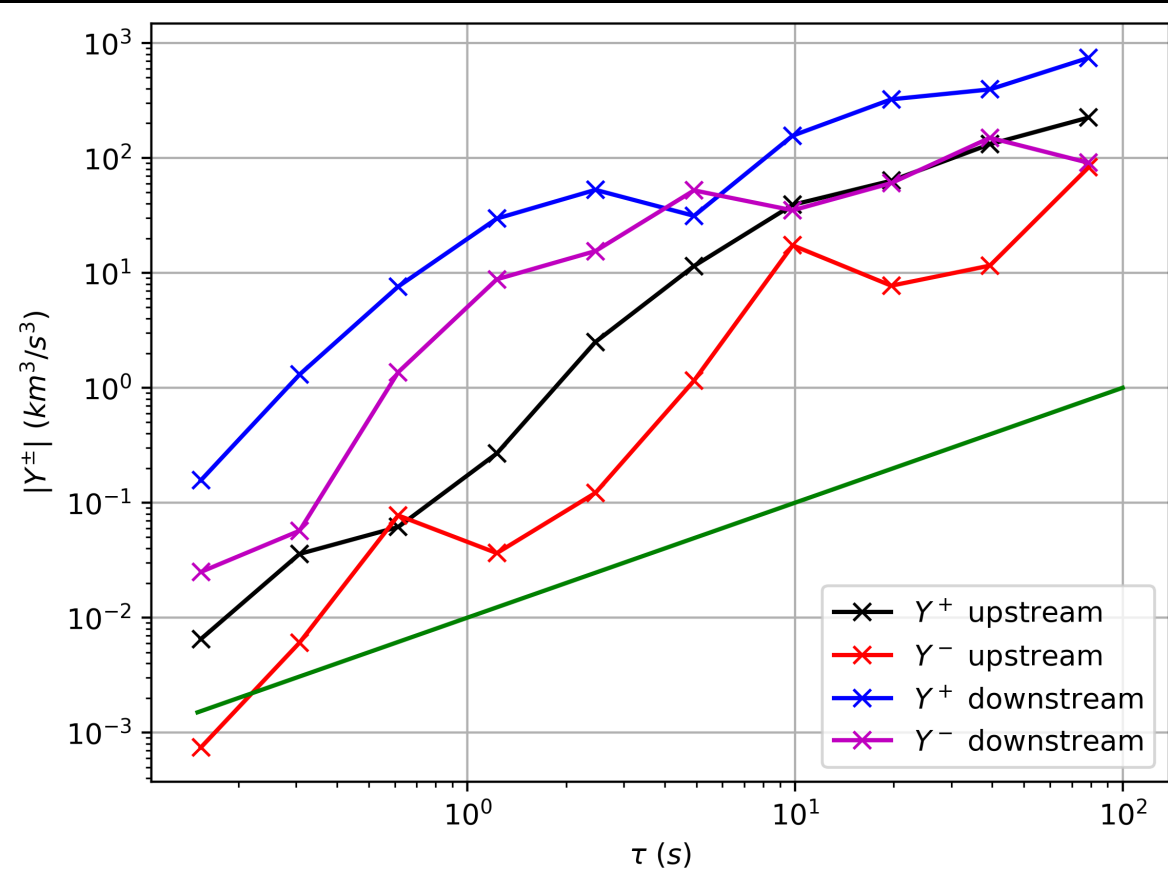
$$Y^{\pm}(\tau) = \left\langle |\Delta \mathbf{Z}^{\pm}(\tau; t)|^2 \Delta Z_R^{\mp}(\tau; t) \right\rangle, \quad \mathbf{Z}^{\pm} = \mathbf{v}_i \pm \mathbf{b}$$

- Energy containing scale: Von Kármán Analysis

$$\epsilon^{\pm} = -\frac{d(Z^{\pm})^2}{dt} = \alpha_{\pm} \frac{(Z^{\pm})^2 Z^{\mp}}{L_{\pm}}$$

- Mean energy transfer rate: $\epsilon = \frac{\epsilon^{+} + \epsilon^{-}}{2}$

Energy Transfer Rate: Inertial scale



| | ϵ (kJ kg ⁻¹ s ⁻¹) | |
|------|---|------------|
| | Upstream | Downstream |
| MMS1 | 4.9 ± 0.3 | 12 ± 1 |
| MMS2 | 4.7 ± 0.3 | 14 ± 1 |
| MMS3 | 5.0 ± 0.3 | 13.7 ± 0.9 |
| MMS4 | 4.6 ± 0.3 | 12 ± 1 |

Table 1. Turbulence energy transfer rate in inertial scale at upstream and downstream region of a shock. The values estimated from third-order law

Energy Transfer Rate: Energy containing scale

- Using Von Kármán – Howarth decay law
- $\alpha_+ \cong \alpha_- = 0.055$, due to high Reynolds number

| | $\epsilon \text{ (kJ kg}^{-1} \text{ s}^{-1}\text{)}$ | |
|------|---|------------|
| | Upstream | Downstream |
| MMS1 | 2.8 ± 0.5 | 20 ± 1 |
| MMS2 | 3.0 ± 0.6 | 20 ± 3 |
| MMS3 | 2.9 ± 0.6 | 21 ± 3 |
| MMS4 | 2.8 ± 0.5 | 19 ± 3 |

Table 2. Turbulence energy transfer rate at upstream and downstream region of a shock. The values estimated from Von Kármán analysis

