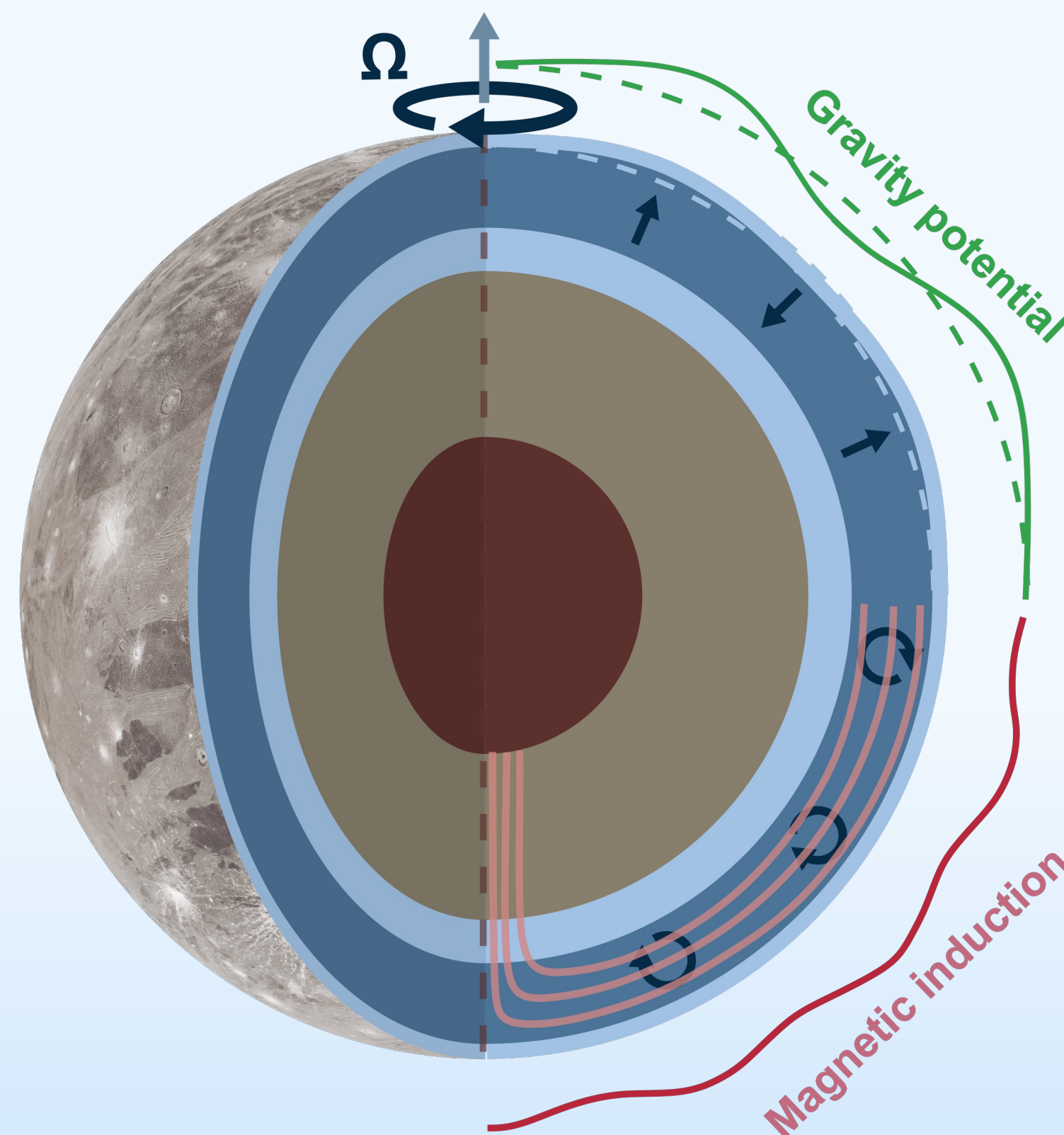


Motivation

- Ganymede has a subsurface ocean** as supported by several evidences including magnetic measurements from Galileo and Juno^a, and auroral observations from the Hubble Space Telescope^b.
- Jupiter's tides can drive periodic flows within the ocean.**
- Tidal flows can induce signals** which are in principle observable. On Earth, satellites can measure the magnetic signature of ocean tidal flows^c.
- Juice is on its way to Ganymede!** The low-altitude orbital phases of the mission (GCO500 and GCO200)^{d,e} could make it possible to detect very weak signals.
- Detection of tidal flows signatures can lead to a better understanding of Ganymede's ocean dynamics and to additional constraints on the moon's interior.**

Question

Can Juice detect the signatures in magnetic and gravity fields induced by Jupiter's eccentricity and obliquity tides, and how could such detections improve constraints on Ganymede's interior?



Summary

Conclusion

Flow

- Tides can excite resonant internal inertial waves within the ocean.
- Hitting a resonance can amplify the response enough to make the flow turbulent.
- Westward obliquity tide excite a resonant Rossby mode. Other tides can hit a resonance depending on the ocean geometry.

Signatures

- Signatures of eccentricity tides appear too small to be detected by Juice.
- Both **magnetic and gravitational signatures of the westward obliquity tide exceed the resolution of Juice's instruments** for a certain range of parameters.
- Although the intensity of the signals are above the resolution of the instruments, separation from ambient gravity and magnetic fields remains challenging.

Characterization

- Successful detection can give additional constraints on ocean depth, ice shell thickness, ocean conductivity and effective viscosity.**

Future studies

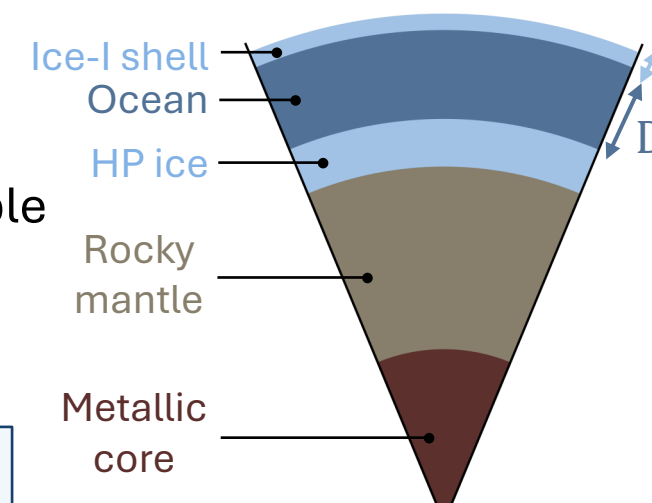
- To overcome the main limitations of the present work, future studies should:
- Include nonlinearities and a better handling of turbulence.
 - Dynamically couple the ocean flow with the elasto-gravitational deformation of the moon in a geometry that is not perfectly spherical.
 - Explore the coupling with other forcing mechanisms such as convection.

