

POLAR HEAT TRANSPOT ENHANCEMENT IN SUB-GLACIAL OCEANS ON ICY MOONS

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EGU General Assembly 2023



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Motivation: Dynamics in sub-glacial oceans on icy moons





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www.scitechdaily.com/images/Enceladus-Plume-768x614.jpg





Direct numerical simulations (DNS) of spherical rotating Rayleigh-Bénard convection (RRBC)

• Control parameters:

- Fluid properties:
- Thermal driving:
- Radius ratio:
- Rotation rate:

$$Pr = \frac{v}{\kappa} = 4.38$$

$$Ra = \frac{\alpha g \Delta T H^3}{\nu \kappa} = 10^6$$

$$\eta = r_i / r_o = \{0.6, 0.8\}$$

$$Ro^{-1} = \frac{2\Omega H}{\sqrt{\alpha g \Delta T H}}$$

$$(Ek^{-1} = \frac{2\Omega H^2}{\nu})$$





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• Heat transport:

Ο

- (Momentum transport:
- $Nu = \frac{qH}{\kappa\Delta T}$ $Re = \frac{UH}{v}$

 $(Ek^{-1} = \frac{2\Omega H^2}{N})$





• Ekman pumping through vertically coherent vortices













Relevance for icy moons

Heat transport enhancement in planar RRBC:

 \circ rotation-affected regime

 \circ Pr > 1

Icy moon oceans:

- o rotation-affected regime
- $Pr \in [10,13] > 1$







Latitudinal vs. global heat transport

Do we observe heat transport enhancement?

• Within the tangent cylinder: Yes!





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Latitudinal vs. global heat transport

Do we observe heat transport enhancement?

- Within the tangent cylinder: Yes!
- $\circ~$ Globally: No! Polar enhancement balanced by equatorial reduction







- **Non-rotating, pure RB** ($\operatorname{Ro}^{-1} = 0$)
 - Radial, buoyant plumes
 - Persistent "large-scale circulation"

Polar view (South):

Equatorial view:

∢X __Z

X Y





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- <u>Non-rotating, pure RB</u> ($Ro^{-1} = 0$) Ο
 - Radial, buoyant plumes
 - Persistent "large-scale circulation"

 $(\text{Ro}^{-1} = 0.3)$ **Weak rotation** Ο

- Radial, buoyant plumes
- Homogeneous distribution \succ





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Equatorial view:



- <u>Non-rotating, pure RB</u> ($Ro^{-1} = 0$) 0
 - Radial, buoyant plumes
 - Persistent "large-scale circulation"
- $(\text{Ro}^{-1} = 0.3)$ Weak rotation Ο
 - Radial, buoyant plumes
 - Homogeneous distribution
- $(Ro^{-1} = 3)$ **Moderate rotation** Ο
 - Taylor-Columns, sheet-like plumes
 - Polar heat transport enhancement



Equatorial view:

X Y



All figures from Hartmann et al. (in prep.)



- <u>Non-rotating, pure RB</u> ($Ro^{-1} = 0$)
 - Radial, buoyant plumes
 - Persistent "large-scale circulation"
- <u>Weak rotation</u> $(Ro^{-1} = 0.3)$
 - Radial, buoyant plumes
 - Homogeneous distribution
- <u>Moderate rotation</u> $(Ro^{-1} = 3)$
 - Taylor-Columns, sheet-like plumes
 - Polar heat transport enhancement

o <u>Strong rotation</u>

$$(\text{Ro}^{-1} = 15.9)$$

- Taylor-Columns, sheet-like plumes
- Equatorial dominated convection



 (Towards onset of)
 Diffusion-free behavior
 [Gastine et al. (JFM,2016) Wang et al. (GRL,2021)]



X Z

X Y

Equatorial view:





 $\Theta = \langle \Theta \rangle$ = 0.3

- 0.2

- 0.1

- 0

-0.1

-0.2



All figures from Hartmann et al. (in prep.)



- Radius ratio for sub-glaical oceans on icy moons:
- \circ $\,$ Increasing the radius ratio:

 $\eta = 0.6 \rightarrow \eta = 0.8$

[Soderlund (GRL, 2019), Vance et al. (JGR Planets, 2018)]

Europa

0.92 - 0.94

Ganymede

0.80 - 0.99



η

Enceladus

0.74 – 0.95

Titan

0.83 - 0.96



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- Increasing the radius ratio: Ο
 - $\eta = 0.6 \rightarrow \eta = 0.8$



Lat=53.1

1.5 H

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4H

Lat=36.9



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• **Ekman pumping** to supply the vortices with cold/hot fluid **most efficient** for equal thicknesses of thermal BL λ_{Θ} and kinetic BL λ_u :

$$\Rightarrow \lambda_{\Theta}/\lambda_u \approx 1$$



From Stevens et al. (NJP 12, 2010)