Discussion Conclusion

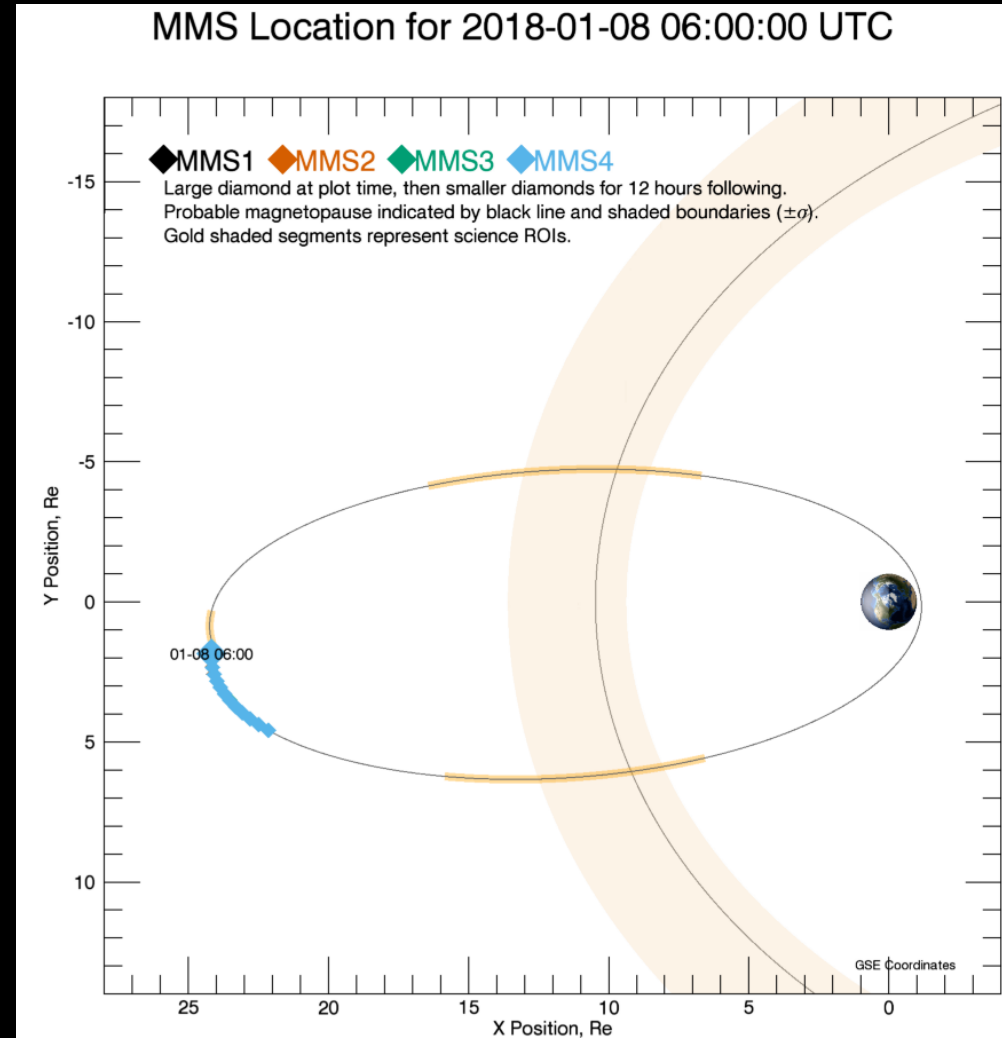
- In this work, we analyzed the energy dissipation rate of turbulence in the regions upstream and downstream of the interplanetary shock (IP) observed by the MMS on 8 January 2018.
- The turbulent energy transfer rate is calculated from both the third-order law (for the inertial range) and von Karman decay law (for the energy-containing range)
- **The analysis shows that the downstream region has a higher transfer rate, and the turbulence has developed to smaller scales compared to the upstream region.**
- Potentially further research should examine different IP shocks with a longer duration of burst mode data from MMS or other missions.