Method

Deformation

Step 1. Jupiter's tides deform Ganymede

- Method:** we solve the classical **elasto-gravitational equations**^f to compute degree 2 static tidal **Love numbers**.
- Model and Assumptions:**
 - 5-layer interior model of Ganymede.
 - Homogeneous, non-rotating, and incompressible layers.
 - Perfectly elastic solid layers and inviscid liquid layers.
- Output:** **Radial deformation** at the top of the ocean



$$\xi = h_2 \frac{\Phi_T}{g_0}$$

Eccentricity ($m = 0$)

Obliquity ($m = 1$)

Eccentricity ($m = 2$)

$\xi(E/W) = 1.3 \text{ m}$

$0.4 \text{ m} / 0.4 \text{ m}$

$3.8 \text{ m} / 0.5 \text{ m}$

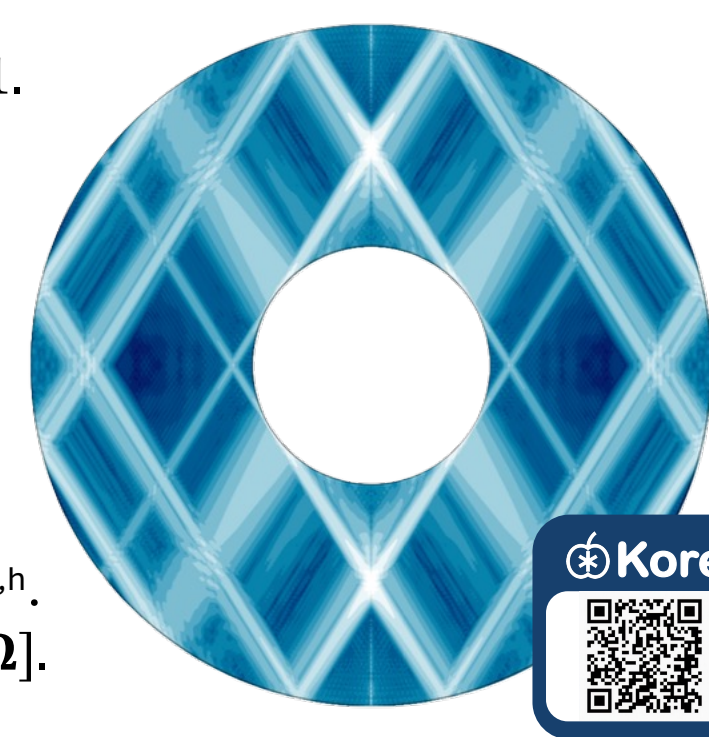
Tidal flow

Step 2. The periodic deformation drives the flow

- Method:** we solve the **linearized Navier-Stokes equations** in the ocean by imposing the displacement at the interface computed in Step 1 as **radial boundary forcing**.
- Model and Assumptions:**
 - Perfectly spherical shell model of the ocean.
 - Newtonian, homogeneous, rotating and incompressible fluid.
 - Molecular Ekman number $Ek = \nu/\Omega R^2 \sim 10^{-14}$.
 - Linearization:** Rossby number $Ro = U/\Omega R \ll 1$.
- Numerical implementation:**
 - Spectral solver: **Kore**.
 - Orders m are decoupled ($m = 0, 1, 2$).
 - Numerically accessible Ekman number up to $Ek = 10^{-10}$.
 - Aspect ratio $\eta = (R - D)/R \in [0.8 - 0.99]$.
- Tidal flow:**
 - Tides can excite linear internal **inertial modes**^{g,h}.
 - Inertial modes have frequencies $\omega \in [-2\Omega - 2\Omega]$.
 - Characterized by intense **shear layers**.

$$i\omega \mathbf{u} + 2\mathbf{e}_z \times \mathbf{u} = -\nabla p + Ek \nabla^2 \mathbf{u}$$

$$\mathbf{u}(r = R) = i \Omega \xi_{\text{tides}}$$



Signatures

Step 3. The tidal flow induce signals in gravity and magnetic fields

- Magnetic field**
 - Motion of a conductive ocean ($\sigma \in [0.1 - 5] \text{ S/m}$) in a background magnetic field **induce a secondary magnetic field**.
 - Linearized induction equation** with insulating BC

$$i\omega \mathbf{b} = \nabla \times (\mathbf{u} \times \mathbf{B}_0) + Em \nabla^2 \mathbf{b}$$

$$\text{with } Em = 1/\mu_0 \sigma \Omega R^2$$

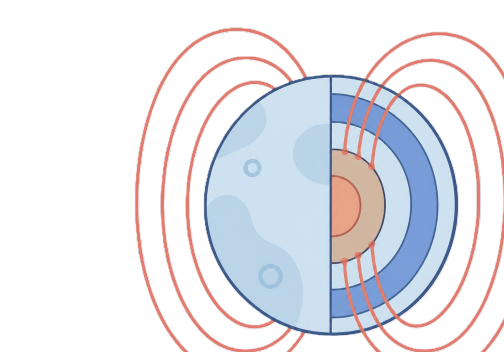
- Extrapolation towards **Juice GCO200** altitude.

2) Gravity field

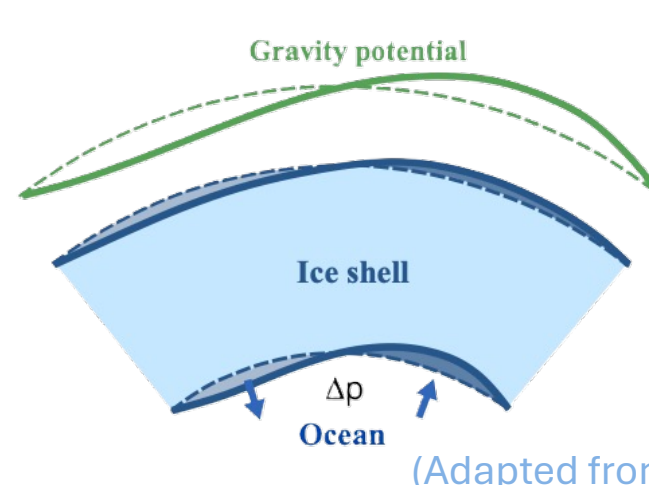
- Global **elasto-gravitational deformations** due to **pressure field variations** associated with tidal flows at the ice-ocean interface.

