

Decomposing Realistic Ocean Flow in Balanced and Unbalanced Parts Using a Novel Balancing Approach Silvano Rosenau¹, Manita Chouksey², Carsten Eden¹

1. Flow Decomposition



Why do we need flow decomposition?

• To understand and quantify energy transfer between different scales (e.g. dissipation of mesoscale eddies)

What high accuracy decomposition method exists?

- Nonlinear normal mode decomposition (Machenhauer, 1977)
- Optimal Balance (Viúdez et al., 2004; Masur et al., 2020)

2. Optimal Balance (OB)

- Transforms nonlinear decomposition problem into linear problem
- Linear problem: Projection onto linear geostrophic eigenspace
- Solve linear problem with spectral decomposition
- Spectral decomposition require idealized model setups:
 - \Rightarrow no topography, no lateral boundaries
 - \Rightarrow constant coriolis parameter.
- We replace spectral decomposition with time averaging

Existing high accuracy decomposition methods are limited to idealized models (e.g. no topography).

• We modify the "Optimal Balance" method, to make it applicable to realistic model setups

Fig. 1:

Decomposition of a 2D shallow water flow field in balanced and unfields using balanced optimal balance. The colors show the layer thickness.

Animation:

3. Modification of OB: Geostrophic projection with time averaging

Define cumulative linear time average:

$$\langle z \rangle_T = \frac{1}{T} \int_{t_0}^{t_0 + T} z \text{ stands for } u, v, w, \text{ etc.}$$

Linearized shallow water / non-hydrostatic boussinesq equations on f-plane can be decomposed into two modes^{*}:

$$z(\mathbf{k}, t) = \underbrace{z_{geo}(\mathbf{k})}_{\text{geostrophic mode}} + \underbrace{z_{wave}(\mathbf{k})}_{\text{inertia-gradients}}$$

\Rightarrow Time average converge to geostrophic mode:

$$\langle z \rangle_T = z_{geo} + \delta_T \cdot z_{wave}$$
 with

• Faster convergence with multiple time averages:



* There are actually two inertia-gravity wave modes which are here simplified to a single one

Theory:

balanced state

z(t)dt(1)linear evolution (no advection)

 $(\mathbf{k})e^{-i\omega(\mathbf{k})t}$ (2)

avity wave mode

$$\delta_T = \frac{i}{\omega T} \left(e^{-i\omega T} - 1 \right) \quad (3)$$

$$\langle \dots \langle z \rangle_{T_1} \dots \rangle_{T_n} = Z_{geo} + (\delta_{T_1} \cdot \dots \cdot \delta_{T_n}) Z_{wave}$$
 (4)

Fig. 2: The projection error is the norm of difference between the time averaged state and the geostrophic mode obtained with spectral decomposition. Experiments are per-formed in the shallow water model with the initial condition shown in Fig. 1.

 $\Rightarrow \mathcal{B}_{OBTA}$ converges to \mathcal{B}_{OB} when δ converges to zero \Rightarrow Apply \mathcal{B}_{OBTA} iteratively to reduce the balancing error

Numerical Experiments:

- Two different models
 - 2D shallow water model
 - 3D non-hydrostatic boussinesq model
- Different initial conditions (unstable jet, random phase)
- Rossby numbers ranging from 0.03 to 0.5



- All numerical results show general agreement with Eq. (5)



GitHub:



Fig. 3:

FRIDOM:

Ocean models

Framework for Idealized

The balancing error is the norm f difference between the balanced state of the corresponding method and $\mathcal{B}_{OB}(z)$ for ten iterations. Experiments are performed in the shallow water model with a Rossby number of 0.1. The number n refers to the number of time averages.

• Differences between $\mathcal{B}_{OBTA}(z)$ and $\mathcal{B}_{OB}(z)$ get extremely small

New method OBTA can be applied to models with topography and lateral boundaries! Solution of OBTA approaches solution of OB