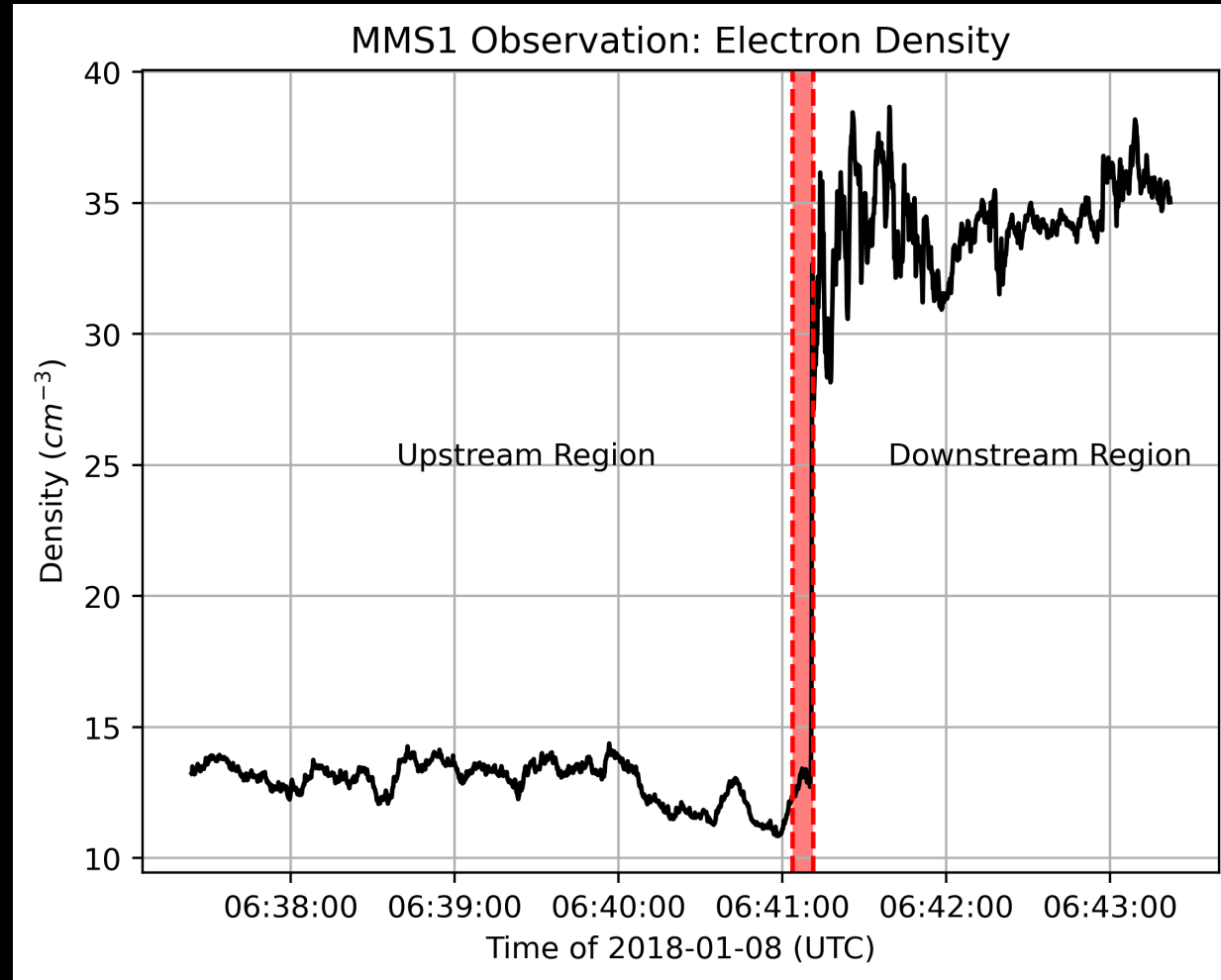
Appendix

*Not encouraged to take photo or share results after this slide

MMS position during the IP shock



Density profile at the IP shock

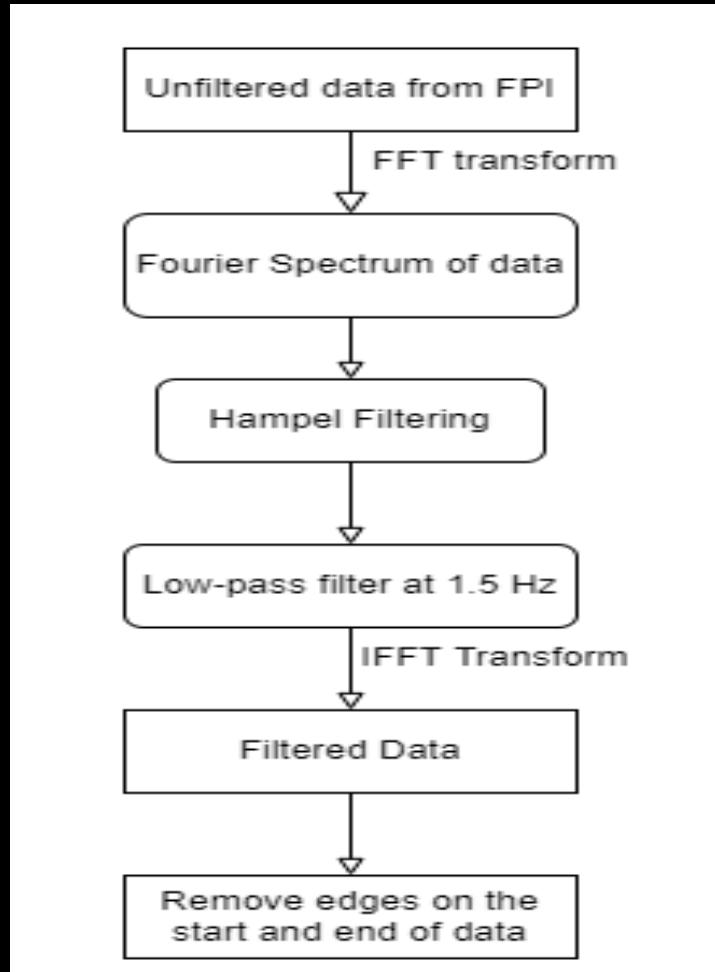


SW angle

- Angle between Solar wind velocity and X_GSE (Radial): $\theta_x = \frac{V_{i,x}}{|V_i|}$
- MMS1
 - θ_x upstream = 177.08 deg, θ_x downstream = 172.03 deg
- MMS2
 - θ_x upstream = 177.16 deg, θ_x downstream = 172.04 deg
- MMS3
 - θ_x upstream = 177.26 deg, θ_x downstream = 172.11 deg
- MMS4
 - θ_x upstream = 177.28 deg, θ_x downstream = 172.09 deg

Cleaning up ion velocity data

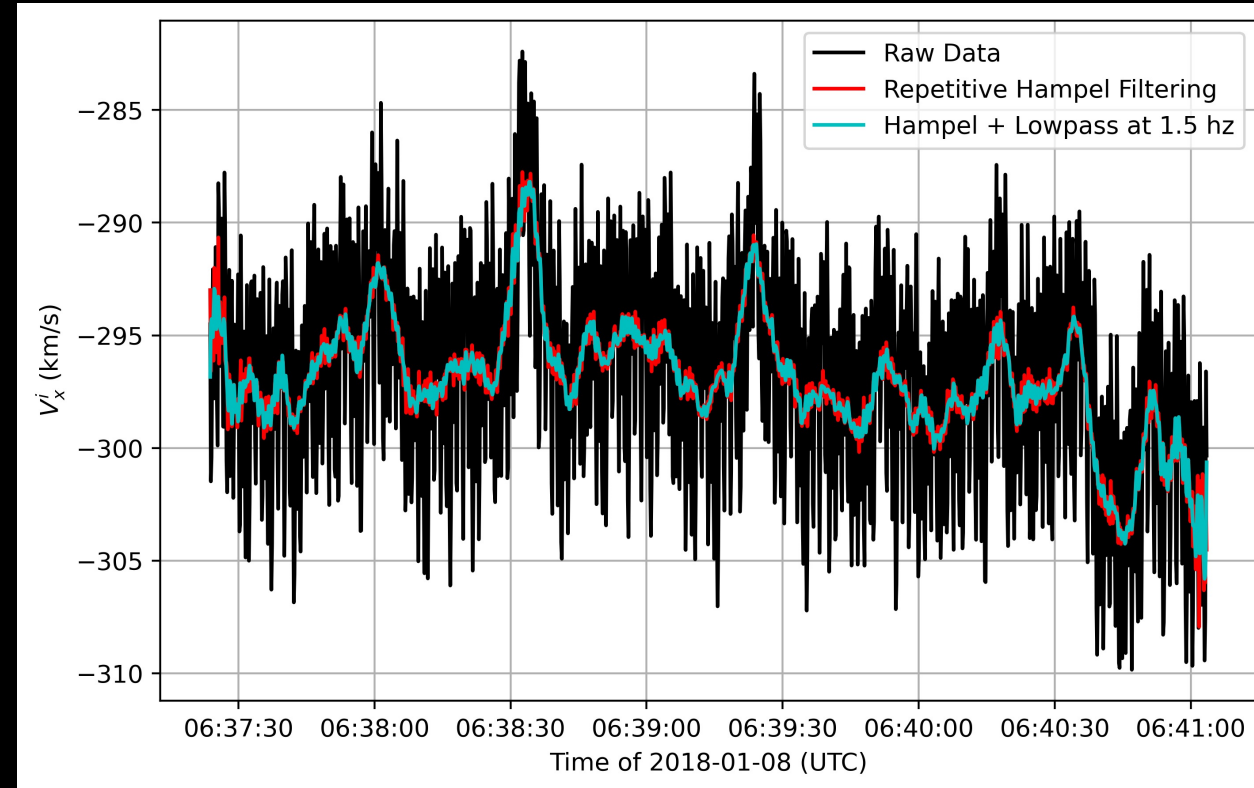
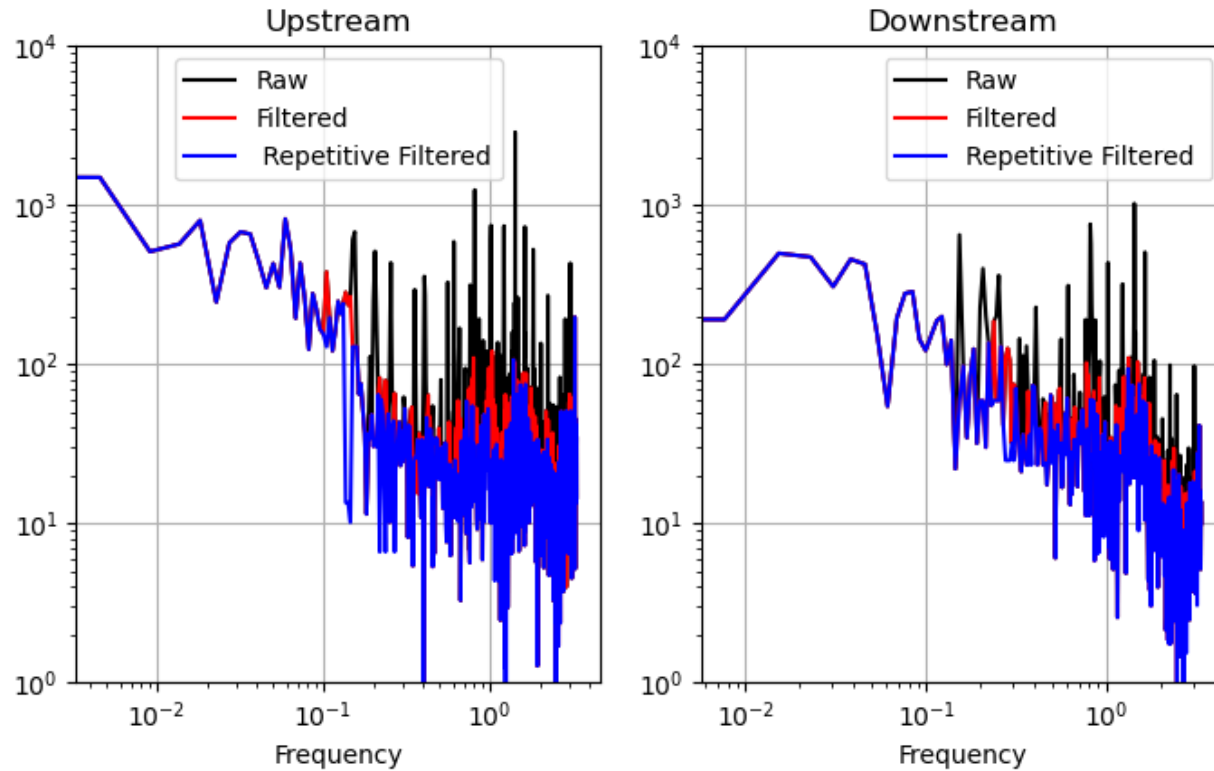
$$y_i = \begin{cases} x_i & \text{if } |x_i - m_i| \leq n_\sigma S_i, \\ m_i & \text{if } |x_i - m_i| > n_\sigma S_i, \end{cases}$$