- Pressure Love numbers**ⁱ computed from the elasto-gravitational equations with the 5-layer Ganymede model.

$$\Phi_{lm}^p = -\frac{k_l^p}{\rho} P_{lm}$$

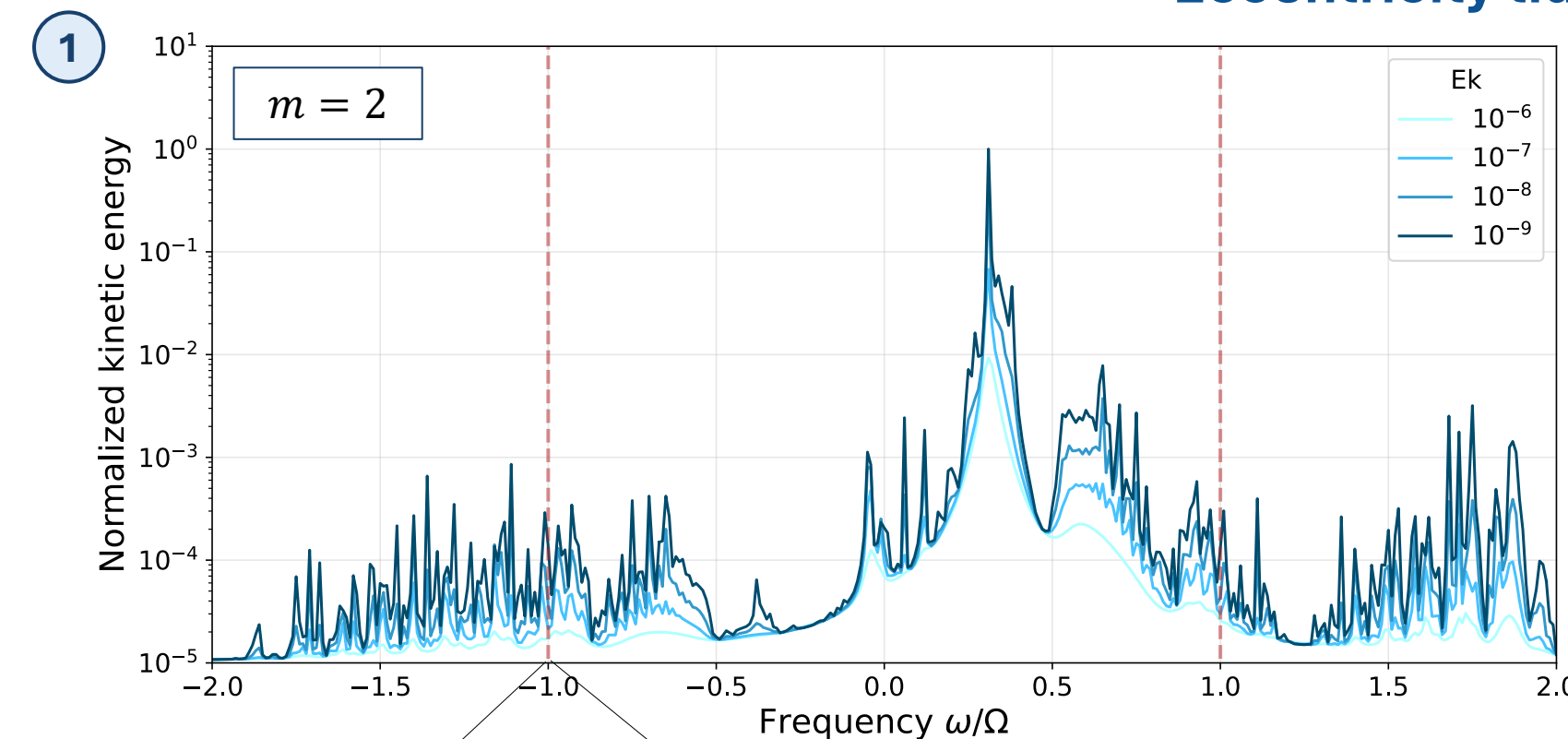


$$\mathbf{B}_0 = \nabla \times \nabla \times \left[\frac{B_0}{2r^2} Y_1^0(\theta, \varphi) \mathbf{r} \right]$$

