POTSDAM INSTITUTE FOR CLIMATE IMPACT RESEARCH

An Intermediate Complexity Approach to the Dynamics of Localized Extreme Heatwaves in the Mid-Latitude Atmosphere for moist-convective environments using Aeolus 2.0

Authors

Sullyandro Guimarães^{1,2}, Masoud Rostami^{1,3}, Stefan Petri¹

BACKGROUND AND CONCEPTS

• Examine the **evolution** of the **large-scale** localized For global dynamical simulations, the two-layer **Moistbuoyancy anomalies** in mid-latitude regions. **Convective Thermal Rotating Shallow Water (mcTRSW)** model Aeolus 2.0 with intermediate complexity was • Investigate the adjustments in the atmosphere for employed. The concept of two interacting layers moist-convective environments, comparing with enabled the study of the dynamics of **localized extreme** the dry case. heatwaves in baroclinic and barotropic situations. The model initialization comprises daily averaged velocity and potential temperature variables from **ERA5** data. onal Mear



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Contact Sullyandro Guimarães Telegrafenberg A62 | D-14412 Potsdam sullyandro@pik-potsdam.de

www.pik-potsdam.de

RESEARCH GOALS

The condensed liquid water content (**CLWC**) anomaly evolution shows that **baroclinic** localized buoyancy perturbation should play an important role for increased cloud formation and condensation, as a result of the heatwave propagation in the atmosphere for those extreme forcings.



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RESULTS AND CONCLUSIONS

Presence of a **circular positive** buoyancy anomaly in the **lower** layer, while the upper layer shows opposite circular rotation wind movement for some of the cases analyzed.

Comparing the **strong** and **weak** buoyancy anomalies results, we can notice the **prolonged effects** of **baroclinic** initial condition over the **barotropic** case.



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(1) Potsdam Institute for Climate Impact Research, Germany (2) University of Potsdam, Potsdam, Germany (3) Laboratoire de Météorologie Dynamique (LMD), Sorbonne University (SU), Paris, France









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Thermal Rotating Shallow Water (TRSW) Models

An extension of the classical shallow water equations, incorporating rotation (via Coriolis forces) and thermodynamics (typically via buoyancy or temperature/density variations). It captures the dynamics of a thin layer of fluid where horizontal scales are much larger than the vertical scale, and where **thermal effects drive motion** in addition to mechanical forces.

Key Features:

- Thermal Forcing: Allows buoyancy-driven flows (e.g., convection, density fronts) by linking the fluid layer thickness or internal energy to temperature/density fields.
- Wave Dynamics: Supports both barotropic (depth-independent) and **baroclinic** (depth-dependent, thermally influenced) modes.
- **Conservation Laws:** Satisfies energy, momentum, and potential vorticity conservation, critical for realistic long-term behavior.

Typical Applications:

- Modeling ocean thermocline dynamics.
- Studying weather system formation like cyclones and anticyclones.
- Investigating equatorial superrotation in planetary atmospheres.
- . Designing **reduced-order models** for climate prediction and large-scale flow.



• Atmospheric and Oceanic Dynamics: Models key features of weather systems, ocean fronts, equatorial waves, and large-scale circulations.

• Understanding Energy Transfer: Provides insights into how thermal energy converts into kinetic energy and vice versa in rotating systems.

What It Is Good For:

• Instability and Turbulence Studies: Analyzes **baroclinic instability**, an essential process for mid-latitude weather and ocean eddies.

Exoplanet and Planetary Atmospheres:

Simulates rotating, stratified atmospheres with differential heating, relevant for Earth, Jupiter, and exoplanets.

• Simplified Testing Ground: Acts as a computationally affordable "laboratory" to develop and test theories before moving to full 3D Navier-Stokes simulations.

Aeolus 2.0

Equations with Moist-Convection

 $(\boldsymbol{v}_1)_t + \boldsymbol{v}_1 \cdot \boldsymbol{\nabla} \boldsymbol{v}_1 + f \hat{\boldsymbol{z}} \wedge \boldsymbol{v}_1 = (\boldsymbol{v}_2)_t + \boldsymbol{v}_2 \cdot \boldsymbol{\nabla} \boldsymbol{v}_2 + f \hat{\boldsymbol{z}} \wedge \boldsymbol{v}_2 = (h_1)_t + \boldsymbol{\nabla} \cdot (h_1 \boldsymbol{v}_1) = \frac{1}{b_1} \gamma(C - C)$ $(h_2)_t + \boldsymbol{\nabla} \cdot (h_2 \boldsymbol{v}_2) = \frac{1}{h_2} \gamma(-C \cdot$ $(b_1)_t + \boldsymbol{v}_1 \cdot \boldsymbol{\nabla} b_1 = \frac{1}{h_1} (-C + D)$ $(b_2)_t + \boldsymbol{v}_2 \cdot \boldsymbol{\nabla} b_2 = \frac{1}{h_2} (C - \mu E)$ $Q_t + \boldsymbol{\nabla} \cdot (Q\boldsymbol{v}_2) = -C + E,$

Hamiltonian Str

 $\partial_t \varphi^a = \left\{ \varphi^a, \mathcal{H} \right\} = \mathbb{J}^{ab} \frac{\delta \mathcal{H}}{\delta \varphi^b}. \quad (\varphi) = (u_1, v_1, h_1, b_1, u_2, v_2, h_2, b_2)^T \text{ is the phase space.}$