- Apply Hampel Filtering - Based Technique (Bandyopadhyay et al. 2018)
 - Using Moving Window Length = 20 s
 - n-sigma = 1
 - Apply Hampel filtering two times in same data set
 - Low pass filter at 1.5 Hz
 - Remove 40 data points at edge
- Based on Bandyopadhyay et al. 2018

Effectiveness of Hampel Filtering

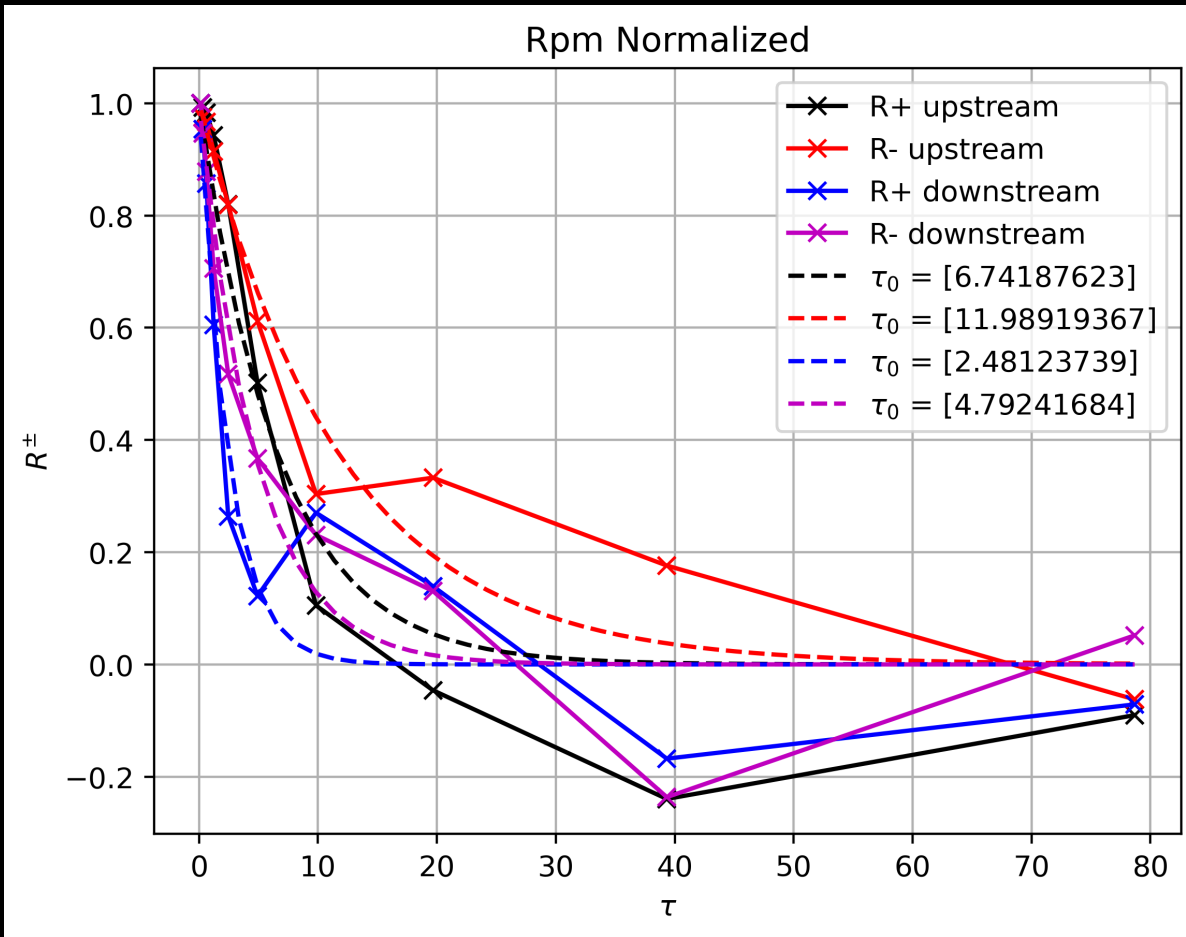
Fourier Spectrum - Ion Velocity Magnitude - Use Hampel Filtering (1 sigma)



Elsasser Increment Calculation

- Find $\Delta \mathbf{Z}^{\pm} = \mathbf{Z}^{\pm}(t + \tau) - \mathbf{Z}^{\pm}(t)$
- $\mathbf{Z}^{\pm} = \mathbf{v}_i \pm \mathbf{b}$; $\mathbf{v}_i = \mathbf{V}_i - \langle \mathbf{V}_i \rangle$, $\mathbf{b} = \mathbf{B} - \langle \mathbf{B} \rangle$ (in Alfven unit)
- τ choice: $\tau = 2^n * T_{spin}/32$
 - $n = (-2, -1, \dots, 6, 7)$
 - T_{spin} (spin period of S/C) = 19.67
 - Reduced more spin tone effect.