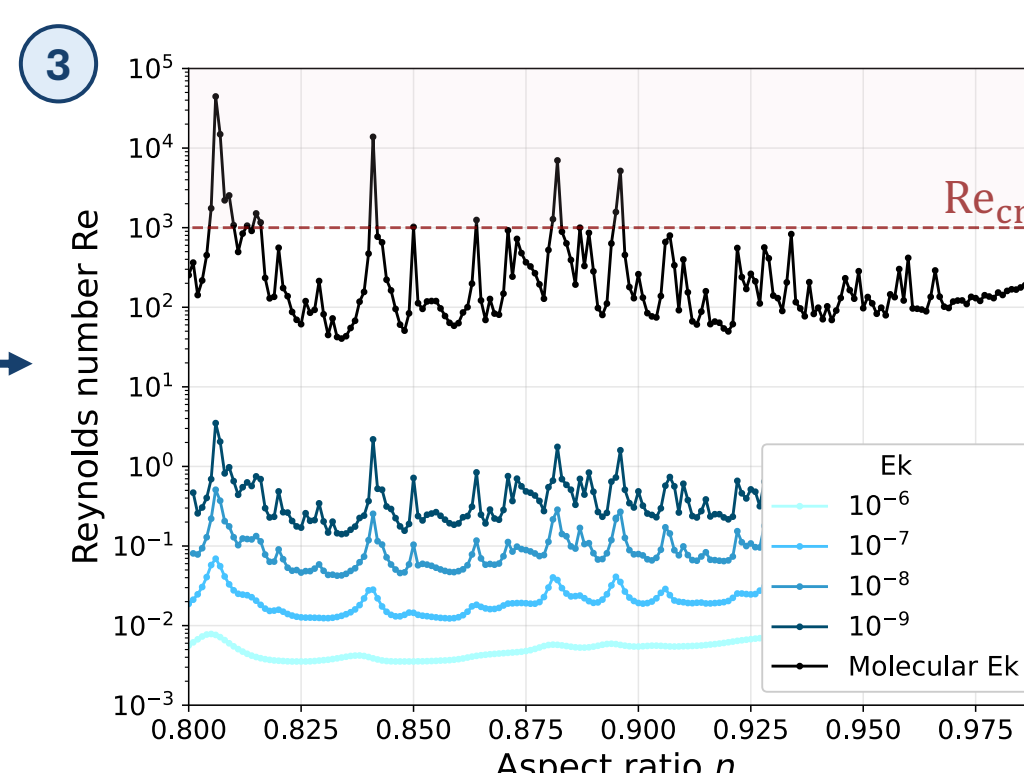
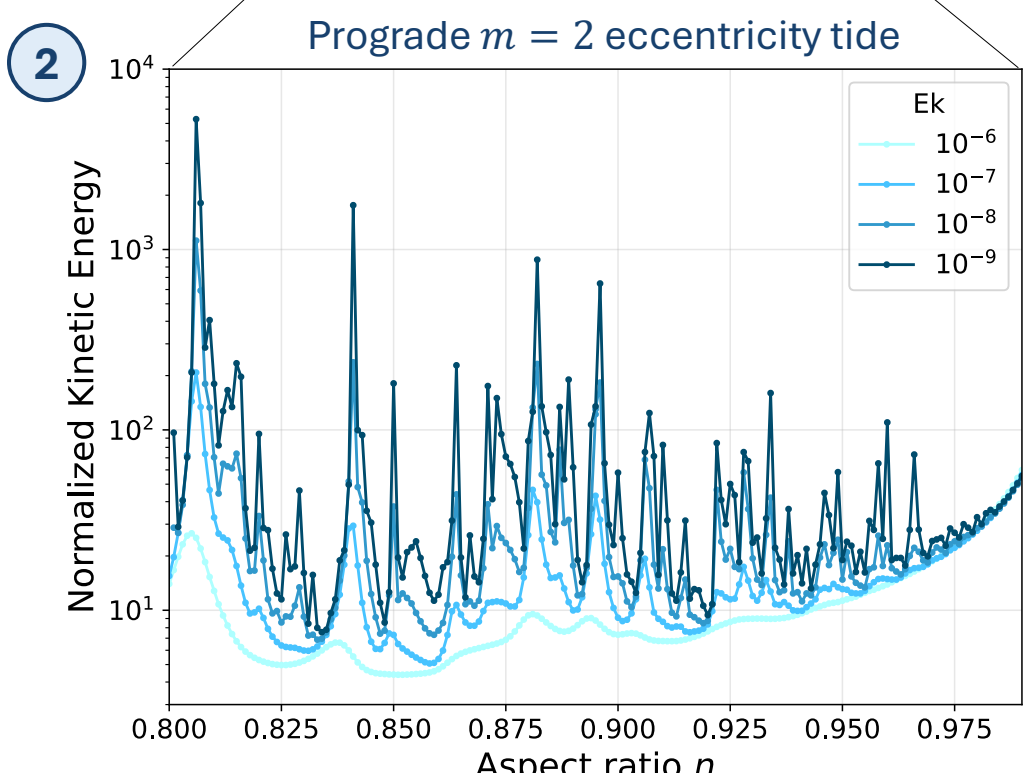