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○ Heat transport maximum Nu_{max} predicted at: $0.8 \approx \lambda_{\Theta}/\lambda_u \propto Ra \ Ek^{3/2} \Rightarrow Ro_{opt}^{-1}$ [King et al., JFM 691, 2012,] [Yang et al., PRF 5, 2020]



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In spherical RRBC:

- Perfect agreement between observed and predicted Ro_{opt}^{-1} (in the tangent cylinder)
- $\circ \quad \mbox{Perfect agreement between observed and} \\ \mbox{predicted } \lambda_\Theta / \lambda_u \mbox{ at } {\rm Ro}_{\rm opt}^{-1} \\ \end{tabular}$



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In spherical RRBC:

- Perfect agreement between observed and predicted Ro_{opt}^{-1} (in the tangent cylinder)
- Perfect agreement between observed and predicted $\lambda_{\Theta}/\lambda_u$ at $\mathrm{Ro}_{\mathrm{opt}}^{-1}$
- Enhancement efficiency still controlled by the boundary layer ratio $\lambda_{\Theta}/\lambda_u$





Radial gravity profile: $g(r) \propto r^{\gamma}$ \circ So far:g(r) = const.

 \circ Now: $g(r) \propto r$ $g(r) \propto 1/r^2$

Pr = 4.38
$Ra = 10^{6}$
$\eta = 0.6$

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No (significant) influence on:

Influence on:





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- \circ Optimal rotation rate Ro_{opt}^{-1}

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- Enhancement Nu/Nu₀ for Ro⁻¹ < Ro⁻¹_{opt}
 (buoyancy-dominated & rotation-affected regimes)

Influence on:





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Influence on:

- Enhancement Nu/Nu₀ for Ro⁻¹ > Ro⁻¹_{opt} (rotation-dominated regime)
- Onset/breakdown of convection (Ra_c)





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- \circ Boundary layer crossing $\lambda_{\Theta}/\lambda_{u}$

Influence on:

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Pr = 4.38

 $Ra = 10^{6}$



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BY

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 - Ekman pumping through axially aligned Taylor columns within the tangent cylinder
 - Globally balanced by equatorial reduction





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> Potential reason for latitudinal variations of the crustal thickness on icy moons



References

- [1] Kunnen R.P.J., Stevens R.J.A.M., Overkamp J., Sun C., van Heijst G.J.F., and Clercx H.J.H., "The role of Stewartson and Ekman layers in turbulent rotating Rayleigh-Bénard convection", J. Fluid Mech. 688 (2011), 422-442.
- [2] Hartmann R., Stevens R.J.A.M., Lohse D., and Verzicco R., "Polar heat transport enhancement in sub-glacial oceans on icy moons", in prep. (2023)
- [3] Soderlund K.M., "Ocean Dynamics of Outer Solar System Satellites", Geophys. Res. Lett. 46 (2019), 8700-8710.
- [4] Gastine T., Wicht J., and Aubert J., "Scaling regimes in spherical shell rotating convection", J. Fluid Mech. 808 (2016), 690-732.
- [5] Wang G., Santelli L., Lohse D., Verzicco R., and Stevens R.J.A.M., "Diffusion-free scaling in rotating spherical Rayleigh-Bénard convection", Geophys. Res. Lett. 48 (2021).
- [6] Vance S.D., Panning M.P., Stahler S., Cammarano F., Bills B.G., Tobie G., and Banerdt B., "Geophysical investigations of habitability in ice-covered ocean worlds", JGR: Planets 123 (2018), 180–205.
- [7] Stevens R.J.A.M., Clerx H., and Lohse D., "Optimal Prandtl number for heat transfer in rotating Rayleigh–Bénard convection", New J. Phys. 12 (2010), 075005.
- [8] King E.M., Stellmach S., and Aurnou J.M., "Heat transfer by rapidly rotating Rayleigh-Bénard convection", J. Fluid Mech. 691 (2012), 568.
- [9] Yang Y., Verzicco R., Lohse D., and Stevens R.J.A.M., "What rotation rate maximizes heat transport in rotating Rayleigh-Bénard convection with Prandtl number larger than one?" Phys. Rev. Fluids 5 (2020), 053501.

Thank you for your interest!

Acknowledgments

We acknowledge all resources supporting this work, namely: ERC Starting Grant *UltimateRB* No. 804283 for financial funding, PRACE for awarding access to MareNostrum4 at Barcelona Supercomputing Center (BSC), Spain and IRENE at Très Grand Centre de Calcul (TGCC) du CEA, France (project 2021250115).



Background: Ed Hawkins, showyourstripes.info