Correlation Function



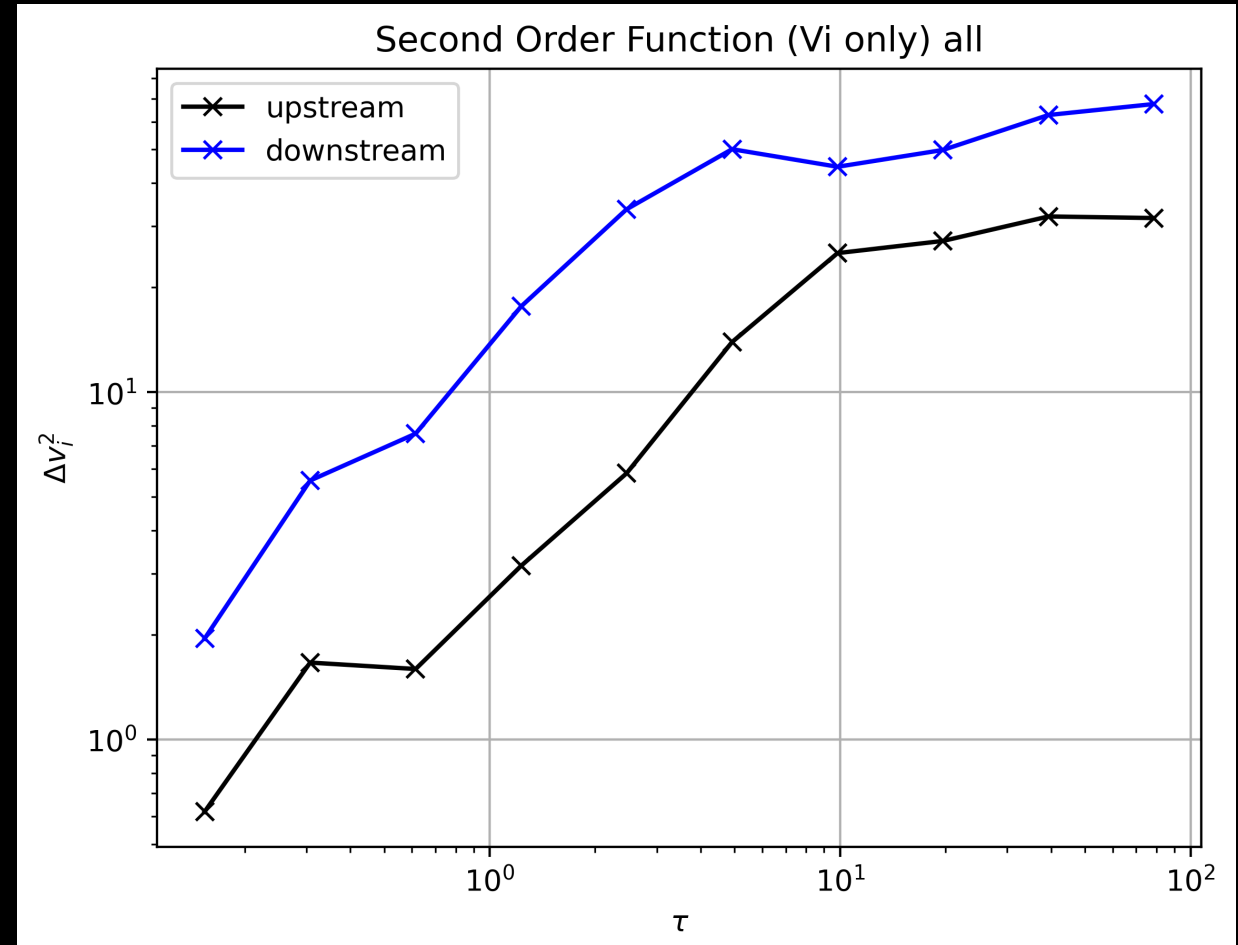
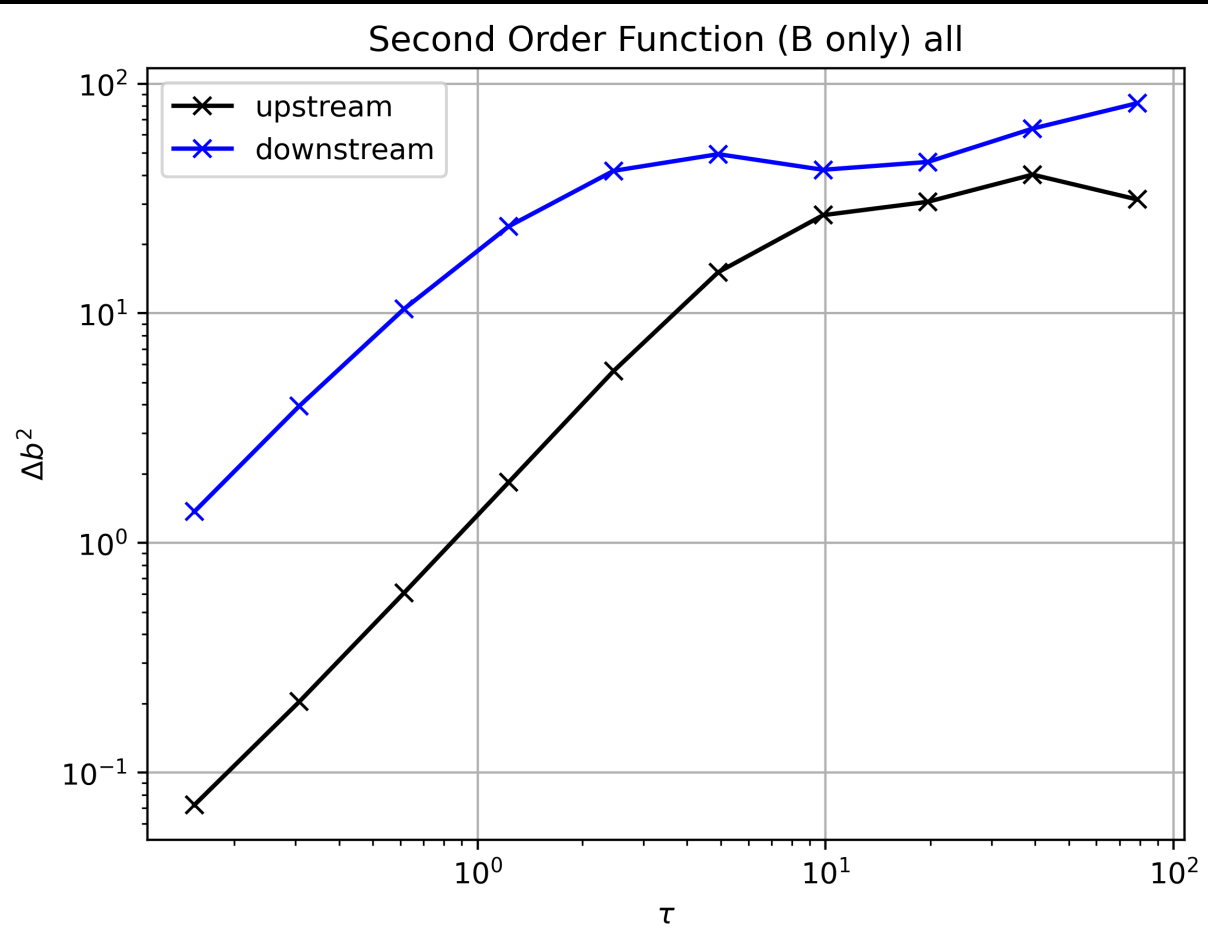
$$R^\pm(\tau) = \langle Z^\pm(t) \cdot Z^\pm(t + \tau) \rangle_T.$$