Results and Discussion

Eccentricity tides

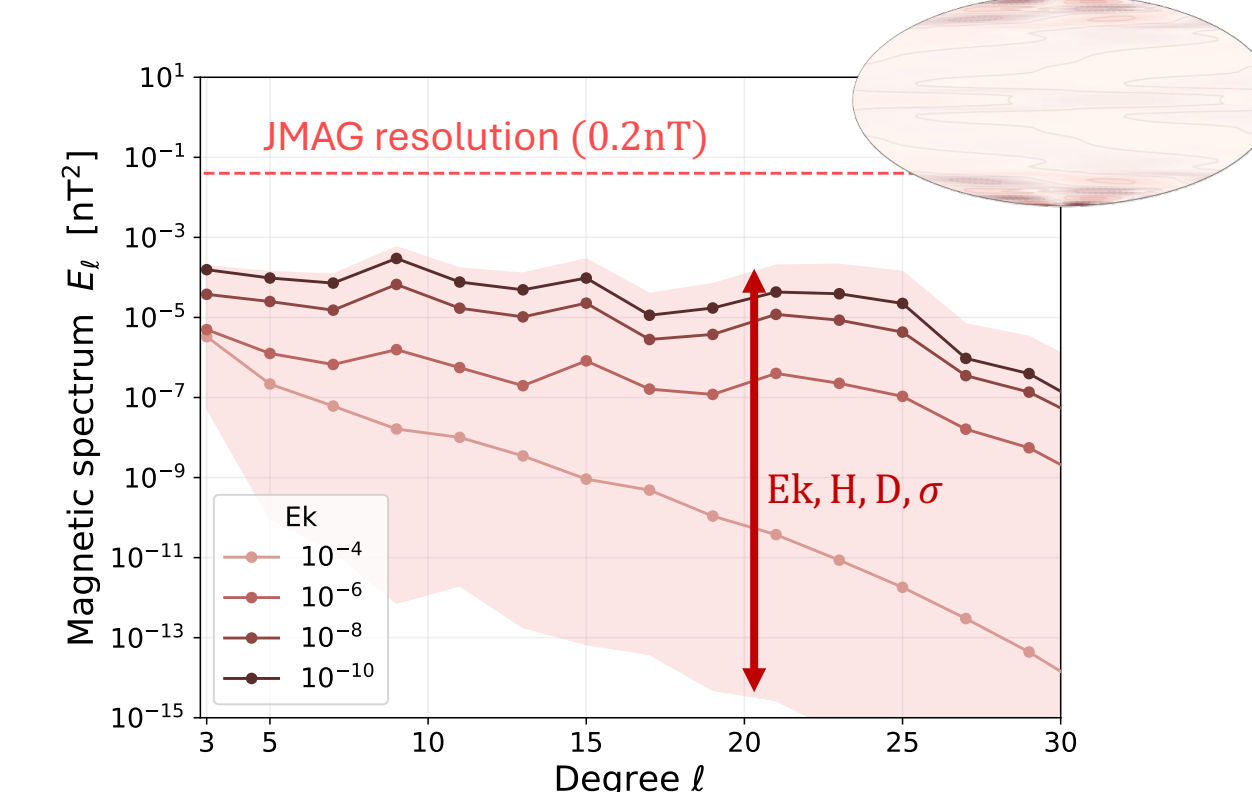
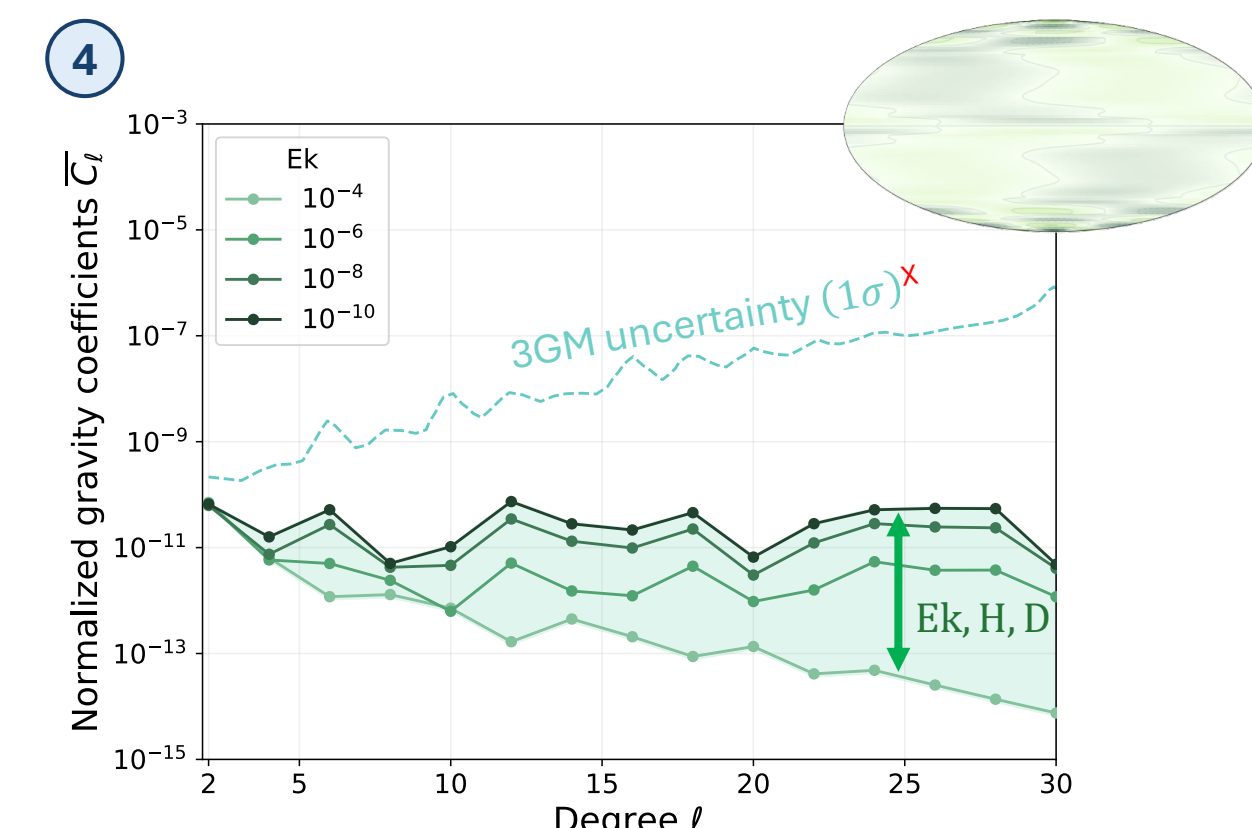


- Frequency spectrum of the response**
 - $m = 2$ forcing ($m = 0$ not shown but equivalent).
 - Spiky due to **inertial modes** excitation.
 - Eccentricity tides at diurnal frequency ($\omega = \pm \Omega$).
- Spectrum of the prograde eccentricity tide**
 - Inertial modes excitation if the aspect ratio is **resonant**.
 - Resonance amplify the flow intensity.
 - Amplification increase with decreasing Ek . Shear layers become thinner and more intense.
- Reynolds number**
 - If the amplification is sufficient, the flow may become **turbulent**. We compute the **Reynolds number** and extrapolate towards molecular Ek to assess the fluid regime (chosen critical value: 10^3).



The flow becomes turbulent

- Linearity assumption breaks
- Turbulence modelled by an effective viscosity
- $Ek^{\text{eff}} \in [10^{-4} - 10^{-10}]$ to account for the uncertainty