H represents the Hamiltonian, and J denotes the Poisson tensor. Hamiltonian = total energy

The functional derivatives of the Hamiltonian H are:



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e Heatwaves in the	(1) Potsdam Institute for Climate Impact Research, Germany		
2.0	(2) University of Potsdam, Potsdam, Germany		
	(3) Laboratoire de Météorologie Dynamique (LMD) Sorbonne University (SU), Paris France		
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$-\left(b_1\nabla(h_1+Z)+b_2\nabla h_2+\frac{h_1}{2}\nabla b_1+h_2\nabla b_2\right)+\frac{v_2-v_1}{h_1b_1}\gamma(C-D),$	(1)	$(\partial_t + \boldsymbol{v}_1 \cdot \boldsymbol{\nabla}) \boldsymbol{v}_1 + f \hat{\boldsymbol{z}} \times \boldsymbol{v}_1 = -\langle \boldsymbol{\nabla} p_1 \rangle,$	(2.11a)
$-\Big(b_2\boldsymbol{\nabla}h_1+b_2\boldsymbol{\nabla}(h_2+Z)+\frac{h_2}{2}\boldsymbol{\nabla}b_2\Big),$	(2)	$(\partial_t + oldsymbol{v}_2 \cdot oldsymbol{ abla})oldsymbol{v}_2 + f \hat{oldsymbol{z}} imes oldsymbol{v}_2 = -\langle oldsymbol{ abla} p_2 angle - rac{1-\gamma}{b_2 h_2} (oldsymbol{v}_2 - oldsymbol{v}_1)$	$(\mathscr{C}-\mathscr{D}),$
D),	(3)	1	(2.11b)
+D),	(4)	$\partial_t h_1 + oldsymbol{ abla} \cdot (h_1 oldsymbol{v}_1) = \; rac{1}{b_1} [(1-\gamma)(-\mathscr{C} + \mathscr{D}) - (1-\gamma^{\mathbb{F}})]$	$\mathbb{F}_1],$
)).	(5)	1	(2.11c)
· / ,	(6)	$\partial_t h_2 + \mathbf{\nabla} \cdot (h_2 \boldsymbol{v}_2) = \frac{1}{b_2} [(1 - \gamma)(+\mathscr{C} - \mathscr{D}) - (1 - \gamma^{\mathbb{F}})]$	$\mathbb{F}_2],$
.),	(0)		(2.11d)
	(7)	$\partial_t b_1 + oldsymbol{v}_1 \cdot oldsymbol{ abla} \ b_1 = \ rac{1}{h_1} [(+ \mathscr{C} - \mu \mathscr{E}) + \mathbb{F}_1],$	(2.11e)
		$\partial_t b_2 + oldsymbol{v}_2 \cdot oldsymbol{ abla} \ b_2 = \ rac{1}{h_2} [(-\mathscr{C} + \mathscr{D}) + \mathbb{F}_2],$	(2.11f)
ucture		$\partial_t q_1 + oldsymbol{ abla} \cdot (q_1 oldsymbol{v}_1) = - \mathscr{C} + \mathscr{E},$	(2.11g)
$\frac{\mathcal{H}}{\mathcal{H}}$ (12) - (11, 12, h, h, h, 12, h, h, h) ^T is the phase space		$\partial_t q_2 + \boldsymbol{\nabla} \cdot (q_2 \boldsymbol{v}_1) = + \mathscr{C} - \mathscr{D}.$	(2.11h)

$$\mathcal{H} = \int_{\mathcal{D}} d^2 x \left[h_1 \left(\frac{1}{2} \boldsymbol{v}_1^2 + \tilde{h_1} b_1 \right) + h_2 \left(\frac{1}{2} \boldsymbol{v}_2^2 + \tilde{h_2} b_2 \right) \right]. \qquad \qquad \mathcal{E}_{\text{vap}}(T_1)$$
the Hamiltonian H are:

$$\frac{\delta \mathcal{H}}{\delta \boldsymbol{v}_k} = h_k \boldsymbol{v}_k, \quad \frac{\delta \mathcal{H}}{\delta h_k} = \zeta_k, \quad \frac{\delta \mathcal{H}}{\delta b_k} = h_k \tilde{h}_k.$$

 $\mathcal{E}_{\mathrm{vap}}(|u_1^n|)$

Fr. Conv.



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Equations adding Forcings

Sea surface evaporation scheme

$$= \mathbb{H}(Q^s - q_1) \exp\left[\frac{-\Delta H_{\text{vap}}}{R_v} \left(\frac{1}{T_1^{\alpha}} - \frac{1}{T_0^{\alpha}}\right)\right],$$

$$= \mathbb{H}(Q^s - q_1) \exp\left(\frac{|u_1^n|^{\alpha_{1v}}}{\alpha_{2v}}\right),$$

$$= \mathbb{H}(Q^s - q_1) \mathcal{A}_F(Q^s - q_1), \quad \text{if}(|u_1^n| \ll 1 \text{ and } b_1 \ll b_1^{\max}),$$

 $\mathcal{E}(T_1, |\boldsymbol{u}_1^n|, q_1) = \mathbb{H}(Q^s - q_1) \left(\mathcal{A}_T \, \hat{\mathcal{E}}_{\text{vap}}(T_1) + \mathcal{A}_u \, \hat{\mathcal{E}}_{\text{vap}}(|\boldsymbol{u}_1^n|) + \text{Fr. Conv.} \right),$