$$R^\pm(\tau^\pm) = \frac{1}{e},$$

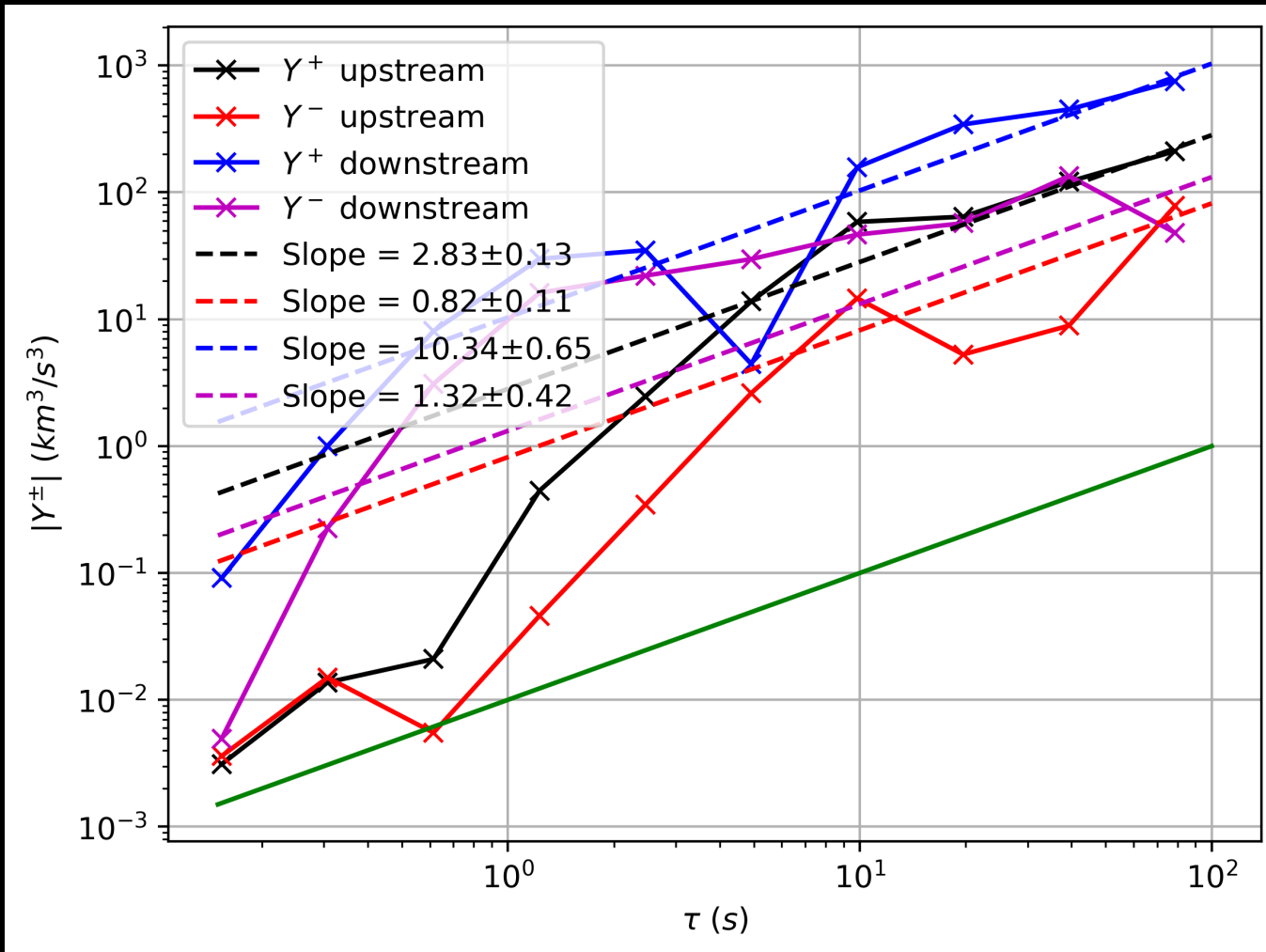
$$L_\pm = |\langle V \rangle| \tau^\pm,$$

- Plot of normalized correlation function with fitted 1/e function
- Correlation scale τ_0
 - From 1/e method
 - τ_0 range from 2 – 12 second
 - Due to availability of data, maybe correlation scale is shorter than it should be for pristine SW.