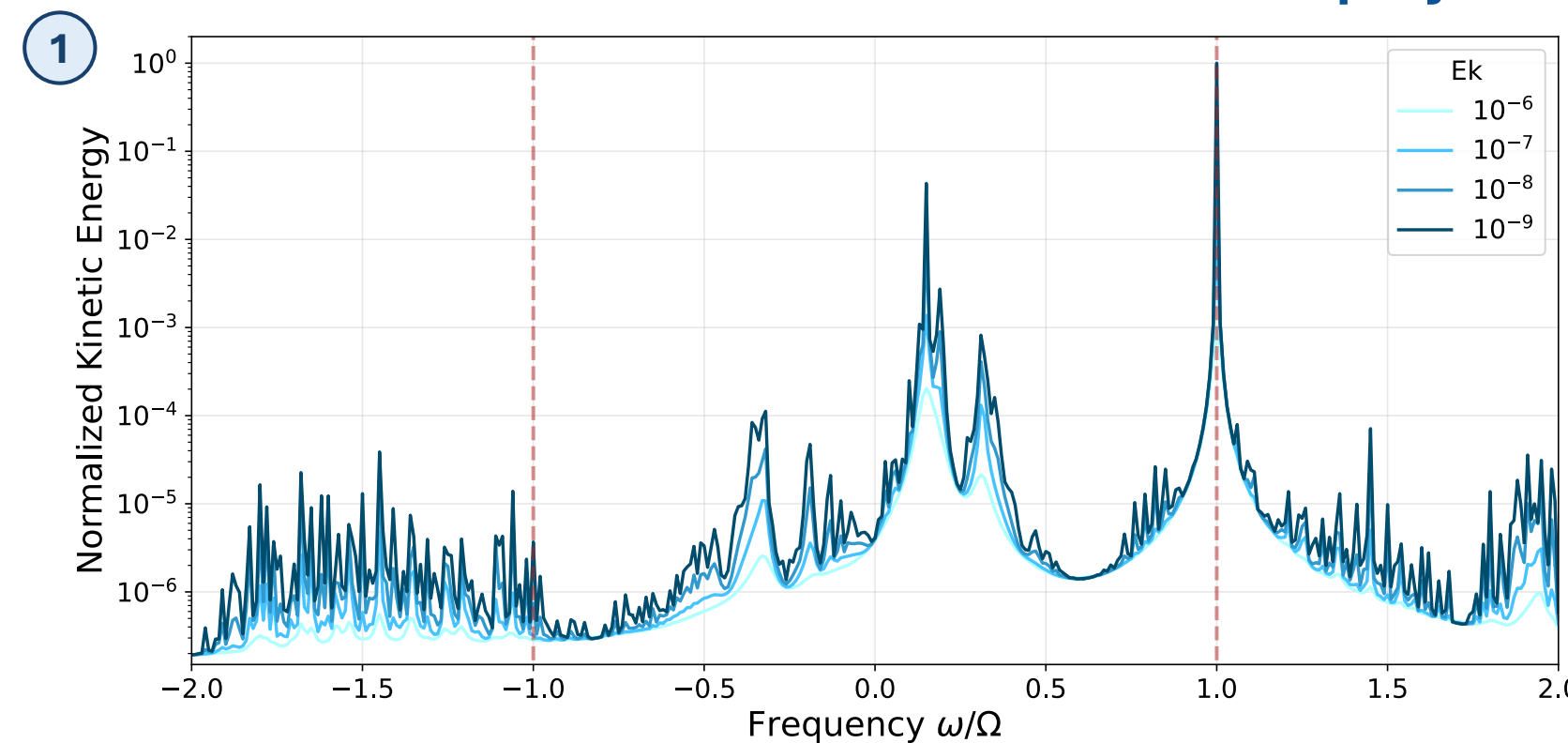
4. Signatures

- Magnetic and gravity signatures of tidal flows compared to Juice's instruments expected accuracy
- Chosen mode $\eta = 0.806$ (Other modes are not shown but the conclusion remain identical).
- The presence of small-scale flow structure associated with shear layers lead to high-degree components
- Shaded area shows signal intensity variability with effective viscosity, ice shell thickness, ocean depth and ocean conductivity
- Signal intensity remain below the detectability threshold for the entire parameter space**

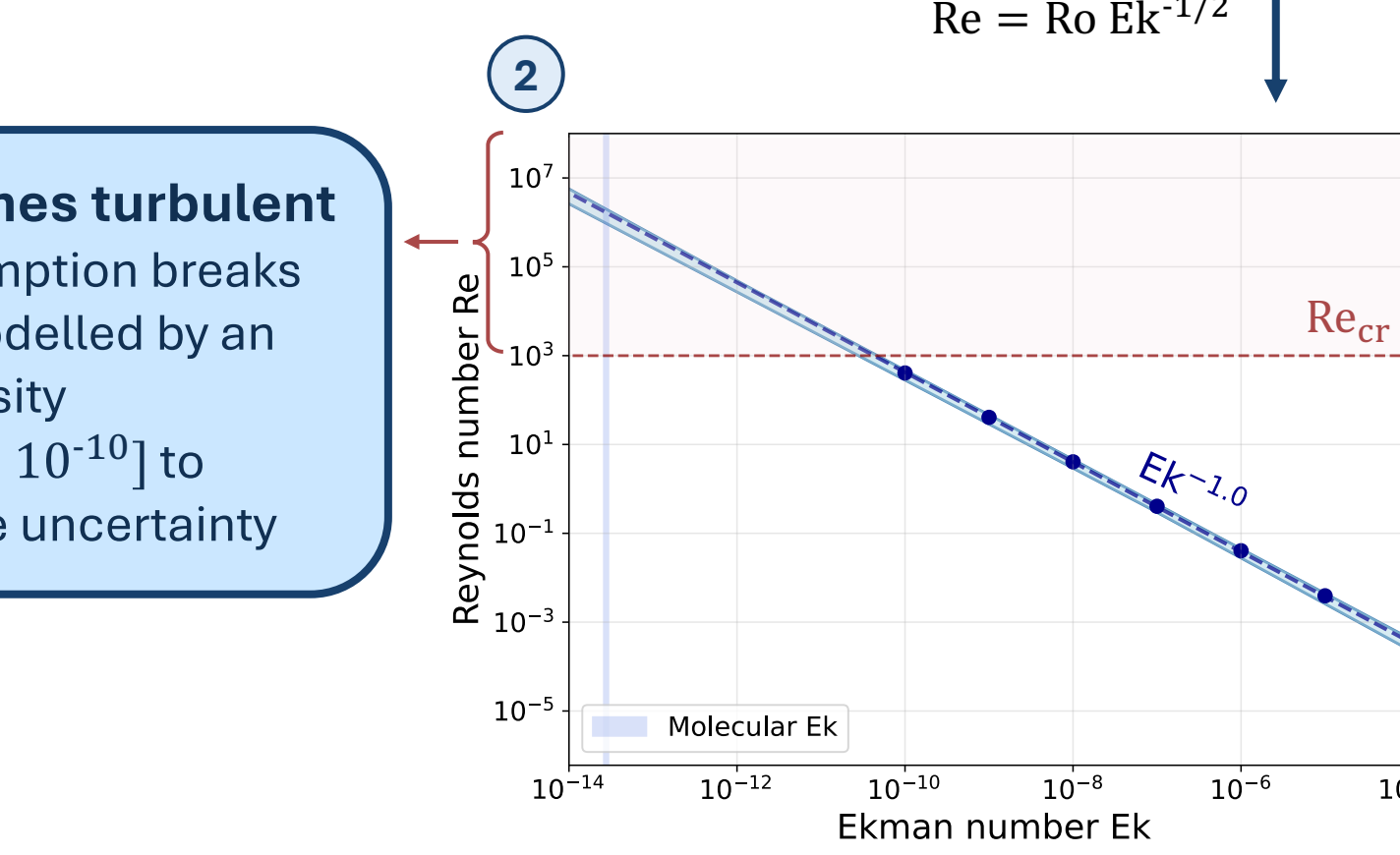
Undetectable signals

- Both magnetic and gravity signals remain below the resolution of Juice's instruments
- Signals associated with eccentricity tides appear too weak to be detected