Second Order Function: B and V_i

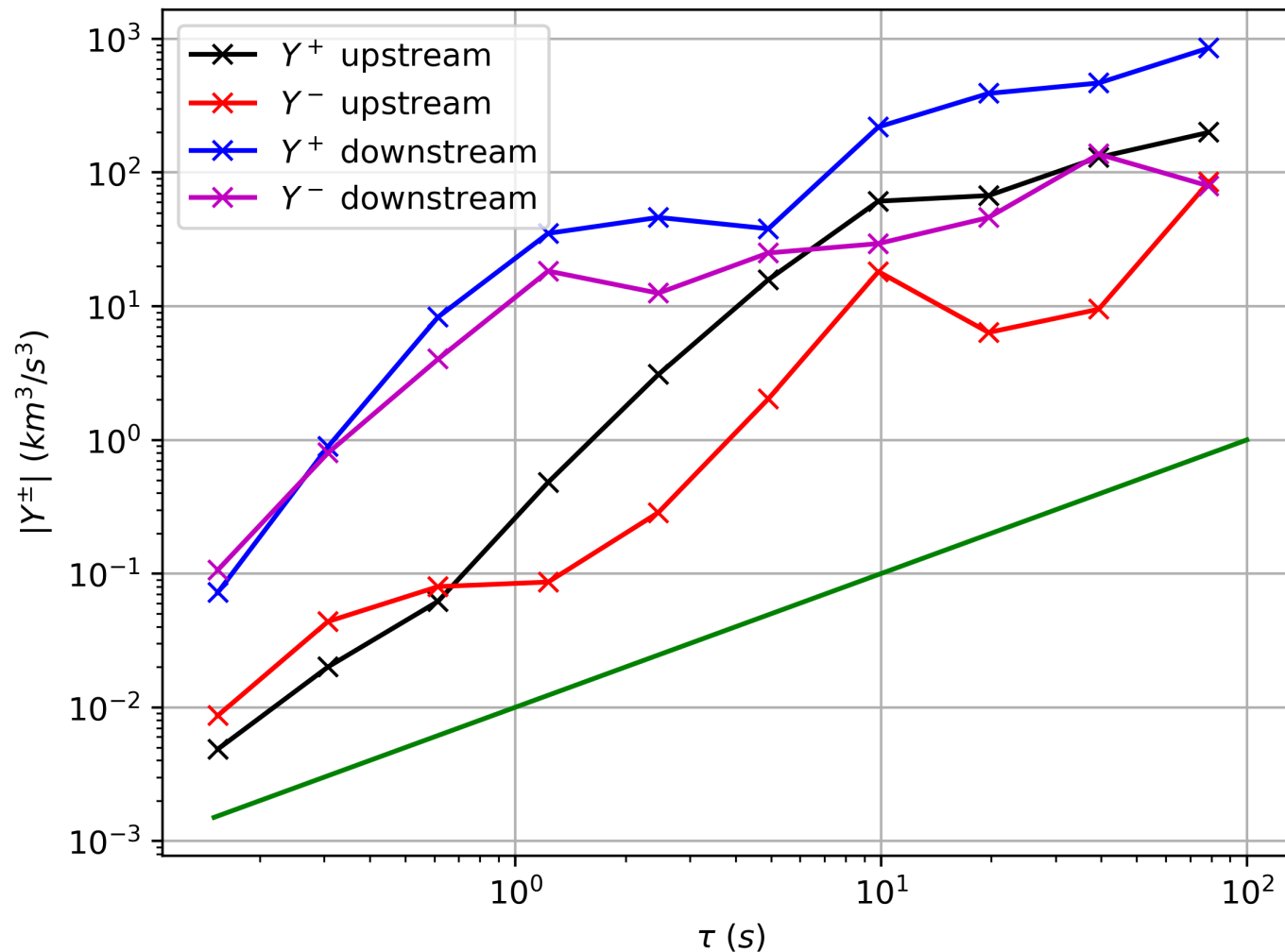


MMS1 – Third Order Law



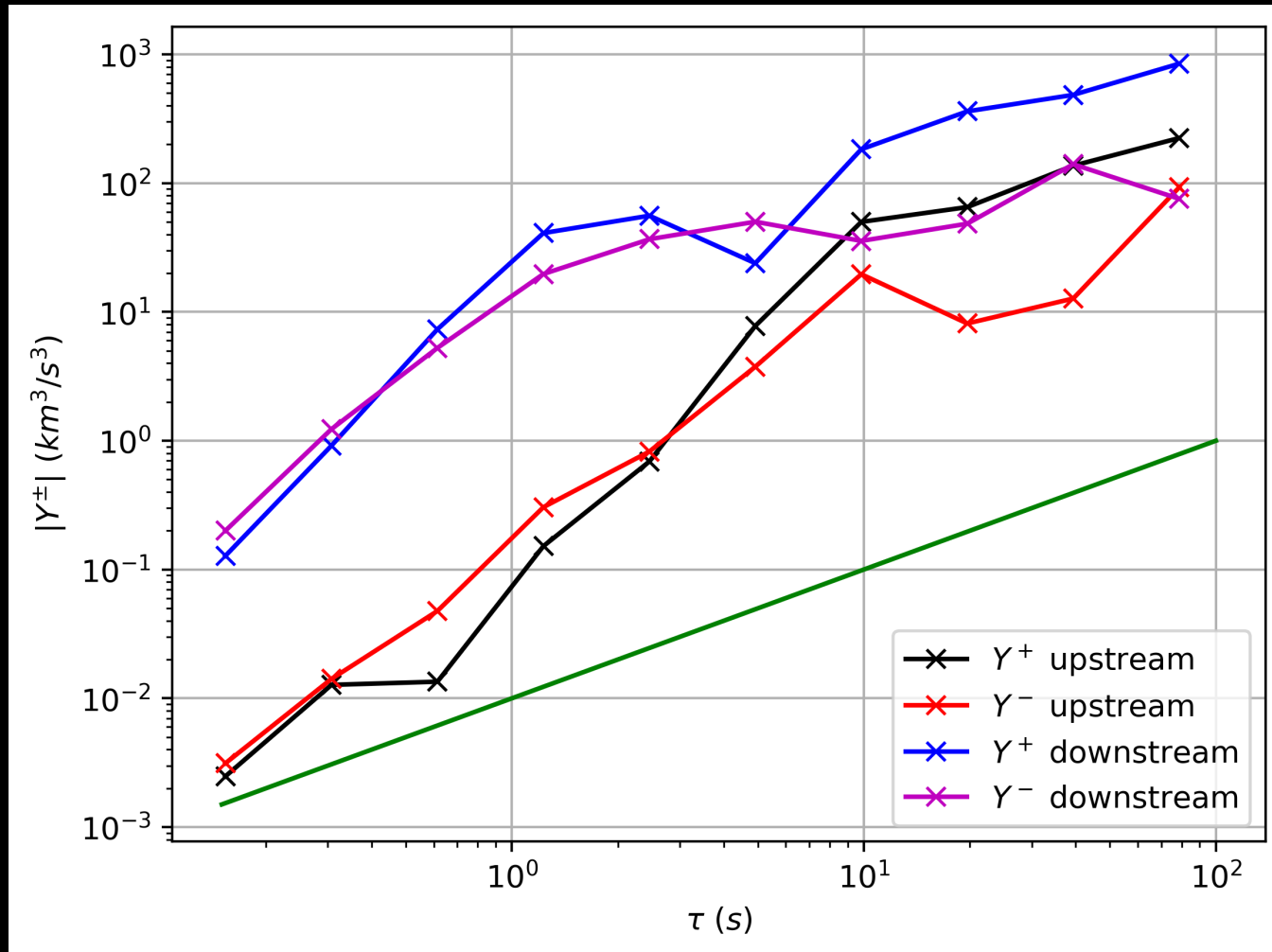
- ϵ^+ upstream = $7.5 \pm 0.2 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- upstream = $2.2 \pm 0.3 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^+ downstream = $21 \pm 1 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- downstream = $3.9 \pm 0.9 \text{ kJ kg}^{-1} \text{ s}^{-1}$