Obliquity tide

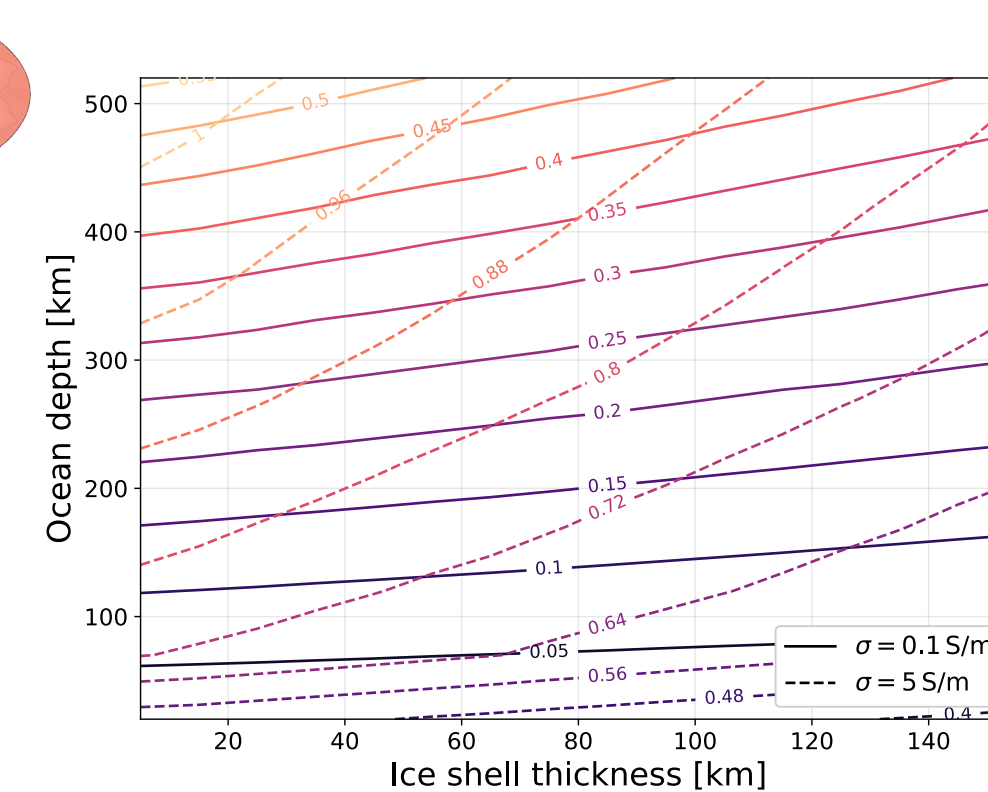
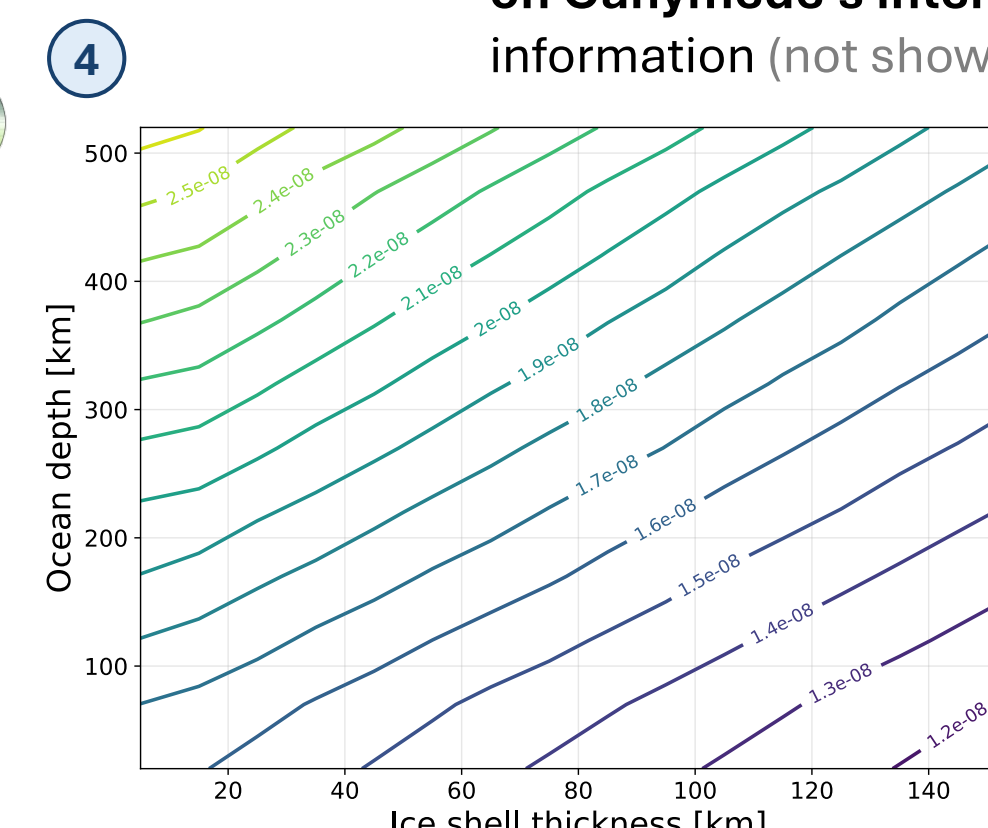
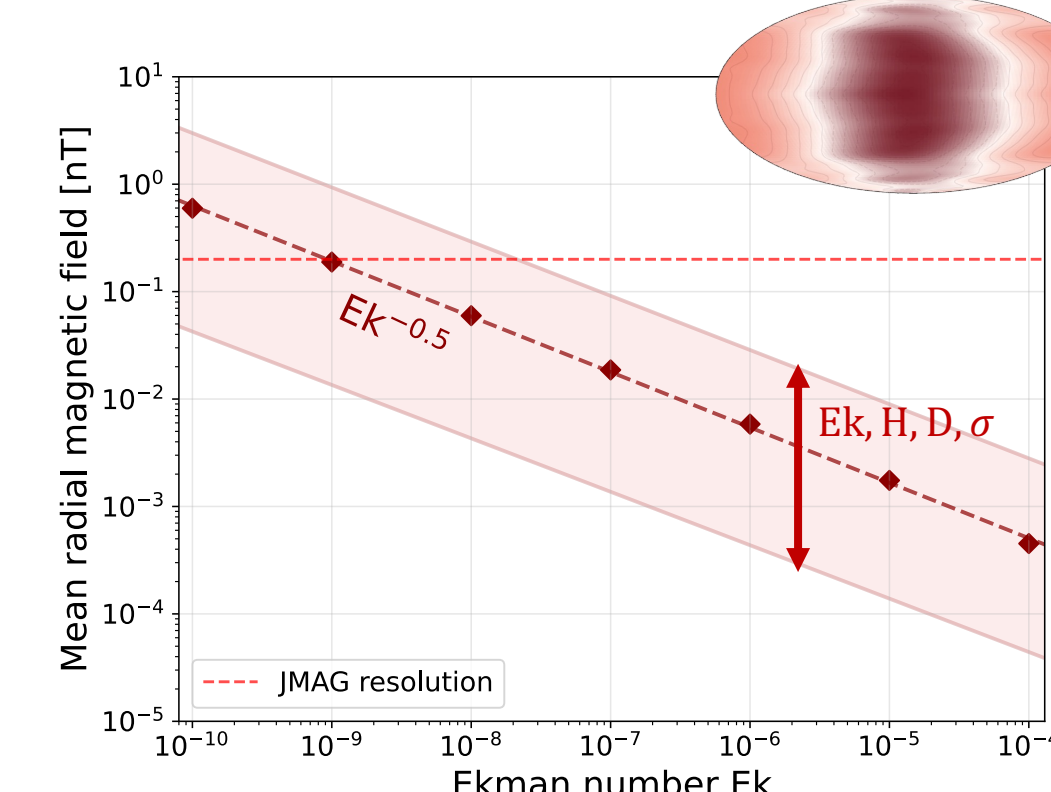
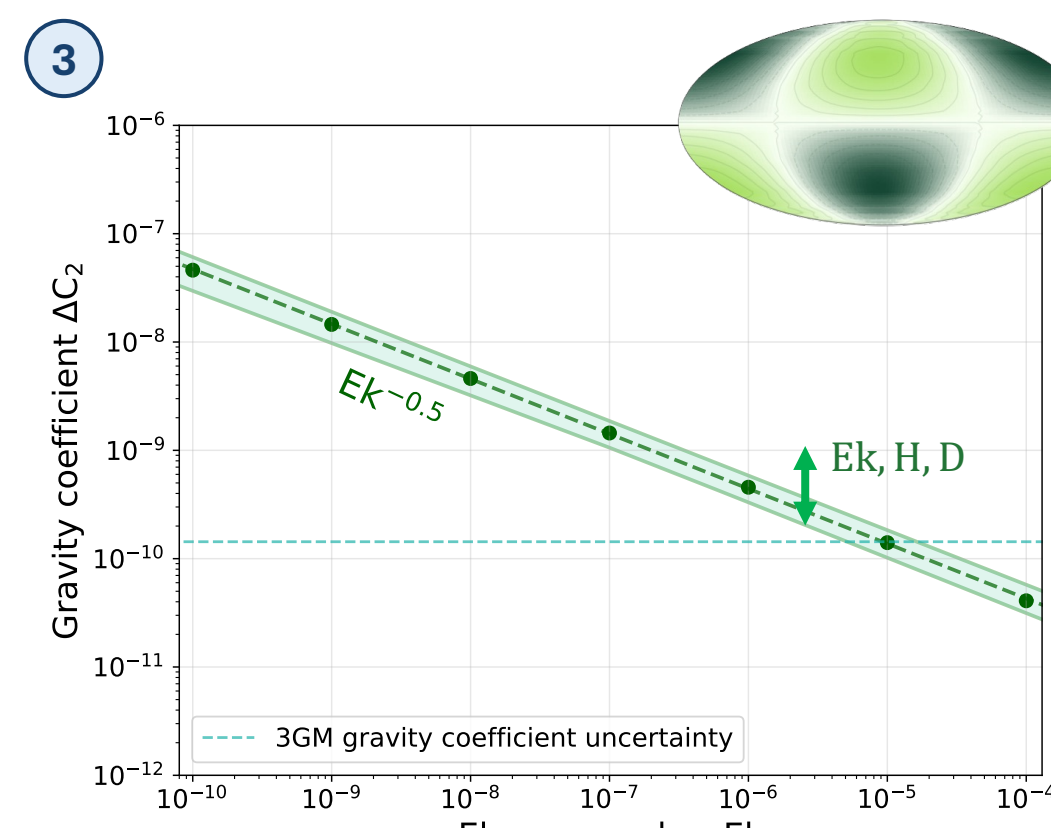