MMS2 - Third Order Law



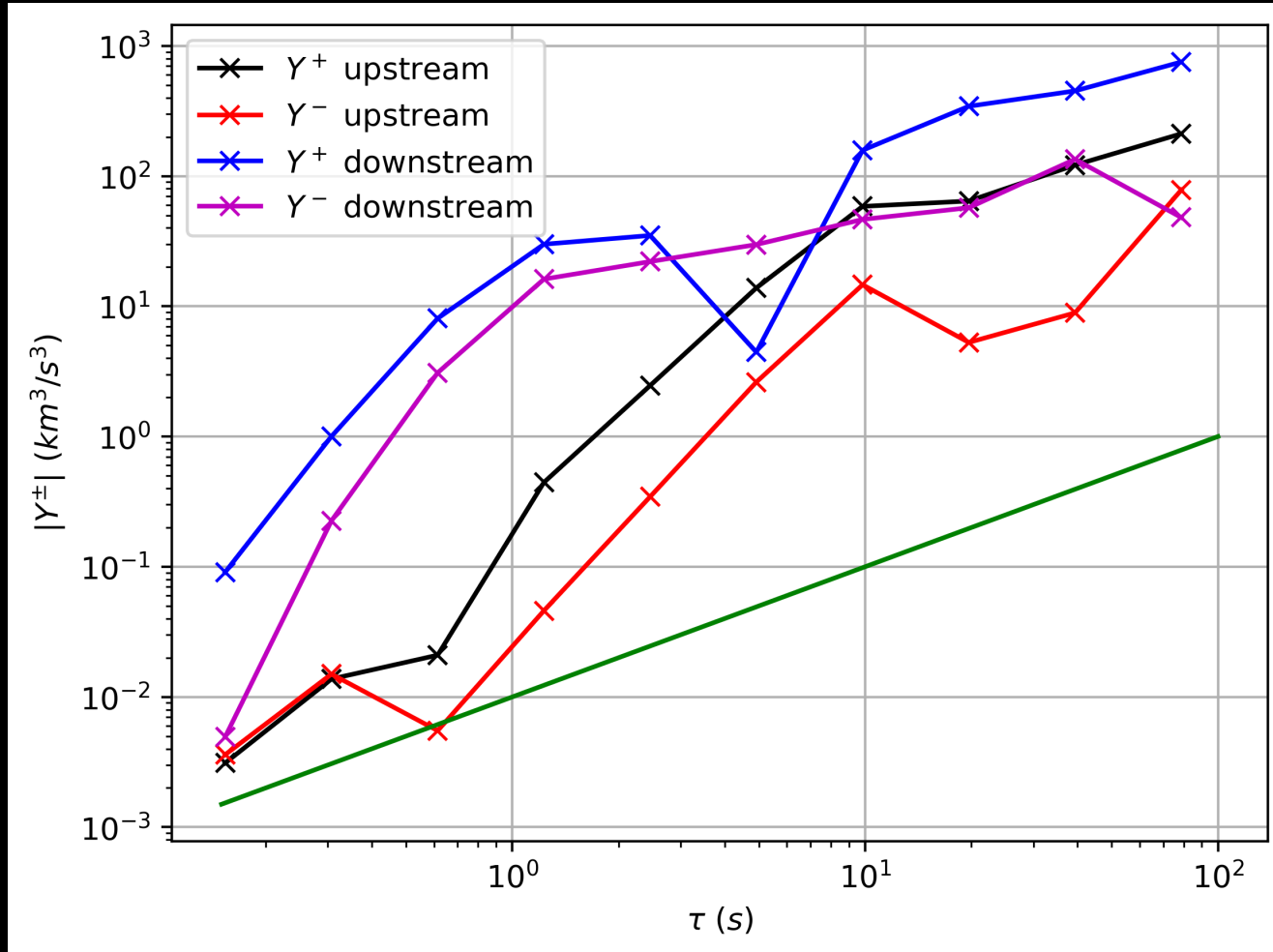
- ϵ^+ upstream = $7.0 \pm 0.4 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- upstream = $2.3 \pm 0.3 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^+ downstream = $25 \pm 2 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- downstream = $3.3 \pm 0.7 \text{ kJ kg}^{-1} \text{ s}^{-1}$

MMS3 - Third Order Law



- ϵ^+ upstream =
 $7.6 \pm 0.3 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- upstream =
 $2.5 \pm 0.3 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^+ downstream =
 $24 \pm 1 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- downstream =
 $3.4 \pm 0.8 \text{ kJ kg}^{-1} \text{ s}^{-1}$

MMS4 - Third Order Law



- ϵ^+ upstream = $7.1 \pm 0.3 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- upstream = $2.1 \pm 0.2 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^+ downstream = $22 \pm 1 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- downstream = $2.8 \pm 0.9 \text{ kJ kg}^{-1} \text{ s}^{-1}$

Von Karman Analysis – MMS1 & 2

- MMS3

- ϵ^+ upstream = $4.6 \pm 0.9 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- upstream = $1.1 \pm 0.1 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^+ downstream = $29 \pm 5 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- downstream = $10 \pm 1 \text{ kJ kg}^{-1} \text{ s}^{-1}$

- MMS4

- ϵ^+ upstream = $4.6 \pm 0.9 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- upstream = $1.3 \pm 0.3 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^+ downstream = $31 \pm 5 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- downstream = $8.8 \pm 0.2 \text{ kJ kg}^{-1} \text{ s}^{-1}$

Von Karman Analysis – MMS3 & 4

- MMS3

- ϵ^+ upstream = $4.5 \pm 0.9 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- upstream = $1.2 \pm 0.2 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^+ downstream = $32 \pm 5 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- downstream = $10 \pm 2 \text{ kJ kg}^{-1} \text{ s}^{-1}$

- MMS4

- ϵ^+ upstream = $4.4 \pm 0.8 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- upstream = $1.2 \pm 0.2 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^+ downstream = $29 \pm 4 \text{ kJ kg}^{-1} \text{ s}^{-1}$
- ϵ^- downstream = $9.8 \pm 2 \text{ kJ kg}^{-1} \text{ s}^{-1}$