- Frequency spectrum of the response**
 - $m = 1$ forcing
 - Spiky due to **inertial modes** excitation.
 - Large resonance at $\omega = \Omega$: spin-over mode
 - Spin-over mode:** Analogous to the purely toroidal Rossby-Haurwitz mode of degree and order 1^l.
 - The westward obliquity tide excite the spin-over mode**
- Reynolds number**
 - Resonant amplification of the flow intensity. The Rossby number scales as $Ek^{-1/2}$ reflecting that the response is dominated by **Ekman layers** (whose thickness scales as $Ek^{1/2}$). Consequently, the Reynolds number scales as Ek^{-1}
 - Extrapolation towards molecular Ek show that the flow becomes **turbulent**



3. Signatures

- Signatures associated with the spin-over are of low-degree
 - The gravity signal is mainly of degree 2.
 - The magnetic signal is mainly of degree 1 (dipolar).
- The intensity of the signals scales as $Ek^{-1/2}$, just like Ro .
- For a certain range of parameters, **the signals exceed the resolution of the instruments**.
- Characterization**
 - Dependence of the signals intensity to Ganymede's physical parameters. Shown for $Ek = 10^{-9}$.
 - A successful detection can lead to **additional constraints on Ganymede's interior**, both from intensity and phase information (not shown).



Detectable signals

- Both gravity and magnetic signatures appear possibly detectable.
- Intensities at $Ek = 10^{-10}$ must be taken as upper bounds
- Successful detection gives constraints on Ganymede's interior
- Detection is challenging as the low-degree signals must be separated from the ambient fields.
- Limitation: the perfect frequency match between the mode and the tide could be affected in a not-perfectly spherical geometry^k.
- Limitation: the constant effective viscosity turbulence model is overly simplistic.

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Abstract information

